## Unit-1

## Theory of Elasticity \& Functional Approximating Methods:

Introduction to Theory of Elasticity: Definition of stress and strain - plane stress - plane strain - stress strain relations in three dimensional elasticity.
Introduction to Variational Calculus: Variational formulation in finite elements - Ritz method Weighted residual methods - Galerkin - sub domain - method of least squares and collocation method - numerical problems
One Dimensional Problems: Discretization of domain, element shapes, discretization procedures, assembly of stiffness matrix, band width, node numbering, mesh generation, interpolation functions, local and global coordinates, convergence requirements, treatment of boundary conditions. Steady state heat transfer analysis : one dimensional analysis
TEXT BOOKS:

1. An introduction to Finite Element Method / JN Reddy / McGraw Hill
2. The Finite Element Methods in Engineering / SS Rao / Pergamon.

## References

1. Tirupathi R. Chandrupatla and Ashok D. Belugundu (2011) Introdution to Finite Elements in Engineering, Prentice Hall.
2. Seshu P., Text Book of Finite Element Analysis, Prentice Hall, New Delhi, 2007.
3. Zienkiewicz O.C., Taylor R.L., Zhu J.Z. (2011), The Finite Element Method: Its basis and fundamentals, Butterworth Heinmann.

## E-RESOURCES: https://nptel.ac.in/courses/112/104/112104193/

https://mecheng.iisc.ac.in/suresh/me237/feaNotes

## Unit-2

Analysis of Trusses:Finite element modelling, coordinates and shape functions, assembly of global stiffness matrix and load vector, finite element equations, treatment of boundary conditions, stress, strain and support reaction calculations.
Analysis of Beams: Element stiffness matrix for Hermite beam element, derivation of load vector for concentrated and UDL, simple problems on beams.

## Unit-3

Two Dimensional Problems: Finite element modelling of two dimensional stress analysis with constant strain triangles CST and treatment of boundary conditions, Higher order and isoparametric elements: Two dimensional four noded isoparametric elements and numerical integration.
Axisymmetric Problems: Formulation of axisymmetric problems.

## Unit-4

Dynamic Analysis: Formulation of finite element model, element consistent and lumped mass matrices, evaluation of eigen values and eigen vectors, free vibration analysis.Steady state heat transfer analysis: one dimensional analysis of a fin.
Introduction to FE software.

- Numerical method for engineering solution.
- Reduces the degrees of freedom from infinite to finite


## Element

- All of the calculations are made at a limited number of points known as nodes.
- The entity joining nodes and forming a specific shape such as quadrilateral or triangular is known as an Element.
""FEM is a numerical technique to find the approximate solutions of partial differential equations. It was originated from the need of solving complex elasticity and structural analysis problems in Civil, Mechanical and Aerospace engineering.""


## Difference between FEM and FEA ??

Finite Element Method (FEM) involves complex mathematical procedures (like a theory manual, lots of equations and mathematics).

Finite Element Analysis (FEA) involves applying FEM to solve real world/ engineering problems.
https://www.amazon.com/DonaldsPractical-Stress-Analysis-Elements-Hardcover-d p-B005KZ2UJE/dp/B005KZ2UJE/ref=mt_other?_encoding=UTF8\&me=\&qid=

In a structural siencelation. FEM helps in producing stiffness \& strength visualizations. H also helps to minimize for material weight it and its cost of the structures. FEM allows for detailed visualization and indicates the distribution of stresses and strains inside the body l of a structure.

In structural analysis, FEM helps in producing stiffness and strength visualizations. It helps in minimising the material and cost of the structures.

Modern FEM packages (ansys, abaqus ), include specific components such as Fluid, thermal, EM and structural working environments

Several modern FEM packages include specific Components sect as fluid, therthal, electromageretic, and structural working environments. FEM allows entire design\$ to be constructed refined and optimized before the design is manufactured.

Numerical Methods which are colmently cased to solve solid and fluid mechanics problems are given bebw.

1. Finite Difference Meftiod
2. Firivie Volume Method.
3. Finite Element Method.
4. Boundary dlenrent Mahod.
5. Meshes Method.
$\rightarrow$ Concepts of elements and Nodes: $\qquad$ Element Library
Ore Dinamisoan (ID) Ekemens of pieces with very small dimensions. These small pieces/ sub domains of finite dimension axe called "Finite Elements". These elements are connected through a number of joints cotich axe called. 'Nodes'. Whtrile disciretizing' the structural systems, it is assumed that the elements are attached to the adjacent elements only at the nodal points. Each, element contains the material and geometrical properties. The material properties inside an element are assumed to be 1, constant. The elements nay be 10,20 $\varepsilon, 30$ elements.




Bulk stress and strain


Bulk or hydrostatic stress, also known as volumetric stress is a component of stress which contains uniaxial stresses, but not shear stresses. A specialized case of hydrostatic stress, contains isotropic compressive stress, which changes only in volume, but not in shape.

$$
\begin{array}{cc}
\text { Bulk } \\
\text { modulus }
\end{array} B=\frac{\Delta p}{\Delta V / V_{0}} \text { stress } \quad \text { strain } \quad k=\frac{1}{B}
$$

$\sigma_{h}=\frac{I_{i}}{3}=\frac{\sigma_{x x}+\sigma_{y y}+\sigma_{z z}}{3}$
Hydrostatic stress is equivalent to the average of the uniaxial stresses along three orthogonal axes

* Siltizan,:
is knowen as strain dhatge in dimension,

$$
\text { Strair }=\frac{\text { ohagge in dionension }}{\text { Original donension }}=\frac{\Delta L}{L}(\text { No Unis) }
$$



Eps v=Eps x+Eps y+Eps z

For cylinder-----------?
And sphere ...........?
$*$ Relation between Stress strain st shear \$tress, e shear strain:


$$
1 \quad \begin{aligned}
& \sigma \alpha \in \\
& \sigma=E \in
\end{aligned}
$$

1. Here $E=$ Modulus of elasticity or Young's Modulus. Q2) Stieax $y^{\prime}=\begin{aligned} & \text { stress (Normal), } \in \text { Normal strain. } \\ & 1\end{aligned}$

(3) Bulk o Volume

$$
\tau=G^{\gamma}
$$

$$
\begin{array}{r}
1 \\
\sigma_{v} \propto \epsilon_{v} \\
\sigma_{v}=k \epsilon^{\prime}
\end{array}
$$

Here $\tau=$ shear stress, ? shear strain
$G=$ shear Modulus ot Modulus of Rigidity
Here $\sigma_{v}=$ Bulk Stress, $\epsilon_{v}=$ Balk, Strain $1 / y$

$$
k_{v}=\text { Bulk Modulus. }
$$

So. Here Ext G., K, \& u (Parson's Ratio) are called elastic. constants because, the values of the above doesnit change with the change in size or shape, but changes with the material. In

Equilibrium of an Elastic Body : Consider a body with surface ' $S$ ' occupying a volume ' $V$ ', in which points are located by a coordinate system ( $x, y$ and $z$ coordinates) shown in figure (1). At particular part, boundary is constrained. Some other region of boundary is subjected to traction ' $T$ ' (uniformly distributed force or load per unit area). Deformation of the body occurs due to force and is specified at constrained part of the boundary. Displacement of any point $x=(x, y, z)^{T}$ is given by $u=(u, v, w)^{T}$.


Figure (1)

Types of forces acting on the body are surface loads (friction, viscous drag) which exists whenever one body moves past to other body in contact $T=\left(T_{x}, T_{y z}, T_{z}\right)^{T}$, body loads (forces distributed on volume of body like self weight, inertia, centrifugal forces, temperature, etc.) $f=\left(f_{x}, f_{y}, f_{z}\right)^{T}$ and point loads (loads concentrated on a point in continuum like tensile or compressive loads) $P_{i}=\left(P_{x}, P_{y}, P_{z}\right)^{T}$.


$$
\begin{equation*}
\frac{\partial \tau_{x z}}{\partial x}+\frac{\partial \tau_{y z}}{\partial y}+\frac{\partial \sigma_{z}}{\partial z}+f_{z}=0 \tag{3}
\end{equation*}
$$

(ii) Strain-displacement Relations

Considering the deformation of the $d x-d y$ face,


Figure (3): Deformed Elemental Surface
Then, corresponding strains are given by,

$$
\varepsilon=\left[\varepsilon_{x^{2}}, \varepsilon_{y,}, \varepsilon_{z^{2}}, \gamma_{z^{2}}, \gamma_{z^{2}}, \gamma_{x x^{2}}\right]^{T}
$$

Where,

$$
\begin{aligned}
& \varepsilon_{x}=\varepsilon_{y}, \varepsilon_{z}-\text { Normal strains } \\
& \gamma_{y z}, \gamma_{x z}, \gamma_{x y} \text { - Engineering shear strains }
\end{aligned}
$$

## Figure |2): Equilibriun of Elanemtal Volum *

There are three sets of equations of equilibrium in theory of elasticity, they are,
(a) Differential equations of equilibrium.
(ii) Strain-displacement relations.
(iii) Stress-strain relations.
(1) Differential Equations of Equilibrium

For elemental volume $a \mathrm{~F}, 3 \times 3$ symmetric matrix is used to specify the components of the stress tensor. Bat for instant six independent components are employed to represent the stress.

$$
\sigma=\left[\sigma_{s}, \sigma_{,}, \sigma_{N}, \tau_{y s} \tau_{v z} \tau_{s}\right]^{z}
$$

Where,

$$
\begin{aligned}
& \sigma_{y T} \sigma_{y}, \sigma_{z}-\text { Normal stresses } \\
& \tau_{8 \pi} ; \tau_{z s}, \tau_{y y}-\text { Shear stresses }
\end{aligned}
$$

Considering equilibrium of an elemental volume, $\Sigma f_{2}=$ $0, \Sigma_{y}=0, \Sigma_{z}=0$ and considering $d V=d a d y t$, and by using, forces are obtained by multiplying stresses with their corresponding areas, equilibrinm equations are given by;

$$
\begin{equation*}
\frac{\partial r_{\tau}}{\partial x}+\frac{\partial \tau_{s y}}{\partial y}+\frac{\partial \tau}{\partial z}+f_{\pi}=0 \tag{1}
\end{equation*}
$$

$\frac{\partial \tau_{x y}}{\partial x}+\frac{\partial \sigma_{y}}{\partial y}+\frac{\partial \tau_{\varepsilon}}{\partial z}+f_{y}=0$

Considering all the faces for small deformations.

$$
\varepsilon=\left[\frac{\partial u}{\partial x^{\prime}} \frac{\partial v}{\partial y} \cdot \frac{\partial w}{\partial z} \cdot \frac{\partial v}{\partial z}+\frac{\partial w}{\partial y} \cdot \frac{\partial u}{\partial z}+\frac{\partial w}{\partial x} \cdot \frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}\right]^{t}
$$

(iii) Stress-strain Relations

Applying Hooke's law to elemental vohame ' $d V$, stramis in terms of stress and material properties of isotropic maternals (elastic modulus and Poisson's ratio) are obtained as,

$$
\begin{aligned}
& \Sigma_{x}=\frac{\sigma_{x}}{E}-\frac{\sigma_{y}}{E} \mu-\frac{\sigma_{z}}{E} \mu \\
& s_{y}=-\frac{\sigma_{z}}{E} \mu+\frac{\sigma_{y}}{E}-\frac{\sigma_{z}}{E} \mu \\
& \varepsilon_{x}=-\frac{\sigma_{x}}{E} \mu-\frac{\sigma_{y}}{E} \mu+\frac{\sigma_{z}}{E} \\
& \gamma_{y=}=\frac{\tau_{y z}}{G} \\
& y_{x=}=\frac{\tau_{z}}{G} \\
& y_{x}=\frac{\tau_{x}}{G}
\end{aligned}
$$

Where,

$$
\begin{aligned}
& E \text { - Young's modulus } \\
& \mu \text { - Possion'a ratio } \\
& \text { G-Modulus of rigidity }
\end{aligned}
$$

And,

$$
G=\frac{E}{2(1+\mu)}
$$



FIGURE 1.1 Three-dimensional body.


FIGURE 1.2 Equilibrium of elemental volume.

$$
\begin{aligned}
& \frac{\partial \sigma_{x}}{\partial x}+\frac{\partial \tau_{x y}}{\partial y}+\frac{\partial \tau_{x z}}{\partial z}+f_{x}=0 \\
& \frac{\partial \tau_{x y}}{\partial x}+\frac{\partial \sigma_{y}}{\partial y}+\frac{\partial \tau_{y z}}{\partial z}+f_{y}=0 \\
& \frac{\partial \tau_{x z}}{\partial x}+\frac{\partial \tau_{y z}}{\partial y}+\frac{\partial \sigma_{z}}{\partial z}+f_{z}=0
\end{aligned}
$$

$$
\left[\begin{array}{ccc}
\sigma_{x x} & \tau_{x y} & \tau_{x z} \\
\tau_{y x} & \sigma_{y y} & \tau_{y z} \\
\tau_{z x} & \tau_{z y} & \sigma_{z z}
\end{array}\right]
$$



FIGURE 1.4 Deformed elemental surface.

$$
\boldsymbol{\epsilon}=\left[\epsilon_{x}, \epsilon_{y}, \epsilon_{z}, \gamma_{y z}, \gamma_{x z}, \gamma_{x y}\right]^{\mathrm{T}}
$$

$\boldsymbol{\epsilon}=\left[\frac{\partial u}{\partial x}, \frac{\partial v}{\partial y}, \frac{\partial w}{\partial z}, \frac{\partial v}{\partial z}+\frac{\partial w}{\partial y}, \frac{\partial u}{\partial z}+\frac{\partial w}{\partial x}, \frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}\right]^{\mathrm{T}}$

$$
\begin{aligned}
& \varepsilon_{x}=\frac{\sigma_{x}}{E}-\frac{\sigma_{y}}{E} \mu-\frac{\sigma_{z}}{E} \mu \\
& \varepsilon_{y}=-\frac{\sigma_{x}}{E} \mu+\frac{\sigma_{y}}{E}-\frac{\sigma_{z}}{E} \mu \\
& \varepsilon_{z}=-\frac{\sigma_{x}}{E} \mu-\frac{\sigma_{y}}{E} \mu+\frac{\sigma_{z}}{E} \\
& \gamma_{y z}=\frac{\tau_{y z}}{G} \\
& \gamma_{x z}=\frac{\tau_{x z}}{G} \\
& \gamma_{x y}=\frac{\tau_{x y}}{G}
\end{aligned}
$$



$$
\begin{aligned}
\varepsilon_{x} & =\frac{\sigma_{x}}{E}-\frac{\sigma_{y}}{E} \mu-\frac{\sigma_{z}}{E} \mu \\
\varepsilon_{y} & =-\frac{\sigma_{x}}{E} \mu+\frac{\sigma_{y}}{E}-\frac{\sigma_{z}}{E} \mu \\
\varepsilon_{z} & =-\frac{\sigma_{x}}{E} \mu-\frac{\sigma_{y}}{E} \mu+\frac{\sigma_{z}}{E} \\
\epsilon_{x} & +\epsilon_{y}+\epsilon_{z}=\frac{(1-2 v)}{E}\left(\sigma_{x}+\sigma_{y}+\sigma_{z}\right)
\end{aligned}
$$

$$
\begin{aligned}
& \varepsilon_{x}=\frac{\sigma_{x}}{E}-\frac{\sigma_{y}}{E} \mu-\frac{\sigma_{z}}{E} \mu \\
& \varepsilon_{y}=-\frac{\sigma_{x}}{E} \mu+\frac{\sigma_{y}}{E}-\frac{\sigma_{z}}{E} \mu \\
& \varepsilon_{z}=-\frac{\sigma_{x}}{E} \mu-\frac{\sigma_{y}}{E} \mu+\frac{\sigma_{z}}{E} \\
& \gamma_{y z}=\frac{\tau_{y z}}{G} \\
& \gamma_{x z}=\frac{\tau_{x z}}{G} \\
& \gamma_{x y}=\frac{\tau_{x y}}{G}
\end{aligned}
$$

$$
\left[\begin{array}{l}
\varepsilon_{x x} \\
\varepsilon_{y y} \\
\varepsilon_{z x} \\
\gamma_{y z} \\
\gamma_{z x} \\
\gamma_{x y}
\end{array}\right]=\left[\begin{array}{cccccc}
\frac{1}{E} & -\frac{\nu}{E} & -\frac{\nu}{E} & 0 & 0 & 0 \\
-\frac{\nu}{E} & \frac{1}{E} & -\frac{\nu}{E} & 0 & 0 & 0 \\
-\frac{\nu}{E} & -\frac{\nu}{E} & \frac{1}{E} & 0 & 0 & 0 \\
0 & 0 & 0 & \frac{1}{G} & 0 & 0 \\
0 & 0 & 0 & 0 & \frac{1}{G} & 0 \\
0 & 0 & 0 & 0 & 0 & \frac{1}{G}
\end{array}\right]\left[\begin{array}{c}
\sigma_{x x} \\
\sigma_{y x} \\
\sigma_{z x} \\
\sigma_{y z} \\
\sigma_{z x} \\
\sigma_{x y}
\end{array}\right]
$$

$$
\left[\begin{array}{c}
\varepsilon_{z x} \\
\varepsilon_{y y} \\
\varepsilon_{x y} \\
\gamma_{y z} \\
\gamma_{x x} \\
\gamma_{x y}
\end{array}\right]=\left[\begin{array}{cccccc}
\frac{1}{E} & -\frac{\nu}{E} & -\frac{\nu}{E} & 0 & 0 & 0 \\
-\frac{\nu}{E} & \frac{1}{E} & -\frac{\nu}{E} & 0 & 0 & 0 \\
-\frac{\nu}{E} & -\frac{\nu}{E} & \frac{1}{E} & 0 & 0 & 0 \\
0 & 0 & 0 & \frac{1}{G} & 0 & 0 \\
0 & 0 & 0 & 0 & \frac{1}{G} & 0 \\
0 & 0 & 0 & 0 & 0 & \frac{1}{G}
\end{array}\right]\left[\begin{array}{c}
\sigma_{x x} \\
\sigma_{y x} \\
\sigma_{z z} \\
\sigma_{y z} \\
\sigma_{z x} \\
\sigma_{x y}
\end{array}\right]
$$

$\left[\begin{array}{c}\sigma_{x x} \\ \sigma_{y y} \\ \sigma_{x z} \\ \sigma_{y z} \\ \sigma_{z x} \\ \sigma_{x y}\end{array}\right]=\frac{E}{(1+\boldsymbol{v})(1-2 \boldsymbol{v})}\left[\begin{array}{cccccc}1-\boldsymbol{v} & \boldsymbol{v} & \boldsymbol{v} & 0 & 0 & 0 \\ \boldsymbol{v} & 1-\boldsymbol{v} & \boldsymbol{v} & 0 & 0 & 0 \\ \boldsymbol{v} & \boldsymbol{v} & 1-\boldsymbol{v} & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{(1-2 \boldsymbol{v})}{2} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{(1-2 \boldsymbol{v})}{2} & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{(1-2 \boldsymbol{v})}{2}\end{array}\right] /\left[\begin{array}{c}\varepsilon_{x x} \\ \varepsilon_{y y} \\ \varepsilon_{x x} \\ \gamma_{y z} \\ \gamma_{x x} \\ \gamma_{x y}\end{array}\right]$
D = Constitutive matrix

One Dimension :
In One dimension, Weir havel Norminalishess or
Special cases along $x$ and and the corresponding nounal strain. $e$.

Stress - strain relation, $\sigma=E \in$
$\rightarrow$ Two Dimensions;
In two dimensions, the problems ore modeled as plane stress \& plane strain.

Plane sacs:
A state of plane stress is sard, to dxiots when the elastic body is very thin and/ there are no loads, applied in the coordinate direction parallel to the thickness "In other cords, for some twa l dimensional objects the stresses can be produced only in (two directions and not possible in the third direction.
rvieplane stress andysis includes problems such as Plates with holes, piles do other changes in geometry.


Special cases
' strata in $x$-directions. ask 11 its sot .11

strain in $y$-direction :
hd on $\epsilon_{y}=\theta \frac{\sigma_{x}}{E}+\frac{\sigma_{y}}{E} \quad \epsilon_{4}=-\theta \frac{\sigma_{z}}{E_{1}}-\frac{v \sigma_{y}}{E}+0$
a) We have

$$
\begin{aligned}
& \varepsilon_{x}=\frac{\sigma_{x}}{E}-\frac{\sigma_{y}}{E} \mu-\frac{\sigma_{z}}{E} \mu \\
& \varepsilon_{y}=-\frac{\sigma_{x}}{E} \mu+\frac{\sigma_{y}}{E}-\frac{\sigma_{z}}{E} \mu \\
& \varepsilon_{z}=-\frac{\sigma_{x}}{E} \mu-\frac{\sigma_{y}}{E} \mu+\frac{\sigma_{z}}{E} \\
& \gamma_{y z}=\frac{\tau_{y z}}{G} \\
& \gamma_{x z}=\frac{\tau_{x z}}{G} \\
& \gamma_{x y}=\frac{\tau_{x y}}{G}
\end{aligned}
$$



Af $\lambda_{x, y}=\frac{x y}{G} \frac{E}{E}$ ashes


$$
\begin{aligned}
& \epsilon_{x}=\frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E} \\
& \epsilon_{y}=-v \frac{\sigma_{y}}{E} \lambda+v \frac{\sigma_{y}}{E} \\
& \epsilon_{z}=-v \sigma_{z}-v \frac{\sigma_{y}}{E} \\
& q_{z y}=\frac{2(1+v)}{E} \tau_{x y}
\end{aligned}
$$

Now Coring to the, stress - Strain Relation ton plane
15) 4
lime tor 1



Special cases
formation of Matrix, wy ? il

$$
\begin{aligned}
& \left\{\begin{array}{l}
\epsilon_{x} \\
\epsilon_{y} \\
\partial_{x y}
\end{array}\right]=\left[\begin{array}{ccc}
\frac{1}{E} & 1, \frac{\nu}{E} & 0 \\
\frac{-\nu}{E} & \frac{1}{E} & 0 / \\
0 & 0 & \frac{2(1+\nu)}{E^{\prime}}
\end{array}\right]\left[\begin{array}{l}
\sigma_{x} \\
\sigma_{y} \\
\tau_{x y}
\end{array}\right\} \\
& \left\{\begin{array}{l}
\epsilon_{x} \\
\epsilon_{y} \\
\nabla_{x y}
\end{array}\right\}=\left[\begin{array}{ccc}
\frac{1}{E} & \frac{\nu}{E} & 0 \\
\frac{v_{1}}{E} & \frac{1}{E} & 0 \\
0 & 0 & 0 \\
\frac{\partial(1+v)}{E}
\end{array}\right]\left[\begin{array}{l}
\sigma_{x} \\
\sigma_{y} \\
\tau_{z y}
\end{array}\right] \\
& \left\{\begin{array}{l}
\epsilon_{x} \\
\epsilon_{y} \\
\nu_{x y}
\end{array}\right\}=\frac{1}{2}\left[\begin{array}{ccc}
1 & -v & 0 \\
-v & 1 & 0 \\
0 & 0 & 0 \\
12(1+v)
\end{array}\right]\left[\begin{array}{c}
\sigma_{x} \\
\sigma_{y} \\
\sigma_{x y}
\end{array}\right] \\
& \because\left\{\begin{array}{l}
\sigma_{x} \\
\sigma_{y} \\
\tau_{x y}
\end{array}\right\}=E\left[\begin{array}{ccc}
1 & -v & 0 \\
-v & 1 & 1 \\
0 & 0 & 2(1+v)
\end{array}\right]^{-1}\left[\begin{array}{l}
\sigma_{x} \\
\epsilon_{y} \\
\gamma_{x y}
\end{array}\right]
\end{aligned}
$$

$$
\begin{aligned}
& \text { Let } A=\left[\begin{array}{ccc}
1 & -v & 0 \\
-1 & 1 & 0 \\
0 & 0 & 2(1+v)
\end{array}\right] \\
& |A|=2(1+v)+v(-2 v(1+v))_{s} \mid \\
& =(1+v) \geq\left[\begin{array}{ll}
1 & -a
\end{array}\right]
\end{aligned}
$$

$$
\begin{aligned}
& =2(1+v)\left[\begin{array}{llc}
1 & v & 0 \\
v & 1 & 0 \\
0 & 9 & \frac{(1-v)}{2}
\end{array}\right] \\
& \therefore A^{-1}=\frac{c^{T}}{1 A 1}=\frac{\phi(1 / v)\left[\begin{array}{lll}
1 & v & 0 \\
v & 1 & 0 \\
0)^{2} & 0 & \frac{(i-v)}{2}
\end{array}\right]}{\sqrt{2((1) v)\left(1-v^{2}\right)}}
\end{aligned}
$$

Special cases

$$
\begin{aligned}
& A^{-1}=\left[\begin{array}{ccc}
1 & v & 0 \\
v & 1 & 0 \\
0 & 0 & \frac{-v}{2}
\end{array}\right] \\
& \because\left\{\left\{\begin{array}{l}
\sigma_{x} \\
\sigma_{y} \\
\tau_{x y}
\end{array}\right\}=\frac{E}{1-v^{2}}\left[\begin{array}{ccc}
1 & v & 0 \\
v & 1 & 0 \\
0 & 0 & \frac{-v}{2}
\end{array}\right]\left\{\begin{array}{c}
\epsilon_{x} \\
\epsilon_{y} \\
\lambda_{x y}
\end{array}\right\}\right. \\
& \sigma=D E \\
& D=\text { Stress - strain Relationship. } \\
& \Rightarrow D^{\prime}=\frac{E}{\left(1-v^{2}\right)}\left[\begin{array}{ccc}
10 & v & 0 \\
v, & 1 & 0 \\
0 & 0 & \frac{1-v}{2}
\end{array}\right](v, i)(v, y)
\end{aligned}
$$

Plane Strain:
The tate of plane strain occurs in members that are not free to expand in the direction perpendicular to the plane of applied forces loads. Plane, strain condition, refers, the occurrence of strain in the body in two directions only and in the third direction the strain is negligible \&s equal to
Zero,
Examples, :- Darns subjected to horizontal loading
Pipes sukjected'to versicle loading


$$
\begin{aligned}
& \text { Here } \\
& \left\{\begin{array}{l}
\sigma_{z} \\
\sigma_{y} \\
\tau_{x y}
\end{array}\right\}=\frac{E}{(1+\nu)(v-2 v)}\left[\begin{array}{ccc}
(1-v) & v & 0 \\
v & (1-v) & 0 \\
0 & 0 & \left(\frac{(-2 v}{s / 2}\right)
\end{array}\right] \cdot\left[\begin{array}{c}
\epsilon_{x} \\
\epsilon_{y} \\
Q_{z y}
\end{array}\right\} \\
& \{\sigma\}=[D]\{\epsilon\} \\
& \frac{E}{(1+v)(1-2 v)}\left[\begin{array}{ccc}
1-p & 0 & 0 \\
0 & 1-v & 0 \\
0 & 0 & \frac{1-2 v}{2}
\end{array}\right]=\begin{array}{l}
\text { Congtatative Mature } \\
0, \\
\text { dor Plane Shin }
\end{array}
\end{aligned}
$$

## Plane Strain:

$\leftrightharpoons$ The tate of plane strain occurs in members that are not free to expand in the direction perpendicular to the plane of applied forces loads. Plane, stain condition, refers, the occurrence of strain in the body in two directions only and in the third direction the strain, is negligible \& equal to Zero,

$$
\boldsymbol{\sigma}=\left[\sigma_{x}, \sigma_{y}, \sigma_{z}, \tau_{y z}, \tau_{x z}, \tau_{x y}\right]^{\mathrm{T}}
$$

For 3 D problems we have, $\boldsymbol{\sigma}=\mathbf{D} \boldsymbol{\epsilon}$

$$
\boldsymbol{\epsilon}=\left[\epsilon_{x}, \epsilon_{y}, \epsilon_{z}, \gamma_{y z}, \gamma_{x z}, \gamma_{x y}\right]^{\mathrm{T}}
$$

$$
\left\{\begin{array}{c}
\sigma_{x} \\
\sigma_{y} \\
\sigma_{z} \\
\tau_{y z} \\
\tau_{x z} \\
\tau_{x y}
\end{array}\right\}=\frac{E}{(1+v)(1-2 v)}\left[\begin{array}{cccccc}
1-v & v & v & 0 & 0 & 0 \\
v & 1-v & v & 0 & 0 & 0 \\
v & v & 1-v & 0 & 0 & 0 \\
0 & 0 & 0 & 0.5-v & 0 & 0 \\
0 & 0 & 0 & 0 & 0.5-v & 0 \\
0 & 0 & 0 & 0 & 0 & 0.5-v
\end{array}\right]\left\{\begin{array}{c}
\boldsymbol{\epsilon}_{x} \\
\boldsymbol{\epsilon}_{y} \\
\boldsymbol{\epsilon}_{z} \\
\gamma_{y z} \\
\gamma_{x z} \\
\gamma_{x y}
\end{array}\right\}
$$

$$
\begin{aligned}
& \epsilon_{z}=0 \\
& \text { Plane Strain } \gamma_{y z} \\
&=0 \\
& \gamma_{x z}=0
\end{aligned}
$$



$$
\left\{\begin{array}{c}
\sigma_{x} \\
\sigma_{y} \\
\tau_{x y}
\end{array}\right\}=\frac{E}{(1+v)(1-2 v)}\left[\begin{array}{ccc}
1-v & v & 0 \\
v & 1-v & 0 \\
0 & 0 & 0.5-v
\end{array}\right]\left\{\begin{array}{c}
\epsilon_{x} \\
\epsilon_{y} \\
\gamma_{x y}
\end{array}\right\}
$$

Plane stress:
$\mathbf{D}=\frac{E}{1-\nu^{2}}\left(\begin{array}{llc}1 & v & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \frac{1-v}{2}\end{array}\right)$
Plane strain: $\quad \mathbf{D}=\frac{E}{(1+v)(1-2 v)}\left(\begin{array}{ccc}1-v & v & 0 \\ v & 1-v & 0 \\ 0 & 0 & \frac{1-2 v}{2}\end{array}\right)$


### 1.4. If a displacement field is described by

$$
\begin{aligned}
& u=\left(x^{2}+4 y^{2}-16 x y\right) 10^{-4} \\
& v=\left(y^{2}-5 x+8 y\right) 10^{-4}
\end{aligned}
$$

determine $\epsilon_{x}, \epsilon_{y}, \gamma_{x y}$ at the point $x=1, y=0$.
1.3. In a plane strain problem, we have

$$
\begin{gathered}
\sigma_{x}=30,000 \mathrm{psi}, \sigma_{y}=-15,000 \mathrm{psi} \\
E=30 \times 10^{6} \mathrm{psi}, v=0.3
\end{gathered}
$$

Determine the value of the stress $\sigma_{z}$.
1.6. A displacement field

$$
\begin{aligned}
& u=2+2 x+4 x^{2}+3 x y^{2} \\
& v=x y-8 x^{2}
\end{aligned}
$$

is imposed on the square element shown in Fig. P1.6.


FIGURE P1. 6
(a) Write down the expressions for $\epsilon_{p}, \epsilon_{p}$, and $\gamma_{\lambda y}$
(b) Plot contours of $\epsilon_{x} \epsilon_{y}$, and $\gamma_{x y}$, using, say, MATLAB software.
(c) Find where $\epsilon_{A}$ is maximum within the square.
1.3 Plane strain condition implies that

$$
\varepsilon_{,}=0=-v \frac{\sigma_{x}}{E}-v \frac{\sigma_{y}}{E}+\frac{\sigma_{.}}{E}
$$

which gives

$$
\sigma_{x}=v\left(\sigma_{x}+\sigma_{y}\right)
$$

We have, $\sigma_{x}=20000 \mathrm{psi} \quad \sigma_{y}=-10000 \mathrm{psi} \quad E=30 \times 10^{6} \mathrm{psi} \quad \mathrm{v}=0.3$
On substituting the values,

$$
\sigma_{i}=3000 \mathrm{psi}
$$

1.4 Displacement field

$$
\begin{aligned}
& \begin{array}{l}
u=10^{-1}\left(-x^{2}+2 y^{2}+6 x y\right) \\
v=10^{-4}\left(3 x+6 y-y^{2}\right) \\
\frac{\partial u}{\partial x}=10^{-4}(-2 x+6 y) \quad \frac{\partial u}{\partial y}=10^{-4}(4 y+6 x) \\
\frac{\partial v}{\partial x}=3 \times 10^{-4} \quad \frac{\partial v}{\partial y}=10^{-4}(6+2 y) \\
\varepsilon=\left\{\begin{array}{c}
\frac{\partial u}{\partial x} \\
\frac{\partial v}{\partial y} \\
\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}
\end{array}\right\} \\
\text { at } x=1, y=0 \\
\varepsilon=10^{-4}\left\{\begin{array}{c}
-2 \\
6 \\
9
\end{array}\right\}
\end{array}
\end{aligned}
$$

$$
\begin{aligned}
& u=0.1 y+4 \\
& v=0
\end{aligned}
$$

It is then easy to see that

$$
\begin{aligned}
& \varepsilon_{x}=\frac{\partial u}{\partial x}=0 \\
& \varepsilon_{y}=\frac{\partial v}{\partial y}=0 \\
& \gamma_{x y}=\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}=0.1
\end{aligned}
$$

1.6 The displacement field is given as

$$
\begin{aligned}
& u=1+3 x+4 x^{3}+6 x y^{2} \\
& v=x y-7 x^{2}
\end{aligned}
$$

(a) The strains are then given by

$$
\begin{aligned}
& \varepsilon_{x}=\frac{\partial u}{\partial x}=3+12 x^{2}+6 y^{2} \\
& \varepsilon_{y}=\frac{\partial v}{\partial y}=x \\
& \gamma_{x y}=\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}=12 x y+y-14 x
\end{aligned}
$$

(b) In order to draw the contours of the strain field using MATLAB, we need to create a script file, which may be edited as a text file and save with ". m " extension. The file for plotting $\varepsilon_{\alpha}$ is given below
file "problp5b.m"
$[X, Y]-$ meshgrid $(-1: .1: 1,-1: .1: 1) ;$ $Z=3 .+12,{ }^{+} X,{ }^{\wedge} 2+6, * Y . \wedge 2$; $\{\mathrm{C}, \mathrm{h}\}=\operatorname{contour}(\mathrm{X}, \mathrm{Y}, Z)$; clabel $(0, h)$;

On running the program, the contour map is shown as follows:


Contours of $\varepsilon_{x}$
Contours of $\varepsilon_{y}$ and $\gamma_{x y}$ are obtained by changing $Z$ in the script file. The numbers on the contours show the function values.
(c) The maximum value of $\varepsilon_{\chi}$ is at any of the corners of the square region. The maximum value is 21 .

## Q6. Define Weighted - Residual method?

## Ans:

Generally, the solutions obtained by solving most of the problems of engineering field are approximate solutions and it is difficult to get accurate solutions, as error exists in the solution.

Weighted residual methods are used to reduce these errors in the problems. This method consists of substituting a trial function in the differential equations formulated for the system and residual obtained is equated to zero. Thus, a solution which is very close to the exact solution, is obtained.

## Q9. Why is variational formulation referred to as weak formulation?

## Ans:

[Nov./Dec.-18, (R13), Q2| Model Paper-III, Q1]
Weakening or reducing the governing differential equation of the problem by the process of integration is referred as weak formulation. In the variational formulation, differential equation of the physical problem is rewritten in the form of an equivalent integral. Thereby, upon integration, differential equation gets reduced, i.e., reduction of double differentiation into single differentiation occurs. Therefore, variational formulation referred to as weak formulation. Using variational formulation it is possible to obtain approximate solutions of spatially continuous type with less difficulty and evaluated approximate solutions forms continuous type functions of coordinates of position within the domain.

General expression for weighted residual methods,

$$
\int_{\Omega} R(x) w_{i}(x) d x=0 \quad i=1,2, \ldots, n
$$

Weighting functions associated with different weighted residual techniques are,

| S.No | Weighted Residual <br> Technique | Expression | Weighting Function |
| :---: | :--- | :--- | :--- |
| 1. | Point Collocation Method | $\int_{\Omega} \delta\left(x-x_{i}\right) R\left(x, a_{i}\right) d x=0$ | $w_{i}=\delta\left(x-x_{i}\right)$ |
| 2. Sub-Domain Method | $\int_{\Omega} R\left(x, a_{i}\right) d x=0$ | $w_{i}=1$ |  |
| 3. | Least square Method | $I=\int_{\Omega}\left[R\left(x, a_{i}\right)\right]^{2} d x=$ minimum | $w_{i}=1$ |
| 4. | Galerkin's Method | $\int_{\Omega} y(x) R\left(x, a_{i}\right) d x=0$ | $w_{i}=y(x)$ |

## Q10. Write about the concept of potential energy?

Ans: It states that, the total potential energy within the body becomes stable or minimum, when the displacement equations satisfy the equilibrium equations. These displacement equations that satisfy equations of equilibrium basically fulfills the boundary conditions and are internally compatible.
Q11. Write the potential energy for beam of span $L$ simply supported at both ends, subjected to a concentrated load $\mathbf{P}$ at mid span. Assume El as constant.

Ans: Potential energy ' $\pi$ ' for a beam of span ' $L$ ', simply supported at both ends and subjected to a concentrated load ' $P$ ' at mid span is given by,

$$
\begin{aligned}
& \pi=U-H \\
& \pi=\frac{P^{2} L^{3}}{\pi^{4} E I}
\end{aligned}
$$

Where,
$E I$ - Flexural rigidity of the beam (constant)
$U$ - Strain energy
$H$ - Work potential.

## Q12. Mention the basic steps of Rayleigh-Ritz method.

Ans: The basic steps of Rayleigh-Ritz method are,

1. Assumption of a displacement field
2. Determination of the total potential
3. Solving the system of equations.

Boundary Value Problem : If a governing equation which is formulated over the domain, consists of dependent variables which compulsorily takes and its partial derivative which may probably takes particular values on the boundary of domain, then such equation describes 'boundary value problem'.
Examples of Boundary Value Problem: Analysis of axial deformation of concrete pier, study of steady-state heat flow in a bar, drawn from solid mechanics and heat transfer.
Initial Value Problem : A governing equation, formulated over a domain is said to define a initial value problem, if dependent variables are compulsorily needed and its partial derivatives are probably needed to specified initially, that is at time $t=0$. Generally, initial value problems depends on time.
Example of Initial Value Problem: Analysis of linear motion of simple pendulum drawn from dynamics.
Boundary and Initial Value problem : If the differential equation formulated for a problem, contain dependent variables which are, needed to take specific values on the boundary and required to specified initially, then such problem is said to be both boundary and initial value problem.
Example of Boundary and Initial Value Problem: Unsteadystate heat transfer in a bar drawn from heat transfer.

Eigen Value Problem : A problem is said to be eigen value problem, if an unknown parameter exists in formulated governing equation in addition to unknown dependent variable. In eigen value problem it is needed to determine both unknown parameter and dependent variable. while satisfying differential equation and related boundary conditions.
Example of Eigen Value Problem: Analysis of axial vibrations of a bar drawn from dynamics.

Finite element method was initially developed for the analysis of aircraft structures, but the wider nature of the theory enables it to be applied for variety of boundary value problems in engineering, where the solution has to be obtained in the region or domain of a systern subject to the fulfillment of certain boundary conditions. The application of finite element methods is more in the following categories of boundary value problems.
(i) Steady-state or equilibrium or time independent problems
(ii) Eigen-value problems
(iii) Transient or propagation problems

The following table gives details about the specific application of FEM in different categories.

| S.No. | Field of Application | Type of Boundary Value Problem |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  |  | Equilibrium Problems | Eigenvalue Problems | Propagation Problems |
| 1. | Aircraft structure | In static analysis of a aircraft wings, fins, missile and rocket structures. | Natural frequencies, flutter, and in the stability of rocket, spacecraft and missile structures. | In the study of response of aircraft structures to random loads, dynamic response of aircraft and space craft to periodic loads. |
| 2. | Mechanical design | In stress analysis of pistons, pressure vessels, gears, composite materials and linkages. | Natural frequencies and stability of gears, machine tools and linkages. | Problems of crack and fracture under dynamic loading. |

weighted residual method is given by,

$$
\begin{aligned}
& \int_{\Omega} W_{i} R\left(x, a_{i}\right) d x=0 \quad i=1,2,3, \ldots, n \\
& \text { Where, } \\
& \quad W_{i}-\text { Weighting function } \\
& \quad \Omega \text { - Domain } \\
& \quad a_{i} \text { - Unknown coefficients }
\end{aligned}
$$

weighted residual method is given by,

$$
\int_{\Omega} W_{i} R\left(x, a_{i}\right) d x=0 \quad i=1,2,3, \ldots, n
$$

Weighting functions associated with different weighted residual techniques are,

| S.No | Weighted Residual <br> Technique | Expression | Weighting Function |
| :---: | :--- | :--- | :--- |
| 1. | Point Collocation Method | $\int_{\Omega} \delta\left(x-x_{i}\right) R\left(x, a_{i}\right) d x=0$ | $w_{i}=\delta\left(x-x_{i}\right)$ |
| 2. Sub-Domain Method | $\int_{\Omega} R\left(x, a_{i}\right) d x=0$ | $w_{i}=1$ |  |
| 3. | Least square Method | $I=\int_{\Omega}\left[R\left(x, a_{i}\right)\right]^{2} d x=\operatorname{minimum}$ | $w_{i}=1$ |
| 4. | Galerkin's Method | $\int_{\Omega} y(x) R\left(x, a_{i}\right) d x=0$ | $w_{i}=y(x)$ |

For the differential equation $\frac{d^{2} y}{d^{2} x}+500 x^{2}=0$ for
$0<X<10$ and with boundary conditions $y(0)=0$ and $y(10)=0$, find the solution of this problem using any two weighted residual methods.

For the differential equation $\frac{d^{2} y}{d^{2} x}+500 x^{2}=0$ for $0<X<10$ and with boundary conditions $y(0)=0$ and $y(10)=0$, find the solution of this problem using any two weighted residual methods.

Given that,
Differential equation,

$$
\frac{d^{2} y}{d x^{2}}+500 x^{2}=0 ; 0 \leq x \leq 10
$$

Boundary conditions, $y(0)=0$ and $y(10)=0$
Consider a trial function,

$$
\begin{aligned}
& y=a_{1} x(10-x) \\
& y=10 a_{1} x-a_{1} x^{2}
\end{aligned}
$$

Differentiating with respect to ' $x$ ',

$$
\frac{d y}{d x}=10 a_{1}-2 a_{1} x
$$

Again differentiating with respect to ' $x$ ',

$$
\frac{d^{2} y}{d x^{2}}=-2 a_{1}
$$

Substituting in equation (1),

$$
\text { Residual, } R=-2 a_{1}+500 x^{2}
$$

## 1. Point Collocation Method

In this method residual is set to 0 .

$$
\begin{array}{r}
\text { i.e., } R=0 \\
-2 a_{1}+500 x^{2}=0
\end{array}
$$

One collocation point is required, since there exist one unknown coefficient in the residual. Collocation point should lied between 0 and 10 .

Assume, collocation point, $x=5$

$$
\begin{aligned}
-2 a_{1}+500(5)^{2} & =0 \\
500 \times 25 & =2 a_{1} \\
\therefore a_{1} & =6250 \\
\therefore \text { Trial function, } y & =6250 \times(10-x)
\end{aligned}
$$

2. Sub-domain Collocation Method

This method involves setting, integral of residual over sub-domain to zero.

$$
\begin{aligned}
& \int_{0}^{10} R d x=0 \\
& \int_{0}^{10}\left(-2 a_{1}+500 x^{2}\right) d x=0 \\
& -2 a_{1}(x)_{0}^{10}+500 \times\left(\frac{x^{3}}{3}\right)_{0}^{10}=0 \\
& -2 a_{1}(10)+\frac{500}{3}\left(10^{3}\right)=0 \\
& -20 a_{1}+\frac{500}{3} \times 10^{3}=0 \\
& 20 a_{1}=\frac{500 \times 10^{3}}{3} \\
& a_{1}=\frac{25000}{3}
\end{aligned}
$$

$\therefore$ Trial function, $y=\frac{25000}{3} x(10-x)$

## 3. Least Square Method

In this method, integral of square of weighted residual over the domain is minimum

$$
\begin{aligned}
I & =\int_{0}^{10} R^{2} d x \\
& =\int_{0}^{10}\left(-2 a_{1}+500 x^{2}\right)^{2} d x \\
& =\int_{0}^{10}\left(4 a_{1}^{2}+250000 x^{4}-2000 x^{2} a_{1}\right) d x \\
& =4 a_{1}^{2}(x)_{0}^{10}+250000\left(\frac{x^{5}}{5}\right)_{0}^{10}-2000\left(\frac{x^{3}}{3}\right)_{0}^{10} a_{1} \\
I & =4 a_{1}^{2} \times 10+50000 \times 10^{5}-\frac{2000}{3} \times 10^{3} a_{1}
\end{aligned}
$$

For stationary value of ' $I$ ',

$$
\begin{aligned}
& \frac{\partial I}{\partial a_{1}}=0 \\
& 8 a_{1} \times 10-\frac{2000}{3} \times 10^{3}=0
\end{aligned}
$$

$$
\begin{gathered}
8 a_{1}=\frac{2000}{3} \times 10^{2} \\
a_{1}=\frac{250}{3} \times 10^{2}
\end{gathered}
$$

Trial function, $y=\frac{25000}{3} x(10-x)$

## 4. Galerkin's Method

In Galerkin's method, the domain integral which is the product of trial function and residual is set to zero.

$$
\begin{aligned}
& \int_{0}^{10} y R d x=0 \\
& \int_{0}^{10} y\left(-2 a_{1}+500 x^{2}\right) d x=0 \\
& \int_{0}^{10} a_{1} x(10-x)\left(-2 a_{1}+500 x^{2}\right) d x=0 \\
& \int_{0}^{10}\left(10 a_{1} x-a_{1} x^{2}\right)\left(-2 a_{1}+500 x^{2}\right) d x=0 \\
& \int_{0}^{10}\left(-20 a_{1}^{2} x+5000 a_{1} x^{3}+2 a_{1}^{2} x^{2}-500 a_{1} x^{4}\right) d x=0
\end{aligned}
$$

$$
\begin{aligned}
& \int_{0}^{10}\left(-20 a_{1}^{2} x+5000 a_{1} x^{3}+2 a_{1}^{2} x^{2}-500 a_{1} x^{4}\right) d x=0 \\
& -20 a_{1}^{2}\left(\frac{x^{2}}{2}\right)_{0}^{10}+5000 a_{1}\left(\frac{x^{4}}{4}\right)_{0}^{10}+2 a_{1}^{2}\left(\frac{x^{3}}{3}\right)_{0}^{10}-500 a_{1}\left(\frac{x^{5}}{5}\right)_{0}^{10}=0 \\
& -10 a_{1}^{2} \times 10^{2}+1250 a_{1} \times 10^{4}+\frac{2}{3} a_{1}^{2} \times 10^{3}-100 a_{1} \times 10^{5}=0 \\
& \quad-\frac{1000}{3} a_{1}^{2}+25 \times 10^{5} a_{1}=0 \\
& \therefore \quad a_{1}=7500
\end{aligned}
$$

$\therefore \quad$ Trial function, $y=7500 x(10-x)$.

The following differential equation is available for a physical phenomenon.
$\frac{d^{2} y}{d x^{2}}+50=0 \leq x \leq 10$
The Trial function is $y=a_{1} x(10-x)$
The boundary conditions are: $\begin{gathered}y(0)=0 \\ y(10)=0\end{gathered}$
Find the value of the parameter $a_{1}$ by the following methods.
(i) Least square method
(ii) Galerkin's method.

Consider a bar subjected to a uniform axial load as shown in the figure, which can steadily show that the deformation of a body is given by differential equation $A E \frac{d^{2} u}{d x^{2}}+q_{0}=0$ with boundary conditions $u(0)=0, \frac{d u}{d x}=0$ at $x=I$. Find the approximate solution by using weighted residual method.


Figure

Consider a bar subjected to a uniform axial load as shown in the figure, which can steadily show that the deformation of a body is given by differential equation $A E \frac{d^{2} u}{d x^{2}}+q_{0}=0$ with boundary conditions $u(0)=0, \frac{d u}{d x}=0$ at $x=I$. Find the approximate solution by using weighted residual method.


Figure

## Ans:

Given that,
Differential equation,

$$
A E \frac{d^{2} u}{d x^{2}}+q_{0}=0
$$

Boundary conditions,

$$
u(0)=0, \frac{d u}{d x}(l)=0
$$

Let, the trial function be,

$$
u(x)=a_{0}+a_{1} x+a_{2} x^{2}
$$

Differentiating equation (2),

$$
\frac{d u}{d x}=a_{1}+2 a_{2} x
$$

Again differentiating,

$$
\begin{equation*}
\frac{d^{2} u}{d x^{2}}=2 a_{2} \tag{3}
\end{equation*}
$$

Subjecting to boundary condition $u(0)=0$,

$$
\begin{aligned}
& 0=a_{0}+a_{1}(0)+a_{2}(0)^{2} \\
& a_{0}=0
\end{aligned}
$$

Subjecting to boundary condition $\frac{d u}{d x}(l)=0$

$$
\begin{aligned}
& 0=a_{1}+2 a_{2}(l) \\
& a_{1}=-2 a_{2} l
\end{aligned}
$$

On substituting coefficients in equation (2),

$$
\begin{align*}
& u(x)=0+\left(-2 a_{2}\right) x+a_{2} x^{2} \\
& u(x)=a_{2} x^{2}-2 a_{2} l x \tag{4}
\end{align*}
$$

On substituting equation (3) in equation (1),
Residual,

$$
\begin{aligned}
& R=A E\left(2 a_{2}\right)+q_{0} \\
& R=2 a_{2} A E+q_{0}
\end{aligned}
$$

Using point collocation method (a weighted residual method) in which residual is set to zero.

$$
\begin{aligned}
& R=0 \\
& 2 a_{2} A E+q_{0}=0 \\
& a_{2}=-\frac{q_{0}}{2 A E}
\end{aligned}
$$

On substituting ' $a_{2}$ ' value in equation (4), approximate solution for elongation at any distance ' $x$ ' is obtained.

$$
\begin{aligned}
u(x) & =\left(-\frac{q_{0}}{2 A E}\right) x^{2}-2\left(-\frac{q_{0}}{2 A E}\right) / x \\
\therefore u(x) & =-\frac{q_{0}}{2 A E}\left(x^{2}-2 l x\right)
\end{aligned}
$$

Elongation at free end i.e., at $x=l$,

$$
\begin{aligned}
u_{l} & =-\frac{q_{0}}{2 A E}\left(l^{2}-2 l(l)\right) \\
& =-\frac{q_{0}}{2 A E}\left(-l^{2}\right) \\
\therefore u_{i} & =\frac{q_{0} l^{2}}{2 A E}
\end{aligned}
$$

Discuss the following methods to solve the given differential equation: EI $\frac{d^{2} y}{d x^{2}}-M(x)=0$
With the boundary condition $\mathrm{y}(0)=0$ and $\mathrm{y}(\mathrm{l})=0$

1. Sub-domain method
2. Point collocation method.


## Discuss the following methods to solve the Ans:

given differential equation: $\mathbf{E I} \frac{d^{2} y}{d x^{2}}-\mathbf{M}_{(x)}=0$
With the boundary condition $y(0)=0$ and $y(l)=0$

1. Sub-domain method
2. Point collocation method.

Given that,
Integral equation, $E I \frac{d^{2} y}{d x^{2}}-M(x)=0$
Boundary conditions,

$$
y(0)=0 \text { and } y(x)=0
$$

Let, $y=a \sin \left(\frac{\pi x}{l}\right)$ be the trial function for deflection.
Differentiating ' $y^{\prime}$,

$$
\begin{aligned}
\frac{d y}{d x} & =a \frac{\pi}{l} \cos \left(\frac{\pi x}{l}\right) \\
\frac{d^{2} y}{d x^{2}} & =-a \frac{\pi^{2}}{l^{2}} \sin \left(\frac{\pi x}{l}\right)
\end{aligned}
$$

Therefore,

$$
\begin{align*}
& \text { Residual, } R=E I \frac{d^{2} y}{d x^{2}}-M(x) \\
& R=E I\left[-\frac{\pi^{2} a}{l^{2}} \sin \left(\frac{\pi x}{l}\right)\right]-M(x) \tag{1}
\end{align*}
$$

## 1. Sub-domain Method

In this method, integral of residual over domain is set to zero.

$$
\begin{aligned}
& \int_{0}^{l} R d x=0 \\
& \int_{0}^{l}\left[-\frac{\pi^{2} a}{l^{2}} E I \sin \left(\frac{\pi x}{l}\right)-M(x)\right] d x=0 \\
& -\frac{\pi^{2} a}{l^{2}} E I\left[-\cos \left(\frac{\pi x}{l}\right)\left(\frac{l}{\pi}\right)\right]_{o}^{l}-[M x]_{o}^{l}=0 \\
& \frac{\pi a E I}{l}[\cos (\pi)-\cos (0)]-M l=0 \\
& \frac{\pi a E I}{l}(-1-1)-M l=0 \\
& -\frac{2 \pi a E I}{l}=M l \\
& \quad a=-\frac{M l^{2}}{2 \pi E I}=-0.159 \frac{M l^{2}}{E I}
\end{aligned}
$$

Trial function, $y=-0.159 \frac{M l^{2}}{E I} \sin \left(\frac{\pi x}{l}\right)$

## 1. Sub-domain Method

In this method, integral of residual over domain is set to zero.

$$
\begin{aligned}
& \int_{0}^{l} R d x=0 \\
& \int_{0}^{l}\left[-\frac{\pi^{2} a}{l^{2}} E I \sin \left(\frac{\pi x}{l}\right)-M(x)\right] d x=0 \\
& -\frac{\pi^{2} a}{l^{2}} E I\left[-\cos \left(\frac{\pi x}{l}\right)\left(\frac{l}{\pi}\right)\right]_{o}^{l}-[M x]_{o}^{l}=0 \\
& \frac{\pi a E I}{l}[\cos (\pi)-\cos (0)]-M l=0 \\
& \frac{\pi a E I}{l}(-1-1)-M l=0 \\
& -\frac{2 \pi a E I}{l}=M l \\
& a=-\frac{M l^{2}}{2 \pi E I}=-0.159 \frac{M l^{2}}{E I}
\end{aligned}
$$

Trial function, $y=-0.159 \frac{M l^{2}}{E I} \sin \left(\frac{\pi x}{l}\right)$

Trial function, $y=-0.159 \frac{M l^{2}}{E I} \sin \left(\frac{\pi x}{l}\right)$

$$
\begin{aligned}
& \text { At } x=l / 2, y=y_{\max } \\
& \text { i.e., } y_{\max }=-0.159 \frac{M l^{2}}{E I} \sin \left(\frac{\pi l}{2 l}\right) \\
& \therefore y_{\max }=-0.159 \frac{M l^{2}}{E I} \quad\left(\because \sin \frac{\pi}{2}=1\right)
\end{aligned}
$$

## 2. Point Collocation Method

In this method, the residual is set to zero.

$$
\begin{aligned}
& R=0 \\
& -\frac{\pi^{2} a}{l^{2}} E I \sin \left(\frac{\pi x}{l}\right)-M(x)=0 \\
& -\frac{\pi^{2} a}{l^{2}} E I \sin \left(\frac{\pi x}{l}\right)=M
\end{aligned}
$$

At $x=l / 2, y=y_{\max }$. Therefore, substitute $x=l / 2$ in the above equation.

$$
\begin{aligned}
& -\frac{\pi^{2} a}{l^{2}} E I \sin \left(\frac{\pi l}{2 l}\right)=M \\
& -\frac{\pi^{2} a E I}{l^{2}}=M \\
& \therefore a=-\frac{M l^{2}}{\pi^{2} E I}
\end{aligned}
$$

Trial function, $y=-\frac{M l^{2}}{\pi^{2} E I} \sin \left(\frac{\pi x}{l}\right)$

## 2. Point Collocation Method

In this method, the residual is set to zero.

$$
\begin{aligned}
& R=0 \\
& -\frac{\pi^{2} a}{l^{2}} E I \sin \left(\frac{\pi x}{l}\right)-M(x)=0 \\
& -\frac{\pi^{2} a}{l^{2}} E I \sin \left(\frac{\pi x}{l}\right)=M
\end{aligned}
$$

At $x=l / 2, y=y_{\max }$. Therefore, substitute $x=l / 2$ in the above equation.

$$
\begin{aligned}
& -\frac{\pi^{2} a}{l^{2}} E I \sin \left(\frac{\pi l}{2 l}\right)=M \\
& -\frac{\pi^{2} a E I}{l^{2}}=M \\
& \therefore a=-\frac{M l^{2}}{\pi^{2} E I}
\end{aligned}
$$

Trial function, $y=-\frac{M l^{2}}{\pi^{2} E I} \sin \left(\frac{\pi x}{l}\right)$

Trial function, $y=-\frac{M l^{2}}{\pi^{2} E I} \sin \left(\frac{\pi x}{l}\right)$

$$
\text { At } x=l / 2, y=y_{\max }
$$

$$
y_{\max }=\frac{M l^{2}}{\pi^{2} E I} \sin \left(\frac{\pi l}{2 l}\right)
$$

$$
\therefore \quad y_{\max }=-0.101 \frac{M l^{2}}{E I}
$$

$$
\left(\because \sin \frac{\pi}{2}=1\right)
$$

## Variational Methods

Variational method involves rewriting the differential equation of physical problem in the form of equivalent integral.

Obtained integral is termed as functional and is allowed to become stationary. Functional becomes stationary at extremum conditions i.e., minimum or maximum conditions.

Therefore, functional is allowed to reach extremum conditions by using appropriate trial functions.

For a problem, trial function which is employed to make the integral stationary is termed as approximate solution.

Any approximation method which uses the principle such as principle of minimum potential energy and principle of virtual work are said to be variational principles.

## Rayleigh-Ritz method

is a variational method and is employed to evaluate solution of structural problems. Potential energy will be stored in all structures when acted upon by the load. Stored potential energy is considered as functional. An approximate solution or trial function is assumed so that functional becomes minimum when trial function is substituted in it.

## Usage of variational method is

limited to the problems governed by the differential equations with order greater than one. If there exists first derivative in the differential equation, it is not possible to apply variational method for that problem.

## Principle of Minimum Potential Energy

For conservative systems, of all the kinematically admissible displacement fields, those corresponding to equilibrium extremize the total potential energy. If the extremum condition is a minimum, the equilibrium state is stable.

Kinematically admissible displacements are those that satisfy the single-valued nature of displacements (compatibility) and the boundary conditions.

## Principle of Virtual Work

A body is in equilibrium if the internal virtual work equals the external virtual work for every kinematically admissible displacement field $\langle\boldsymbol{\phi}, \boldsymbol{\epsilon}(\phi)\rangle$.

For every displacement field of kinematically admissible type. A body is said to be in equilibrium, if value of the internal virtual work becomes same as that of external virtual work.

$$
\therefore \partial U=\partial \mathrm{W}
$$

The internal virtual strain energy,

$$
\begin{equation*}
\partial U=\int_{V}(\partial \varepsilon)^{T} \sigma d V \tag{1}
\end{equation*}
$$

Where,

$$
U \text { - Strain energy (internal work) }
$$

$W$ - External work
$\varepsilon$ - Strain vector
$\sigma$ - Stress vector
$V$ - Volume of Element.

Total external virtual workdone,

$$
\begin{equation*}
\partial W=(\partial u)^{T} P_{i}+\int_{V}(\partial u)^{T} b d V+\int_{A}(\partial u)^{T} q d A \tag{4}
\end{equation*}
$$

From minimum potential energy principle,
Total potential energy, $\pi=U-W_{p}$
Where,

$$
\begin{aligned}
& U-\text { Strain energy of the system } \\
& W_{p}-\text { Total work potential of the load. }
\end{aligned}
$$

Total potential energy for the general elastic body is given by,

$$
\begin{equation*}
\pi=\frac{1}{2} \int_{V} \sigma^{T} \varepsilon d V-\int_{V} u^{T} b d V-\int_{A} u^{T} q d A-\sum_{i} u_{i}^{T} P_{i} \tag{6}
\end{equation*}
$$

i.e., $\frac{\partial \pi}{\partial x}=0$

Where,
$\pi$ - Total potential energy
$x$ - Displacement field
Also, $\pi=U-W$
Where,
$U$ - Strain energy within the body
$W$ - Work done on the system

## Derivation of Equilibrium Equation

For the body shown in figure,
Let,
$P$ - Load acting on the body
$u$ - Deformation
$K$ - Stiffness


Figure
Workdone on the body,

$$
\begin{aligned}
W & =\text { Applied load } \times \text { Deformation length } \\
& =P \times u
\end{aligned}
$$

Strain energy within the body,

$$
\begin{aligned}
U & =\frac{1}{2}[\text { Force in the body } \times \text { deformation length }] \\
& =\frac{1}{2}[K u \times u] \\
& =\frac{1}{2} K u^{2}
\end{aligned}
$$

But, total potential energy, $\pi=U-W$

$$
=\frac{1}{2} K u^{2}-P \cdot u
$$

Under equilibrium condition, for potential energy to be minimum,

$$
\begin{aligned}
\frac{\partial \pi}{\partial u} & =0 \\
\frac{\partial}{\partial u}\left[\frac{1}{2} K u^{2}-P . u\right] & =0 \\
K u-P & =0 \\
\therefore K u & =P
\end{aligned}
$$

$\therefore \quad$ Equilibrium equation is given by, $K u=P$.

## Finite Element Discretization Corresponding to Variational

 FormulationThe mathematical model of a bar is discretized and assembled to form a model, which comprises of small bar elements. Then, equations for finite elements of the bar is derived by using total potential energy functional.


Figure (2): Bar Member


Figure (3): Discretized Bar Member


Figure (2): Bar Member


Figure (3): Discretized Bar Member

TPE-total potential energy
Consider a bar member, divided into certain number of elements, as shown in figure. TPE functionals are scalar quantities and for a discretized bar member, TPE functional is the summation of functional of each element.
i.e., $\pi=\pi_{1}+\pi_{2}+\ldots .+\pi_{n-1}+\pi_{n}$

Where,

$$
n \text { - Total number of elements }
$$

$$
\text { i.e., } \pi=\pi_{1}+\pi_{2}+\ldots .+\pi_{n-1}+\pi_{n}
$$

Similarly, internal energy, external energy and condition of minimum potential energy principle are formed by summation of the corresponding parameter of each finite element.

Minimum potential energy principle is given by,

$$
\delta \pi=\delta \pi_{1}+\delta \pi_{2}+\ldots .+\delta \pi_{n-1}+\delta \pi_{n}=0
$$

For an element ' $e$ ' as a whole, based on variational calculus fundamentals, the above equation can be written as

$$
\delta \pi_{e}=\delta U_{e}-\delta W_{e}=0
$$

This equation is called variational equation and it is a basic formulation, from which stiffness equations for the elements can be developed, after discretization of displacement field for the bar member.

Rayleigh - Ritz method is a variational method and mostly used to solve structural problems which are complex in nature.

There are two techniques of problem solving in Rayleigh-Ritz method.

1. Rayleigh-Ritz Method Using Minimum Potential Energy Concept
2. Rayleigh-Ritz Method Using Integral Approach

## 1. Rayleigh-Ritz Method Using Minimum Potential Energy Concept

In this method, initially a trial function in terms of Ritz parameters (coefficients) is considered. Structure of assumed trial function may be a polynomial function or trigonometric function.

Structure of polynomial trial function,

$$
\begin{equation*}
y=a_{1}+a_{2} x+a_{3} x^{2}+a_{4} x^{3}+\ldots \tag{1}
\end{equation*}
$$

Structure of trigonometric trial function,

$$
\begin{equation*}
y=a_{1} \sin \frac{\pi x}{l}+a_{2} \sin \frac{3 \pi x}{l}+a_{3} \sin \frac{5 \pi x}{l}+\ldots \tag{2}
\end{equation*}
$$

Where,

$$
a_{1}, a_{2}, a_{3} \ldots a_{n}-\text { Ritz parameters or coefficients }
$$

Then, total potential energy is formulated using anyone of two trial function structures mentioned above.

Total potential energy,

$$
\pi=U-W
$$

Where, ' $U$ ' and ' $W$ ' are strain energy and workdone due to external force respectively and are specified so that they forms the functions of approximated trial function. Then formulation of total potential energy is carried out by deriving the trial function compatible with the formats of ' $U$ ' and ' $W$ '. Finally, approximate solution can be evaluated when total potential energy is made to reach minimum value.

For potential energy to be minimum,

$$
\frac{\partial \pi}{\partial a_{1}}=\frac{\partial \pi}{\partial a_{2}}=\ldots .=\frac{\partial \pi}{\partial a_{n}}=0
$$

Thus, from above equation, Ritz coefficients $a_{1}, a_{2} \ldots a_{\mathrm{n}}$ can be calculated.

## 2. Rayleigh-Ritz Method Using Integral Approach

This method involves rewriting the differential equation of physical problem in the form of equivalent integral. Obtained integral is termed as functional and is allowed to become stationary. Functional becomes stationary at extremum conditions i.e, minimum or maximum conditions. Therefore, functional is allowed to reach extremum conditions by using appropriate trial functions. For a problem, trial function which is employed to make the integral stationary is termed as approximate solution.

Consider a physical problem whose governing differential equation is given by,

$$
K \frac{d^{2} y}{d x^{2}}+L=0 \quad \begin{aligned}
& \text { Where, } \\
& K, L-\text { Constants or variables }
\end{aligned}
$$

An integral equivalent to differential equation is given by,

$$
I=\int_{0}^{1}\left[\frac{1}{2} K\left(\frac{d y}{d x}\right)^{2}-L y\right] d x
$$

$$
I \text { - Functional }
$$

Then anyone of the types of trial functions in equations (1) and (2) is considered and differentiated so that it is suitable for equivalent integral format. Evaluation of the approximate solution is obtained by making the integral to become stationary.

$$
\frac{\partial I}{\partial a_{i}}=0
$$

Where, $i=1,2,3 \ldots n$,

Subjected to boundary conditions,

$$
\begin{aligned}
& y(0)=y_{0} \\
& y(l)=y_{l}
\end{aligned}
$$

## General Steps:

1. Formulate Potential Energy Functional
2. Assume a trial displacement function, which should satisfy boundary condition
3. Substitute admissible trial displacement function into Potential Energy Functional and simplify it
4. Minimize the Potential Energy Functional so as to obtain the equilibrium condition
5. Determine the unknown displacement, hence strain and stress

Explain the potential energy formulation for obtaining element equations in Finite element methods.

Consider a stepped bar with three nodes and two elements.

$F_{1}, F_{2}, F_{3}$ - Concentrated loads at each node
$u_{1}, u_{2}, u_{3}$ - Displacements at each node
Total potential energy $=$ Strain energy + Workdone

$$
\pi=U-W
$$

("-ve' sign due to workdone on the system)
For minimum potential energy,

$$
\frac{\partial \pi}{\partial s}=0
$$

Since, the bar is divided into two elements,
Total potential energy, $\pi=\pi_{1}+\pi_{2}$
Considering element (1),
Element potential energy,

$$
\begin{aligned}
\pi_{1} & =U_{1}-W_{1} \\
& =\left[\frac{1}{2} k_{1}\left(u_{2}-u_{1}\right)^{2}-\left(F_{1} u_{1}+F_{2} u_{2}\right)\right] \\
\pi_{1} & =\frac{1}{2} k_{1}\left(u_{2}-u_{1}\right)^{2}-F_{1} u_{1}-F_{2} u_{2}
\end{aligned}
$$

For potential energy to be minimum at each node of element (1),

$$
\begin{align*}
& \text { i.e., at node 1, } \\
& \qquad \begin{array}{l}
\frac{\partial \pi_{1}}{\partial u_{1}}=0 \\
\frac{\partial}{\partial u_{1}}\left[\frac{1}{2} k_{1}\left(u_{2}-u_{1}\right)^{2}-F_{1} u_{1}-F_{2} u_{2}\right]=0 \\
k_{1}\left(u_{2}-u_{1}\right)(-1)-F_{1}=0 \\
k_{1} u_{1}-k_{1} u_{2}=F_{1}
\end{array}
\end{align*}
$$

And, at node 2,

$$
\begin{align*}
& \frac{\partial \pi_{1}}{\partial u_{2}}=0 \\
& \frac{\partial}{\partial u_{2}}\left[\frac{1}{2} k_{1}\left(u_{2}-u_{1}\right)^{2}-F_{1} u_{1}-F_{2} u_{2}\right]=0 \\
& k_{1}\left(u_{2}-u_{1}\right)(1)-F_{2}=0 \\
& -k_{1} u_{1}+k_{1} u_{2}=F_{2} \tag{3}
\end{align*}
$$

$$
\begin{gathered}
k_{1} \text { - Element stiffness } \\
=\frac{A_{1} E_{\perp}}{l_{1}}
\end{gathered}
$$

Finite element matrix is obtained by writing equations
(2) and (3) in matrix form.

$$
\begin{gather*}
1  \tag{4}\\
2 \\
{\left[\begin{array}{cc}
k_{1} & -k_{1} \\
-k_{1} & k_{1}
\end{array}\right] 1\left\{\begin{array}{l}
u_{1} \\
2 \\
u_{2}
\end{array}\right\}}
\end{gather*}=\left\{\begin{array}{l}
F_{1} \\
F_{2}
\end{array}\right\}
$$

## Considering element (2),

Element potential energy,

$$
\begin{aligned}
\pi_{2} & =U_{2}-W_{2} \\
& =\left[\frac{1}{2} k_{2}\left(u_{3}-u_{2}\right)^{2}-\left(F_{2} u_{2}+F_{3} u_{3}\right)\right] \\
\pi_{2} & =\frac{1}{2} k_{2}\left(u_{3}-u_{2}\right)^{2}-F_{2} u_{2}-F_{3} u_{3}
\end{aligned}
$$

For potential energy to be minimum at each node of element (2),

$$
\begin{align*}
& \text { i.e., at node 2, } \\
& \frac{\partial \pi_{2}}{\partial u_{2}}=0 \\
& \frac{\partial}{\partial u_{2}}\left[\frac{1}{2} k_{2}\left(u_{3}-u_{2}\right)^{2}-F_{2} u_{2}-F_{3} u_{3}\right]=0 \\
& k_{2}\left(u_{3}-u_{2}\right)(-1)-F_{2}=0 \\
& k_{2} u_{2}-k_{2} u_{3}=F_{2} \tag{5}
\end{align*}
$$

And, at node 3,

$$
\begin{align*}
& \frac{\partial \pi_{2}}{\partial u_{3}}=0 \\
& \frac{\partial}{\partial u_{3}}\left[\frac{1}{2} k_{2}\left(u_{3}-u_{2}\right)^{2}-F_{2} u_{2}-F_{3} u_{3}\right]=0 \\
& k_{2}\left(u_{3}-u_{2}\right)(1)-F_{3}=0 \\
& -k_{2} u_{2}+k_{2} u_{3}=F_{3} \tag{6}
\end{align*}
$$

Finite element matrix is obtained by writing equations
(5) and (6) in matrix form.

$$
\left[\begin{array}{cc}
2 & 3 \\
k_{2} & -k_{2}  \tag{7}\\
-k_{2} & k_{2}
\end{array}\right] 2\left\{\begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{l}
F_{2} \\
F_{3}
\end{array}\right\}
$$

Global finite element matrix is given by,

$$
[K]\{u\}=\{F\}
$$

It obtained by adding equations (4) and (7).

$$
\begin{align*}
& {\left[\begin{array}{cc}
1 & 2 \\
k_{1} & -k_{1} \\
-k_{1} & k_{1}
\end{array}\right] 1\left[\begin{array}{l}
u_{1} \\
u_{1} \\
u_{2}
\end{array}\right\}=\left\{\begin{array}{l}
F_{1} \\
F_{2}
\end{array}\right\}} \\
& {\left[\begin{array}{cc}
k_{2} & -k_{2} \\
-k_{2} & k_{2}
\end{array}\right] 2\left\{\begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{l}
F_{2} \\
F_{3}
\end{array}\right\}}
\end{align*}
$$


$\left[\begin{array}{cc}\left(k_{2}\right) & k_{2} \\ \hdashline-k_{2} & k_{2}\end{array}\right]$

Global finite element matrix is given by,

$$
[K]\{u\}=\{F\}
$$

It obtained by adding equations (4) and (7).

$$
\left.\begin{array}{l}
\text { ceflon } 20 c c \\
{\left[\begin{array}{ccc}
k_{1} & -k_{1} & 0 \\
-k_{1} & k_{1}+k_{2} & -k_{2} \\
0 & -k_{2} & k_{2}
\end{array}\right]} \\
\hline
\end{array}\right]\left[\begin{array}{l}
u_{1} \\
3
\end{array}\right\}=\left\{\begin{array}{c}
F_{1} \\
F_{2} \\
F_{3} \\
u_{3}
\end{array}\right\}
$$

Where,

$$
\begin{aligned}
& {[K] \text { - Global stiffness matrix }} \\
& \{u\} \text { - Global displacement vector } \\
& \{F\} \text { - Load vector }
\end{aligned}
$$

Q. Consider a bar subjected to a uniform axial load as shown in the figure, which can steadily show that the deformation of a body is given by differential equation $A E \frac{d^{2} u}{d x^{2}}+q_{0}=0$ with boundary conditions $u(0)=0, \frac{d u}{d x}=0$ at $x=I$. Find the approximate solution by using weighted residual method.


Figure
Q. The functional form of a bar clamped at one end and left free at the other end and subjected to uniform axial load $q$ is given by,

$$
\mathrm{I}=\int_{0}^{1}\left[\frac{1}{2} \mathrm{AE}\left(\frac{\mathrm{du}}{\mathrm{dx}}\right)^{2}-\mathrm{qu}\right] \mathrm{dx}
$$

The essential boundary is $u(0)=0$, obtain the approximate solution to the problem by using RayleighRitz method.
(Or)

Example 3.1. A bar under uniform load. Consider a bar clamped at one end and left free at the other end and subjected to a uniform axial load $q_{0}$ as shown in Figure 3.4. The governing differential equation is given by


$$
A E \frac{d^{2} u}{d x^{2}}+q=0
$$

with the boundary conditions $u(0)=0 ;\left.\quad \frac{d u}{d x}\right|_{x=L}=0$.
illustrate the solution using the $\mathrm{R}-\mathrm{R}$ method.
Fig. 3.4 Rod under axial load (Example 3.1).
2. Assume a trial displacement function, which should satisfy boundary condition
3. Substitute admissible trial displacement function into Potential Energy Functional and simplify it
4. Minimize the Potential Energy Functional so as to obtain the equilibrium condition
5. Determine the unknown displacement, hence strain and stress

Strain energy stored in the bar, $U=\int_{0}^{L}\left[\frac{1}{2} A E\left(\frac{d u}{d x}\right)^{2}\right] d x$
Work potential of the external forces, $\mathrm{W}=-\int_{0}^{L} q_{0} u d x$.
$I$ (Functional) or $\Pi_{P}=\int_{0}^{L}\left[\frac{1}{2} A E\left(\frac{d u}{d x}\right)^{2}-q_{0} u\right] d x$

$$
u(x) \approx c_{1} x+c_{2} x^{2}
$$

This satisfies the essential boundary condition that $u(0)=0$. We have

$$
\begin{gathered}
\frac{d u}{d x}=c_{1}+2 c_{2} x \\
\Pi_{p}=\int_{0}^{L}\left[\frac{A E}{2}\left(c_{1}+2 c_{2} x\right)^{2}-q_{0}\left(c_{1} x+c_{2} x^{2}\right)\right] d x \\
=\frac{A E}{2}\left[c_{1}^{2} L+\frac{4 c_{2}^{2} L^{3}}{3}+2 c_{1} c_{2} L^{2}\right]-q_{0} \frac{c_{1} L^{2}}{2}-q_{0} c_{2} \frac{L^{3}}{3} \\
\frac{\partial \Pi_{p}}{\partial c_{i}}=0, \quad i=1,2
\end{gathered}
$$

Therefore,

$$
\begin{aligned}
& \frac{\partial \Pi_{p}}{\partial c_{i}}=0 \Rightarrow \frac{A E}{2}\left(2 c_{1} L+2 c_{2} L^{2}\right)-\frac{q_{0} L^{2}}{2}=0 \\
& \frac{\partial \Pi_{p}}{\partial c_{2}}=0 \Rightarrow \frac{A E}{2}\left(8 c_{2} L^{3} / 3+2 c_{1} L^{2}\right)-\frac{q_{0} L^{3}}{3}=0
\end{aligned}
$$

Solving, we obtain

$$
c_{1}=\frac{q_{0} L}{A E}, \quad c_{2}=-\frac{q_{0}}{2 A E}
$$

Thus,

$$
u(x)=\frac{q_{0}}{A E} x(L-x / 2)=\frac{q_{0}}{2 A E}\left(2 L x-x^{2}\right)
$$

Bars

$$
\begin{aligned}
S E & =\frac{1}{2} p \delta \quad\left(\delta=\frac{P l}{A E}\right) \\
& =\frac{1}{2} p \cdot \frac{p l}{A E} \\
& =\frac{1}{2} \frac{p^{2} l}{A E} \cdot x \frac{A}{A} \\
& =\frac{1}{2}\left(\frac{P}{A}\right)^{2} \frac{l A}{E} \\
& =\frac{1}{2} \sigma^{2} \cdot \frac{V}{E}\left(\because E=\frac{\sigma}{E}\right) \\
& =\frac{1}{2} \sigma^{2} \cdot \frac{V}{(\sigma / E)} \\
& =\frac{1}{2} \sigma \cdot E \cdot V .
\end{aligned}
$$

$$
\begin{aligned}
& d(S E)=\frac{1}{2} \sigma \epsilon d v \\
& d v=A \cdot d x \\
& d(S E)=\frac{1}{2} \sigma \cdot E \cdot A \cdot d x \\
&=\frac{1}{2} \cdot E \epsilon \cdot \epsilon \cdot A d x \\
&=\frac{A E}{2} \epsilon^{2} d x \\
&=\frac{A E}{2}\left(\frac{d u}{d x}\right)^{2} d x \\
& S E=\int d(S E)=\int \frac{A E}{2}\left(\frac{d y}{d x}\right)^{2} d x
\end{aligned}
$$

$$
\begin{aligned}
& \text { Beams } \\
& S E=\frac{1}{2} \sigma \cdot E \cdot V=\frac{\sigma^{2} V}{2 E} \\
& \left(\frac{M}{I}=\frac{\sigma}{y}=\frac{E}{R}\right) \\
& \sigma=\frac{M y}{I} \cdot \\
& S E=\frac{M^{2} y^{2}}{I^{2}} \cdot \frac{V}{2 E} \\
& d x \\
& d(S E)=\frac{M^{2} y^{2}}{2 E I^{2}} A \cdot d x
\end{aligned}
$$

$$
\begin{aligned}
& S E=\int d(S E) \\
&=\int \frac{M^{2} I}{2 E I^{2}} d x \\
&\left(\because \int y^{2} d A=I\right) \\
& S E=\int \frac{m^{2}}{2 E I} d x \\
& E I \frac{d^{2} y}{d x^{2}}=M \Rightarrow \frac{M}{E I}=\frac{d^{2} y}{d x^{2}} \\
& S E=\frac{E I}{2} \int\left(\frac{m}{2}\right)^{2} d x \\
& S E=\frac{E I}{2} \int\left(\frac{d^{2} y}{d x^{2}}\right)^{2} d x
\end{aligned}
$$

Example 3.2. A simply supported beam under uniform load. Consider a simply supported beam under uniformly distributed load $q_{0}$ as shown in Figure 3.5. For a deformation $v(x)$, we have



Fig. 3.5 Simply supported beam under load (Example 3.2).

The strain energy

$$
U=\int_{0}^{L} \frac{1}{2} E I\left(\frac{d^{2} v}{d x^{2}}\right)^{2} d x
$$

The potential of the external forces is

$$
V=-\int_{0}^{L} q_{0} v d x
$$

Thus we have the total potential

$$
\Pi_{p}=\int_{0}^{L}\left[\frac{E I}{2}\left(\frac{d^{2} v}{d x^{2}}\right)^{2}-q_{0} v\right] d x
$$

Assume a displacement field. Let us assume $v(x) \approx c_{1} \sin (\pi x / L)$. This satisfi boundary conditions $v(0)=0=v(L)$. We have

$$
\frac{d^{2} v}{d x^{2}}=-c_{1}\left(\frac{\pi}{L}\right)^{2} \sin \frac{\pi x}{L}
$$

Evaluation of the total potential. The total potential of the system is

$$
\begin{array}{rl}
\Pi_{p} & =\int_{0}^{L}\left[\frac{E I}{2}\left(-c_{1}\left(\frac{\pi}{L}\right)^{2} \sin \frac{\pi x}{L}\right)^{2} d x-q_{0} c_{1} \sin \frac{\pi x}{L}\right] d x \\
& =\frac{\pi^{4} E I}{4 L^{3}} c_{1}^{2}-\frac{2 q_{0} L}{\pi} c_{1} \\
\frac{\partial \Pi_{p}}{\partial c_{1}}=0 & \text { Thus the final solution is } \\
\frac{\pi^{4} E I}{2 L^{3}} c_{1}-\frac{2 q_{0} L}{\pi}=0 & v(x)=0.01307 \frac{q_{0} L^{4}}{E I} \sin \frac{\pi x}{L} \\
& \\
c_{1} & =0.01307 \frac{q_{0} L^{4}}{E I}
\end{array}
$$

## Example 1.2

The potential energy for the linear elastic one-dimensional rod (Fig. E1.2), with body force neglected, is

$$
\Pi=\frac{1}{2} \int_{0}^{L} E A\left(\frac{d u}{d x}\right)^{2} d x-2 u_{1}
$$

$$
\text { where } u_{1}=u(x=1)
$$

$u=0$ at $x=0$ and $u=0$ at $x=2$.
Write the expression for the displacement and stress?


Let us consider a polynomial function

$$
u=a_{1}+a_{2} x+a_{3} x^{2}
$$

This must satisfy $u=0$ at $x=0$ and $u=0$ at $x=2$. Thus,

$$
\begin{aligned}
& 0=a_{1} \\
& 0=a_{1}+2 a_{2}+4 a_{3}
\end{aligned}
$$

Hence,

$$
\begin{aligned}
& a_{2}=-2 a_{3} \\
& u=a_{3}\left(-2 x+x^{2}\right) \quad \text { at } \mathbf{x}=\mathbf{1}, \mathbf{u}=\mathbf{u}_{\mathbf{1}} \\
& u_{1}=-a_{3}
\end{aligned}
$$

Then $d u / d x=2 a_{3}(-1+x)$ and

$$
\begin{aligned}
\Pi & =\frac{1}{2} \int_{0}^{2} 4 a_{3}^{2}(-1+x)^{2} d x-2\left(-a_{3}\right) \\
& =2 a_{3}^{2} \int_{0}^{2}\left(1-2 x+x^{2}\right) d x+2 a_{3} \\
& =2 a_{3}^{2}\left(\frac{2}{3}\right)+2 a_{3}
\end{aligned}
$$

We set $\partial \Pi / \partial a_{3}=4 a_{3}\left(\frac{2}{3}\right)+2=0$, resulting in

$$
a_{3}=-0.75 \quad u_{1}=-a_{3}=0.75
$$

The stress in the bar is given by

$$
\sigma=E \frac{d u}{d x}=1.5(1-x)
$$



## Approximative Methods

## Variational Methods

approximation is based on the minimization of a functional, as those defined in the earlier slides.

- Rayleigh-Ritz Method


## Weighted Residual Methods

start with an estimate of the the solution and demand that its weighted average error is minimized

- The Galerkin Method
- The Least Square Method
- The Collocation Method
- The Subdomain Method


## 1 Dimensional Problems

Ex: Bars
Trusses Beams

## Steps in an FE Analysis

## Preprocessor/Modeling



## Preprocessor / Modeling:

- Identification of the appropriateness of analysis by FEM
- Identification of type of analysis
- Idealization, ie., choice of element type/types


## Preprocessor / Modeling:

- Identification of the appropriateness of analysis by FEM
- Identification of type of analysis
- Idealization, ie., choice of element type/types
- Creation of material behavior model
- Discretization of the solution region (meshing)
- Application of boundary conditions


## 

- Formulation of element stiffness matrices
- Assembly of global stiffness matrix [K]
- formulation of load vector $\{F\}$
- Solution of $[K]\{u\}=\{F\}$ to get nodal displacements $\{u\}$


## 

- Formulation of element stiffness matrices
- Assembly of global stiffness matrix [K]
- formulation of load vector $\{F\}$
- Solution of $[K]\{u\}=\{F\}$ to get nodal displacements $\{u\}$
- (Solution of $[K]-\lambda[M]$ in case of dynamic analysis)
- (Solution of $[K]-\lambda\left[K_{g}\right]$ in case of buckling analysis)
- Calculation of elemental stresses


## Postprocessing / View results:

- View results (displacements, stresses, mode shapes, etc.)
- Interpret and validate results
- If required, re-formulate, and re-analysis


## Basic Steps Involved In FEM:

1. Domain Discretization
2. Selection of displacement functions (interpolation)
3. Formation of elemental (stiffness matrix and load vector)
4. Formation of Global (stiffness matrix and load vector) : $\mathrm{K} U=F$
5. Application of boundary condition
6. Solution of simultaneous equations (for unknown nodal displacements )
7. Calculation of stresses and strains
8. Interpretation of results
9. Domain Discretization: It is performed by using the mesh generating programs (preprocessors). This step involves splitting the structure into number of small regular shaped elements. Generally, a body is discretized by using tetrahedron or hexahedron elements in 3D analysis, whereas, by employing triangular or quadrilateral elements in 2D analysis.

## 2. Selection of displacement functions

(Specifying the interpolation function order i.e, Linear or Quadratic approximation)

1D Elements Line Element

2D Elements Surface Element
1D Elements
Line Element
2D Elements
Surface Element
3D Elements
Element Name
Spring, Damper
Beam, Truss
Shell, Plane2D
Hexder Second Order
Tetrahedral

The loading consists of three types: the body force $f$, the traction force $T$, and the point load $P_{i}$. These forces are shown acting on a body in Fig. 3.1. A body force is a distributed



One-dimensional bar loaded by traction, body, and point loads.

(a)

(b)

FIGURE 3.2 Finite element modeling of a bar.
element model in Fig. 3.2b has five dof. The displacements along each dof are denoted by $Q_{1}, Q_{2}, \ldots, Q_{5}$. In fact, the column vector $\mathbf{Q}=\left[Q_{1}, Q_{2}, \ldots, Q_{5}\right]^{\mathrm{T}}$ is called the global displacement vector. The global load vector is denoted by $\mathbf{F}=\left[F_{1}, F_{2}, \ldots, F_{5}\right]^{\mathrm{T}}$. The vectors $\mathbf{Q}$ and $\mathbf{F}$ are shown in Fig. 3.3. The sign convention used is that a displacement or


FIGURE $3.3 \quad \mathbf{a}$ and $F$ vectors


FIGURE 3.4 Element connectivity.
3. Formation of elemental (stiffness matrix and load vector)

4. Formation of Global (stiffness matrix and load vector)

$$
(\mathbf{K Q}-\mathbf{F})=0
$$

## 5. Application of boundary condition

Apply displacement and forced boundary conditions
6. Solution of simultaneous equations

$$
(\mathbf{K} \mathbf{Q}-\mathbf{F})=0
$$

Interpolation Functions: Interpolation function is also known as approximate function, which is defined to obtain the approximate solution for a given problem, by dividing the domain into smaller elements. It must be selected for each element in such a way that, it must provide better solution for a finite element problem.

## Forms of Interpolation Functions

There are two forms of Interpolation functions. They are,

1. Polynomial form
2. Trigonometric form.

Considering one dimensional element in which ' $\phi$ ' represents a field variable and,
The variation of ' $\phi$ ' in polynomial form is represented as,

$$
\phi=a_{1}+a_{2} x+a_{3} x^{2}+a_{4} x^{3}+a_{5} x^{4}+\ldots
$$

The variation of ' $\phi$ ' in trigonometric form is represented as,

$$
\phi=a_{1} \sin \left(\frac{\pi x}{l}\right)+a_{2} \sin \left(\frac{3 \pi x}{l}\right)+a_{3} \sin \left(\frac{5 \pi x}{l}\right)+\ldots
$$

Then, consider two equations of ' $\phi$ ',

$$
\begin{aligned}
\quad \phi & =a_{1}+a_{2} x+a_{3} x^{2}+a_{4} x^{3}+a_{5} x^{4} \\
\text { And, } \quad \phi & =a_{1}+a_{2} x+a_{3} x^{2}
\end{aligned}
$$

Equations (1) gives the accurate solution because of its higher order, For most of the problems, polynomial form is adopted because of the following advantages,

1. Easy to formulate the equation
2. More accurate results are obtained
3. Simple Structure.


Figure (a): Approximation


Figure (b): Linear Approximation


Figure (c): Quadratic Approximation

## Derivation of 1D linear interpolation function for the displacement function

## Or 2 noded bar element Or Linear bar element



Figure

Function, $u=u(x)$
Consider a linear interpolation formula for a function $u=u(x)$ in the range $u_{1}$ and $u_{2}$ as,

$$
\begin{equation*}
u=a_{1}+a_{2} x \tag{1}
\end{equation*}
$$

Where,

$$
a_{1}, a_{2}-\text { Constants }
$$

Applying boundary conditions,

$$
\text { i.e., } u\left(x_{1}\right)=u_{1} \text { and } u\left(x_{2}\right)=u_{2} \text {. }
$$

Substituting the above values in equation (1),

$$
\text { i.e., } \begin{align*}
u_{1} & =a_{1}+a_{2} x_{1}  \tag{2}\\
u_{2} & =a_{1}+a_{2} x_{2} \tag{3}
\end{align*}
$$

Solving the above two equations,

$$
\begin{aligned}
& a_{1}=\frac{u_{1} x_{2}-u_{2} x_{1}}{x_{2}-x_{1}} \text { and } \\
& a_{2}=\frac{u_{2}-u_{1}}{x_{2}-x_{1}}
\end{aligned}
$$

Substituting ' $a_{1}$ ' and ' $a_{2}$ ' values in equation (1),

$$
\begin{aligned}
u & =\frac{\left(u_{1} x_{2}-u_{2} x_{1}\right)}{\left(x_{2}-x_{1}\right)}+x \frac{u_{2}-u_{1}}{\left(x_{2}-x_{1}\right)} \\
& =\frac{u_{1} x_{2}-u_{2} x_{1}+u_{2} x-u_{1} x}{\left(x_{2}-x_{1}\right)} \\
& =\frac{u_{1}\left(x_{2}-x\right)+u_{2}\left(x-x_{1}\right)}{\left(x_{2}-x_{1}\right)}
\end{aligned}
$$

$$
\therefore \quad u=u_{1} \frac{\left(x_{2}-x\right)}{\left(x_{2}-x_{1}\right)}+u_{2} \frac{\left(x-x_{1}\right)}{\left(x_{2}-x_{1}\right)}
$$

Then, shape functions are given by,

$$
\therefore \quad u(x)=\left[\begin{array}{ll}
N_{1} & N_{2}
\end{array}\right]\left[\begin{array}{l}
u_{1} \\
u_{2}
\end{array}\right]
$$

$$
\begin{aligned}
& \therefore N_{1}=\frac{x_{2}-x}{l}=\frac{x_{2}-x}{x_{2}-x_{1}} \quad\left(\because l=x_{2}-x_{1}\right) \\
& \therefore N_{2}=\frac{x-x_{1}}{l}=\frac{x-x_{1}}{x_{2}-x_{1}}
\end{aligned}
$$



$$
u=N_{1} u_{1}+N_{2} u_{2}
$$



Linear interpolation of the displacement function within an element

## Q. Define the shape function. What are the properties of a shape function?

Ans: The mathematical expression which defines the geometry or shape of the finite element is termed as shape function. They are used to determine the variation of field variables such as displacement, temperature, etc. In finite element method, the problems cannot be solved without using shape functions.

The properties of shape functions are as follows,
(a) The summation of all the shape functions is equal to 1.
(b) The value of each function at its own node is 1 and the value at other node is zero
(c) The shape functions can be linear or quadratic functions, based on the conditions that, first derivative of shape function should be infinite within the element and the displacements across element boundary should be continuous.

Coordinates, penalty approach
Local Coordinates: In local coordinate system, the nodes of various elements of the structure are specified by the origin, which is placed within the element. This type of coordinate system is adopted, in order to minimize the computational efforts while calculating the global stiffness matrix and displacement vectors. Local coordinates may be different for different elements.



Figure (4): Global Coordinate System

Natural coordinate system


Shape functions:
In natural CS:

$$
\begin{aligned}
& S_{i}(\xi)=(1-\xi) / 2 \\
& S_{i}(\xi)=(\xi+1) / 2
\end{aligned}
$$

In LCS:
$S_{i}(x)=(l-x) / l$
$S(x)=x / l$
In GCS:

$$
S_{i}(X)=\left(X_{j}-X\right) /\left(X_{F}-X_{i}\right)
$$

$$
S_{i}(X)=\left(X-X_{i}\right) /\left(X_{F}-X_{i}\right)
$$

```
GCS - {O-X}
LCS-{o-x}
Natural CS - {0'- \xi}
```

Then, shape functions are given by,

$$
\begin{array}{ll}
\therefore N_{1}=\frac{x_{2}-x}{l}=\frac{x_{2}-x}{x_{2}-x_{1}} & N_{1}(\xi)=\frac{1-\xi}{2} \\
\therefore N_{2}=\frac{x-x_{1}}{l}=\frac{x-x_{1}}{x_{2}-x_{1}} & N_{2}(\xi)=\frac{1+\xi}{2}
\end{array}
$$





Linear interpolation of the displacement function within an element

## Or 3 noded bar element Or Quadratic bar element



Figure: Quadratic Bar Element
Boundary Conditions,
At node-1, $u=u_{1}, x_{1}=0$
At node-2, $u=u_{2}, x_{2}=l$
At node- $3, u=u_{3}, x_{3}=\frac{l}{2}$

Consider a quadratic bar element of length ' $l$ '. Let $u_{1}$, $u_{2}, u_{3}$ are the nodal displacement at nodal points $1,2,3$.

The polynomial for one dimensional quadratic bar element is given by,

$$
\begin{equation*}
u=a_{0}+a_{1} x+a_{2} x^{2} \tag{1}
\end{equation*}
$$

Equation (1) can be written in matrix form,

$$
\{u\}=\left[\begin{array}{lll}
1 & x & x^{2}
\end{array}\right]\left\{\begin{array}{l}
a_{0}  \tag{2}\\
a_{1} \\
a_{2}
\end{array}\right\}
$$

Substituting the boundary conditions in eq 1 . and solving for unknowns a0, a1, a2 we get

$$
\{u\}=\left[\begin{array}{lll}
N_{1} & N_{2} & N_{3}
\end{array}\right]\left\{\begin{array}{l}
u_{1} \\
u_{2} \\
u_{3}
\end{array}\right\}
$$

$$
\{u\}=N_{1} u_{1}+N_{2} u_{2}+N_{3} u_{3}
$$

Where,

$$
\begin{aligned}
& N_{1}=1-\frac{3 x}{l}+\frac{2}{l^{2}} x^{2} \\
& N_{2}=\frac{-x}{l}+\frac{2}{l^{2}} x^{2} \\
& N_{3}=\frac{4 x}{l}-\frac{4}{l^{2}} x^{2}
\end{aligned}
$$



Shape Functions for Bar Element


$$
N_{1}=\frac{\xi(g-1)}{2}
$$



$$
N_{2}=\frac{1+\xi}{2}
$$

$N_{2}=\left(1+\xi_{g}\right)\left(1-\xi_{g}\right)$

$$
\begin{aligned}
& N_{1}=-\frac{9}{16}\left(\xi+\frac{1}{3}\right)\left(\xi-\frac{1}{3}\right)(\xi-1) \\
& N_{2}=\frac{27}{16}(\xi+1)\left(\xi-\frac{1}{3}\right)(\xi-1) \\
& N_{3}=-\frac{27}{16}(\xi+1)\left(\xi+\frac{1}{3}\right)(\xi-1) \\
& N_{4}=\frac{9}{16}(\xi+1)\left(\xi+\frac{1}{8}\right)\left(\xi-\frac{1}{3}\right)
\end{aligned}
$$

$$
N_{1} 1
$$

$$
N_{3}=\frac{\xi(\xi+1)}{2}
$$




## Derivation of strain displacement matrix (using 2 noded bar element)

$$
\therefore \quad u(x)=\left[\begin{array}{ll}
N_{1} & N_{2}
\end{array}\right]\left[\begin{array}{l}
u_{1} \\
u_{2}
\end{array}\right]
$$

$$
\begin{aligned}
u & =N_{1} u_{1}+N_{2} u_{2} \\
& =\left(\frac{x_{2}-x}{l}\right) u_{1}+\left(\frac{x-x_{1}}{l}\right) u_{2}
\end{aligned}
$$

The strain for element is defined as,

$$
\varepsilon=\frac{d u}{d x}=\frac{u_{2}-u_{1}}{l}
$$

Also,

$$
\varepsilon=\frac{d u}{d x}=\frac{d N_{1}}{d x} u_{1}+\frac{d N_{2}}{d x} u_{2}
$$

$$
\begin{aligned}
\varepsilon & =\frac{d u}{d x}=\frac{d N_{1}}{d x} u_{1}+\frac{d N_{2}}{d x} u_{2} \\
& =\frac{-1}{l} u_{1}+\frac{1}{l} u_{2} \\
& =\frac{1}{l}\left[\begin{array}{ll}
-1 & 1
\end{array}\right]\left\{\begin{array}{l}
u_{1} \\
u_{2}
\end{array}\right\} \\
\varepsilon & =\frac{1}{\left(x_{2}-x_{1}\right)}\left[\begin{array}{ll}
-1 & 1
\end{array}\right]\left\{\begin{array}{l}
u_{1} \\
u_{2}
\end{array}\right\}
\end{aligned}
$$

$$
\text { i.e., } \varepsilon=[B]\{u\}
$$

Where,

$$
\begin{aligned}
{[B] } & =\frac{1}{\left(x_{2}-x_{1}\right)}\left[\begin{array}{ll}
-1 & 1
\end{array}\right] \\
\therefore[B] & =\frac{1}{l}\left[\begin{array}{ll}
-1 & 1
\end{array}\right]
\end{aligned}
$$

The above equation is known as 'element strain displacement matrix' for one dimensional element.

## Potential Energy Approach

General expression for total potential energy in an elastically loaded structure is,

$$
\pi=\frac{1}{2} \int_{V} \sigma^{T} \varepsilon d V-\int_{V} u^{T} f d v-\int_{A} u^{T} \mathbf{T} d A-\sum_{i} u_{i}^{T} P_{i}
$$

It can be written for the 1D problems as

$$
\Pi=\frac{1}{2} \int_{L} \sigma^{\mathrm{T}} \epsilon A d x-\int_{L} u^{\mathrm{T}} f A d x-\int_{L} u^{\mathrm{T}} T d x-\sum_{i} u_{i} P_{i}
$$

## Element Stiffness Matrix, $\mathbf{K}^{(e)}$

## Derivation of Elemental stiffness matrix

Consider first term in the general expression,
Total strain energy,

$$
\begin{aligned}
& U_{e}=\frac{1}{2} \int_{V} \sigma^{T} \varepsilon d V \\
& U_{e}=\frac{1}{2} \int_{V}(E B u)^{T}(B u) A d x
\end{aligned}
$$

$$
[\because \sigma=E B u \text { and } \varepsilon=B u]
$$

$$
=\frac{1}{2} u^{T} u E_{e} A_{e} B^{T} B \int_{x_{1}}^{\lambda_{2}} d x
$$



Figure (1)
Let, $x_{1}, x_{2}$-Lengths at node-1 and node-2
$u_{1}, u_{2}$ - Displacement vectors

## Where.

$B$ - Strain displacement matrix
$E_{e}$ - Young's modulus of the element
$A_{\text {s }}$ - Cross-sectional area of the element
And.

$$
\begin{aligned}
& \qquad B=\frac{1}{l_{e}}\left[\begin{array}{cc}
-1 & 1
\end{array}\right] \quad\left[\because x_{2}-x_{1}=l_{e}\right] \\
& B^{T}=\frac{1}{l_{e}}\left[\begin{array}{c}
-1 \\
1
\end{array}\right] \\
& B^{T} B=\frac{1}{l_{e}}\left[\begin{array}{c}
-1 \\
1
\end{array}\right] \frac{1}{l_{e}}\left[\begin{array}{ll}
-1 & 1
\end{array}\right] \\
& B^{T} B=\frac{1}{l_{e}^{2}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
& \text { And } \int_{x_{1}}^{x_{2}} d x=[x]_{x_{1}}^{x_{2}}=x_{2}-x_{1}=l_{e}
\end{aligned}
$$

Where,
$l_{\mathrm{e}}$ - Length of the element
$\therefore \quad$ On substituting in equation (1),

$$
\begin{aligned}
& U_{e}=\frac{1}{2} u^{T}\left\{E_{e} A_{e} \frac{1}{l_{e}^{2}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] l_{e}\right\} u \\
& U_{e}=\frac{1}{2} u^{T} K_{e} u
\end{aligned}
$$

Where,
Stiffness matrix, $K_{\mathrm{e}}=\frac{E_{e} A_{e}}{l_{e}}\left[\begin{array}{cc}1 & -1 \\ -1 & 1\end{array}\right]$
And, global stiffness matrix $K_{\text {global }}=\sum_{e} K_{e}$ i.e., summation of individual element stiffness matrices.

## Element Body Load Vector (F ${ }_{\mathrm{e}}$ )

Considering the second term in $\pi_{T}$

$$
\begin{aligned}
\int_{v} u^{T} f d v & =\int_{l}(N u)^{T} f A d x \\
& =\int_{l} u^{T} N^{T} f A d x \\
& =u^{T} A_{a} f \int_{l} N^{T} d x=u^{T} A_{e} f\left\{\begin{array}{c}
\int N_{1} d x \\
\int N_{2} d x
\end{array}\right\} \\
& =u^{T}\left[\frac{A_{e} f l_{e}}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\}\right]\left[\because \int N_{1} d x=\int N_{2} d x=\frac{l_{e}}{2}\right] \\
& =u^{T} . F_{e}
\end{aligned}
$$

Where,

$$
F_{e}=\frac{f A_{e} l_{e}}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\}
$$

## Element Traction Load Vector ( $\mathrm{T}_{\mathrm{e}}$ )

Considering the third term in $\pi_{T}$.

$$
\begin{aligned}
\int_{i} u^{T} T d x & =\int_{i}(N u)^{T} T d x=u^{T}\left\{T \int N^{T} d x\right\} \\
& =u^{T} T\left\{\begin{array}{l}
\int N_{1} d x \\
\int N_{2} d x
\end{array}\right\} \\
\int_{l} u^{T} T d x & =u^{T}\left[T \frac{l_{e}}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\}\right]=u^{T} T_{\mathrm{e}}
\end{aligned}
$$

Where,

$$
\because T_{e}=\frac{T l_{e}}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\}
$$

## Temperature Load Vector

Thermal load= A *hermal stress
If a temperature gradient exists then, temperature load,

$$
\begin{aligned}
\theta & =A E \varepsilon_{0} \\
& =A E \alpha \Delta T
\end{aligned}
$$

Temperature load vector,

$$
\theta_{\mathrm{e}}=A E \alpha \Delta T\left\{\begin{array}{c}
-1 \\
1
\end{array}\right] \begin{gathered}
1 \\
2
\end{gathered}
$$

Refer..Chandrupatla
3.10 TEMPERATURE EFFECTS

$$
U=\int_{L} \frac{1}{2}\left(\epsilon-\epsilon_{0}\right)^{\mathrm{T}} E\left(\epsilon-\epsilon_{0}\right) A d x
$$

$$
F=\sum_{e}\left[F_{e}+T_{e}+\theta_{e}\right]+P_{i}
$$

Element stiffness matrix,

$$
K_{\mathrm{e}}=\frac{E_{e} A_{e}}{l_{e}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]
$$

global stiffness matrix $K_{\text {global }}=\sum_{e} K_{e}$

$$
\mathrm{K}_{\text {(global) }} \mathrm{U}=\mathrm{F}_{\text {(global) }}
$$

Global load vector,

$$
F=\sum_{e}\left(F_{e}+T_{e}+\theta_{e}\right)+P_{i}
$$

Where,

$$
\begin{aligned}
& F_{e}=\frac{f A_{e} l_{e}}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\} \\
& \because T_{e}=\frac{T l_{e}}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\}
\end{aligned}
$$

$$
\theta_{\mathrm{e}}=A_{e} E_{e} \alpha \Delta T\left\{\begin{array}{r}
-1 \\
1
\end{array}\right\}
$$

## Equilibrium Equation

$$
[K]\{u\}=\{F\}
$$

Where,
[K] - Global stiffness matrix

$$
\{u\} \text { - Global displacement vector }
$$

$\{F\}$ - Global load vector
Equation for stress, strain and support reactions are,

$$
\sigma=E \varepsilon, \quad \varepsilon=B u, \quad R=K u=F
$$

If temperature gradient is present then,

$$
\begin{aligned}
\sigma & =E\left(\varepsilon-\varepsilon_{0}\right) \\
& =E[\varepsilon-\alpha \Delta T] \\
\sigma & =E[B u-\alpha \Delta T]
\end{aligned}
$$

## Example



Figure (2)
Temperature gradient, $\Delta T=80^{\circ} \mathrm{C}$

For this problem, calculate nodal displacements, stresses in each bar, Reactions at the supports


## Basic Steps Involved In FEM:

1. Domain Discretization
2. Selection of displacement functions
3. Formation of elemental (stiffness matrix and load vector)
4. Formation of Global (stiffness matrix anc load vector) : K U = F
5. Application of boundary condition
6. Solution of simultaneous equations (for unknown nodal displacements )
7. Calculation of stresses and strains
8. Interpretation of results


Boundary conditions, $u_{1}=u_{4}=0$

Element stiffness matrix,

$$
K_{\mathrm{e}}=\frac{E_{e} A_{e}}{l_{e}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]
$$

Stiffness matrix of element-1,

$$
\begin{aligned}
& K_{1}=\frac{83 \times 10^{3} \times 2400}{800}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
& K_{1}=10^{5}\left[\begin{array}{cc}
1 & 2.49 \\
-2.49 \\
-2.49 & 2.49
\end{array}\right]\left[\begin{array}{l}
1 \\
2
\end{array}\right.
\end{aligned}
$$

Stiffness matrix of element-2,

$$
\begin{aligned}
& K_{2}=\frac{70 \times 10^{3} \times 1200}{600}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
& K_{2}=10^{5}\left[\begin{array}{cc}
1.4 & -1.4 \\
-1.4 & 1.4
\end{array}\right]
\end{aligned}
$$

Stiffness matrix of element -3 ,

$$
\begin{aligned}
& K_{3}=\frac{200 \times 10^{3} \times 600}{400}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
& K_{3}=10^{5}\left[\begin{array}{cc}
3 & 4 \\
3 & -3 \\
-3 & 3
\end{array}\right] \begin{array}{l}
3 \\
4
\end{array}
\end{aligned}
$$



Global stiffness matrix,

$$
\left.\begin{array}{rl}
K & =10^{5}\left[\begin{array}{cccc}
1 & 2 & 3 & 4 \\
2.49 & -2.49 & 0 & 0 \\
-2.49 & 2.49+1.4 & -1.4 & 0 \\
0 & -1.4 & 1.4+3 & -3 \\
0 & 0 & -3 & 3
\end{array}\right] 4 \\
& =10^{5}\left[\begin{array}{cccc}
1 & 2 & 3 & 4 \\
2.49 & -2.49 & 0 & 0 \\
-2.49 & 3.89 & -1.4 & 0 \\
0 & -1.4 & 4.4 & -3 \\
0 & 0 & -3 & 3
\end{array}\right] 1
\end{array}\right] \begin{aligned}
& 1 \\
& 3 \\
& 4
\end{aligned}
$$

Global load vector,

$$
F=\sum_{e}\left(F_{e}+T_{e}+\theta_{e}\right)+P_{i}
$$

Temperature load on element-1,

$$
\begin{aligned}
& \theta_{1}=2400 \times 83 \times 10^{3} \times 18.9 \times 10^{-6} \times 80\left\{\begin{array}{c}
-1 \\
1
\end{array}\right\} \\
& \theta_{1}=10^{5}\left\{\begin{array}{c}
-3.012 \\
3.012
\end{array}\right\} 2
\end{aligned}
$$

Temperature load on element-2,

$$
\begin{aligned}
& \theta_{2}=1200 \times 70 \times 10^{3} \times 23 \times 10^{-6} \times 80\left\{\begin{array}{c}
-1 \\
1
\end{array}\right\} \\
& \theta_{2}=10^{5}\left\{\begin{array}{c}
-1.54 \\
1.54
\end{array}\right\} 3
\end{aligned}
$$

Temperature load on element-3,

$$
\begin{aligned}
& \theta_{3}=600 \times 200 \times 10^{3} \times 11.7 \times 10^{-6} \times 80\left\{\begin{array}{c}
-1 \\
1
\end{array}\right\} \\
& \theta_{3}=10^{5}\left\{\begin{array}{c}
-1.12 \\
1.12
\end{array}\right\} 4
\end{aligned}
$$

Global load vector,

Points loads,
At node- $1=0$

$$
F=\sum_{e}\left(F_{e}+T_{e}+\theta_{e}\right)+P_{i}
$$

At node- $2=-60 \times 10^{3} \mathrm{~N}$
At node-3 $=-75 \times 10^{3} \mathrm{~N}$
At node-4 $=0$

$$
\begin{aligned}
& \{F\}=\left\{\begin{array}{l|l}
-3.012 \times 10^{5} \\
(3.012-1.54-0.6) \times 10^{5} & 1 \\
(1.54-1.12-0.75) \times 10^{5} & 2 \\
1.12 \times 10^{5} & 4 \\
4
\end{array}\right. \\
& \{F\}=10^{5}\left\{\begin{array}{c|c|c}
-3.012 & 1 \\
0.872 & 2 \\
-0.33 & 3 \\
1.12 & 4
\end{array}\right.
\end{aligned}
$$

Then,

$$
\begin{aligned}
{[K]\{u\} } & =\{F\} \\
& 10^{5}\left[\begin{array}{cccc}
2.49 & -2.49 & 0 & 0 \\
-2.49 & 3.89 & -1.4 & 0 \\
0 & -1.4 & 4.4 & -3 \\
0 & 0 & -3 & 3
\end{array}\right]\left[\begin{array}{l}
u_{1} \\
u_{2} \\
u_{3} \\
u_{4}
\end{array}\right\}=10^{5}\left\{\begin{array}{c}
-3.012 \\
0.872 \\
-0.33 \\
1.12
\end{array}\right\}
\end{aligned}
$$

Applying the boundary conditions, $u_{1}=u_{4}=0$.
By elimination approach, the above equation reduces to,

$$
\begin{aligned}
10^{5}\left\{\begin{array}{cc}
3.89 & -1.4 \\
-1.4 & 4.4
\end{array}\right\}\left\{\begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right\} & =10^{5}\left\{\begin{array}{c}
0.872 \\
-0.33
\end{array}\right\} \\
3.89 u_{2}-1.4 u_{3} & =0.872 \\
-1.4 u_{2}+4.4 u_{3} & =-0.33
\end{aligned}
$$

Solving the above two equations,

$$
\begin{aligned}
& u_{2}=0.223 \mathrm{~mm} \\
& u_{3}=-0.00415 \mathrm{~mm}
\end{aligned}
$$

$\therefore \quad$ Global displacement vector, $u=\left[\begin{array}{lll}0 & 0.223 & -0.00415\end{array}\right]^{T} \mathrm{~mm}$

Calculation of stresses,

$$
\sigma=E(B u-\alpha \Delta T)
$$

Stress induced in element-1,

$$
\sigma_{1}=83 \times 10^{3}\left[\frac{1}{800}\left[\begin{array}{ll}
-1 & 1
\end{array}\right]\left[\begin{array}{c}
0 \\
0.223
\end{array}\right]-18.9 \times 10^{-6} \times 80\right]
$$

$\therefore \sigma_{1}=-102.366 \mathrm{MPa}$
Stress induced in element-2,


$$
\begin{aligned}
\sigma_{2} & =70 \times 10^{3}\left[\frac{1}{600}\left[\begin{array}{ll}
-1 & 1
\end{array}\right]\left[\begin{array}{c}
0.223 \\
-0.004
\end{array}\right]-23 \times 10^{-6} \times 80\right] \\
\therefore \sigma_{2} & =-155.28 \mathrm{MPa}
\end{aligned}
$$

Stress induced in element-3,

$$
\begin{aligned}
& \sigma_{3}=200 \times 10^{3}\left[\frac{1}{400}\left[\begin{array}{ll}
-1 & 1
\end{array}\right]\left[\begin{array}{c}
-0.004 \\
0
\end{array}\right]-11.7 \times 10^{-6} \times 80\right] \\
& \therefore \sigma_{3}=-185.2 \mathrm{MPa}
\end{aligned}
$$

Solving for the support reactions,

$$
\begin{aligned}
\mathrm{K} 11 \mathrm{u} 1+K_{12} u_{2}+K_{13} u_{3}+\mathrm{K} 14 \mathrm{u} 4 & =F_{1}+R_{1} \\
\mathrm{~K} 41 \mathrm{u} 1+K_{42} u_{2}+K_{43} u_{3}+\mathrm{K} 44 \mathrm{u} 4 & =F_{4}+R_{4} \\
(-2.49 \times 0.223) \times 10^{5} & =R_{1}+(-3.012) \times 10^{5} \\
\therefore \quad R_{1} & =245.6 \mathrm{kN} \\
(-3 \times-0.004) \times 10^{5} & =R_{4}+\left(1.12 \times 10^{5}\right) \\
\therefore R_{4} & =-110.8 \mathrm{kN}
\end{aligned}
$$

## Example 3.8

An axial load $P=300 \times 10^{3} \mathrm{~N}$ is applied at $20^{\circ} \mathrm{C}$ to the rod as shown in Fig. E3.8. The temperature is then raised to $60^{\circ} \mathrm{C}$.
(a) Assemble the $\mathbf{K}$ and $\mathbf{F}$ matrices.
(b) Determine the nodal displacements and element stresses.


FIGURE E3.8

## Solution

(a) The element stiffness matrices are

$$
\begin{aligned}
& \mathbf{k}^{1}=\frac{70 \times 10^{3} \times 900}{200}\left[\begin{array}{rr}
1 & -1 \\
-1 & 1
\end{array}\right] \mathrm{N} / \mathrm{mm} \\
& \mathbf{k}^{2}=\frac{200 \times 10^{3} \times 1200}{300}\left[\begin{array}{rr}
1 & -1 \\
-1 & 1
\end{array}\right] \mathrm{N} / \mathrm{mm}
\end{aligned}
$$

Thus,

$$
\mathbf{K}=10^{3}\left[\begin{array}{ccc}
315 & -315 & 0 \\
-315 & 1115 & -800 \\
0 & -800 & 800
\end{array}\right] \mathrm{N} / \mathrm{mm}
$$

Now, in assembling $\mathbf{F}$, both temperature and point load effects have to be considered. The element temperature forces due to $\Delta T=40^{\circ} \mathrm{C}$ are obtained from Eq. 3.106b as

$$
\boldsymbol{\theta}^{1}=70 \times 10^{3} \times 900 \times 23 \times 10^{-6} \times 40\left\{\begin{array}{c}
-1 \\
1
\end{array}\right\}_{2}^{\downarrow} \mathrm{N}
$$

and

$$
\boldsymbol{\theta}^{2}=200 \times 10^{3} \times 1200 \times 11.7 \times 10^{-6} \times 40\left\{\begin{array}{c}
-1 \\
1
\end{array}\right\}_{3}^{2} \mathrm{~N}
$$

Upon assembling $\boldsymbol{\theta}^{1}, \boldsymbol{\theta}^{2}$, and the point load, we get

$$
\mathbf{F}=10^{3}\left\{\begin{array}{l}
-57.96 \\
57.96-112.32+300 \\
112.32
\end{array}\right\}
$$

or

$$
\mathbf{F}=10^{3}[-57.96, \quad 245.64, \quad 112.32]^{\mathrm{T}} \mathrm{~N}
$$

(b) The elimination approach will now be used to solve for the displacements. Since dof 1 and 3 are fixed, the first and third rows and columns of $\mathbf{K}$, together with the first and third components of $\mathbf{F}$, are deleted. This results in the scalar equation

$$
10^{3}[1115] Q_{2}=10^{3} \times 245.64
$$

yielding

$$
Q_{2}=0.220 \mathrm{~mm}
$$

Thus,

$$
\mathbf{Q}=\left[\begin{array}{lll}
0, & 0.220, & 0
\end{array}\right]^{\mathrm{T}} \mathrm{~mm}
$$

In evaluating element stresses, we have to use Eq. 3.108b:

$$
\begin{aligned}
\sigma_{1} & =\frac{70 \times 10^{3}}{200}\left[\begin{array}{ll}
-1 & 1
\end{array}\right]\left\{\begin{array}{c}
0 \\
0.220
\end{array}\right\}-70 \times 10^{3} \times 23 \times 10^{-6} \times 40 \\
& =12.60 \mathrm{MPa}
\end{aligned}
$$

and

$$
\begin{aligned}
\sigma_{2} & =\frac{200 \times 10^{3}}{300}\left[\begin{array}{ll}
-1 & 1
\end{array}\right]\left\{\begin{array}{c}
0.220 \\
0
\end{array}\right\}-200 \times 10^{3} \times 11.7 \times 10^{-6} \times 40 \\
& =-240.27 \mathrm{MPa}
\end{aligned}
$$

Concept of assembly


Figure (1): Bar Element
$[K]\{U\}=\{F\}$

$$
\left.\left.\begin{array}{rl}
\frac{A E}{l}\left[\begin{array}{ccccc}
1 & 2 & 3 & 4 & 5 \\
1 & -1 & 0 & 0 & 0 \\
-1 & 1+1 & -1 & 0 & 0 \\
0 & -1 & 1+1 & -1 & 0 \\
0 & 0 & -1 & 1+1 & -1 \\
0 & 0 & 0 & -1 & 1
\end{array}\right] 1 \\
2 \\
5
\end{array} \right\rvert\, \begin{array}{c}
u_{1} \\
u_{2} \\
u_{3} \\
u_{4} \\
u_{5}
\end{array}\right\}=\left\{\begin{array}{c}
F_{1} \\
F_{2} \\
F_{3} \\
F_{4} \\
F_{5}
\end{array}\right\}
$$

Using two finite elements, find the stress distribution in a uniformly tapering bar of cross sectional area $300 \mathrm{~mm}^{2}$ and $200 \mathrm{~mm}^{2}$ at their ends, length 100 mm , subjected to an axial tensile load of 50 N at smaller end and fixed at larger end. Take $E=2 \times 10^{5} \mathrm{~N} / \mathrm{mm}^{2}$.


$$
\begin{aligned}
a_{1} & =300 \mathrm{~mm}^{2} \\
a_{3} & =200 \mathrm{~mm}^{2} \\
a_{2} & =\frac{a_{1}+a_{3}}{2} \\
& =\frac{300+200}{2} \\
a_{2} & =250 \mathrm{~mm}^{2}
\end{aligned}
$$

Given that,
Cross sectional area of bigger end, $a_{1}=300 \mathrm{~mm}^{2}$ Cross sectional area of smaller end, $a_{3}=200 \mathrm{~mm}^{2}$ Length of the bar, $l=100 \mathrm{~mm}$

Axial tensile load, $P=50 \mathrm{~N}$
Young's modulus, $E=2 \times 10^{5} \mathrm{~N} / \mathrm{mm}^{2}$


Figure (1): Tapered Bar


Element (1)


Figure (3): Element (1)

Finite element equation for element (1) is given by,

$$
\begin{align*}
& \frac{A_{1} E_{1}}{l_{1}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]\left\{\begin{array}{l}
u_{1} \\
u_{2}
\end{array}\right\}=\left\{\begin{array}{l}
F_{1} \\
F_{2}
\end{array}\right\} \\
& \frac{275 \times 2 \times 10^{5}}{50}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]\left[\left\{\begin{array}{l}
u_{1} \\
u_{2} \\
2
\end{array}\right\}=\left\{\begin{array}{l}
F_{1} \\
F_{2}
\end{array}\right\}\right. \\
& 11 \times 10^{5}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]\left\{\begin{array}{l}
u_{1} \\
u_{2}
\end{array}\right\}=\left\{\begin{array}{l}
F_{1} \\
F_{2}
\end{array}\right\} \\
& 10^{5}\left[\begin{array}{cc}
1 & 2 \\
11 & -11 \\
-11 & 11
\end{array}\right]\left[\begin{array}{l}
1 \\
u_{1} \\
u_{2} \\
u_{2}
\end{array}\right\}=\left\{\begin{array}{c}
F_{1} \\
F_{2}
\end{array}\right\} \tag{2}
\end{align*}
$$

Element (2)


Figure (4): Element (2)
Finite element equation for element (2) is given by,

$$
\begin{aligned}
& \frac{A_{2} E_{2}}{l_{2}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]\left\{\begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{l}
F_{2} \\
F_{3}
\end{array}\right\} \\
& \frac{225 \times 2 \times 10^{5}}{50}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]\left\{\begin{array}{l}
u_{2} \\
u_{3} \\
3
\end{array}\right\}=\left\{\begin{array}{c}
F_{2} \\
F_{3}
\end{array}\right\} \\
& 9 \times 10^{5}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]\left\{\begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
F_{2} \\
F_{3}
\end{array}\right\} \\
& 10^{5}\left[\begin{array}{cc}
9 & -9 \\
-9 & 9
\end{array}\right]\left[\begin{array}{l}
3
\end{array} \left\lvert\, \begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right.\right\}=\left\{\begin{array}{c}
F_{2} \\
F_{3}
\end{array}\right\}
\end{aligned}
$$

From equation (1) and (2),
Finite element equation is given by,

$$
\begin{gather*}
1 c^{2} 3 \\
10^{5}\left[\begin{array}{ccc}
11 & -11 & 0 \\
-11 & 11+9 & 9 \\
0 & -9 & 9
\end{array}\right]\left[\begin{array}{l}
u_{1} \\
2 \\
1 \\
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{l}
F_{1} \\
F_{2} \\
F_{3}
\end{array}\right\}  \tag{3}\\
10^{5}\left[\begin{array}{ccc}
11 & -11 & 0 \\
-11 & 0 & -9 \\
0 & -9 & 9
\end{array}\right]\left[\begin{array}{l}
u_{1} \\
u_{2} \\
u_{3}
\end{array}\right]=\left\{\begin{array}{l}
F_{1} \\
F_{2} \\
F_{3}
\end{array}\right\}
\end{gather*}
$$

Applying the boundary conditions,

$$
u_{1}=0, F_{1}=0, F_{2}=0, F_{3}=50 \mathrm{~N}
$$

From equation (3),

$$
10^{5}\left[\begin{array}{ccc}
11 & -11 & 0 \\
-11 & 20 & -9 \\
0 & -9 & 9
\end{array}\right]\left[\begin{array}{l}
u_{1} \\
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
0 \\
50
\end{array}\right\}
$$

Eliminating the first row and first column,
Since $u_{1}=0$

$$
\begin{array}{lc}
10^{5}\left[\begin{array}{cc}
20 & -9 \\
-9 & 9
\end{array}\right]\left[\begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
50
\end{array}\right\} & \text { On solving the above equations, } \\
10^{5}\left(20 u_{2}-9 u_{3}\right)=0 & u_{2}=4.545 \times 10^{-5} \\
10^{5}\left(-9 u_{2}+9 u_{3}\right)=50 & u_{3}=10.101 \times 10^{-5}
\end{array}
$$

## Stress Distribution

## For element (1)

$$
\text { Stress, } \begin{aligned}
\sigma_{1} & =E_{1} \times \frac{u_{2}-u_{1}}{I_{1}} \\
& =2 \times 10^{5} \times \frac{\left(4.545 \times 10^{-5}-0\right)}{50} \\
\sigma_{1} & =0.1818 \mathrm{~N} / \mathrm{mm}^{2}
\end{aligned}
$$

## For element (2)

$$
\text { Stress, } \begin{aligned}
\sigma_{2} & =E_{2} \times \frac{u_{3}-u_{2}}{l_{2}} \\
& =2 \times 10^{5} \times \frac{\left(10.10 \times 10^{-5}-4.545 \times 10^{-5}\right)}{50} \\
\sigma_{2} & =0.222 \mathrm{~N} / \mathrm{mm}^{2}
\end{aligned}
$$

What is study state heat transfer analysis? Write its governing Equation?

Steady state heat transfer is defined as the temperature at any point in the medium does not change with time.

For a one dimensional steady state heat transfer,

$$
K \cdot \frac{d^{2} T}{d x^{2}}+q=0
$$

$K$ - Thermal conductivity
$T$ - Temperature
$q$ - Internal heat source per unit volume

## 3D Conduction heat transfer

General 3D conduction Equation:

$$
\frac{\partial}{\partial x}\left(k \frac{\partial T}{\partial x}\right)+\frac{\partial}{\partial y}\left(k \frac{\partial T}{\partial y}\right)+\frac{\partial}{\partial z}\left(k \frac{\partial T}{\partial z}\right)+\dot{q}=\rho c \frac{\partial T}{\partial \tau}
$$

For constant conductivity:

$$
\begin{aligned}
\frac{\partial^{2} T}{\partial x^{2}} & +\frac{\partial^{2} T}{\partial y^{2}}+\frac{\partial^{2} T}{\partial z^{2}}+\frac{\dot{q}}{k}=\frac{1}{\alpha} \frac{\partial T}{\partial \tau} \\
\alpha & =k / \rho c \\
& =\text { Thermal diffusivity of a material }
\end{aligned}
$$

Q. Give the finite element equation for a one dimensional heat conduction element.

Ans: The finite element equation for a one dimensional heat conduction element is given by,

$$
\{F\}=\left[K_{c}\right]\{T\}
$$

$\{F\}$ - Force vector
$=\left\{\begin{array}{c}F_{1} \\ F_{2}\end{array}\right\}$ for a two noded element
$\left[K_{c}\right]$ - Stiffness matrix in case of heat conduction

$$
=\frac{K A}{l}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]
$$

$\{T\}$ - Nodal temperature vector

$$
=\left\{\begin{array}{l}
T_{1} \\
T_{2}
\end{array}\right\} \text { for a two noded element }
$$

## Similar to structural problems

- Formulation of element stiffness matrices
- Assembly of global stiffness matrix [K]
- formulation of load vector $\{F\}$
- Solution of $[\mathrm{K}]\{\mathrm{u}\}=\{\mathrm{F}\}$ to get nodal displacements $\{\mathrm{u}\}$


## Unit 2

Trusses


FIGURE 4.1 A rwo-dimunsional trass



Figure(1): Pin Jointed Bar Element
$x, y$-Global co-ordinates $m, n$-Local co-ordinates
$q_{1}, q_{2}$-Displacement at nodes 1,2 in the local co-ordinate system ( $m, n$ directions)


Figure(2): Components of Nodal Displacements
$v_{1}, u_{1}, v_{2}, u_{2}$-Components of nodal displacements $q_{1}, q_{2}$ in the global co-ordinate system ( $x, y$ directions)

$$
\{q\}=\left\{\begin{array}{l}
q_{1} \\
q_{2}
\end{array}\right\} \quad \begin{aligned}
& q_{1}=u_{1} \cos \theta+v_{1} \sin \theta \\
& q_{2}=u_{2} \cos \theta+v_{2} \sin \theta
\end{aligned}
$$

$$
\{q\}=\left\{\begin{array}{l}
q_{1} \\
q_{2}
\end{array}\right\} \quad \begin{aligned}
& q_{1}=u_{1} \cos \theta+v_{1} \sin \theta \\
& q_{2}=u_{2} \cos \theta+v_{2} \sin \theta
\end{aligned}
$$

$$
\left\{\begin{array}{l}
q_{1} \\
q_{2}
\end{array}\right\}=\left[\begin{array}{cccc}
\cos \theta & \sin \theta & 0 & 0 \\
0 & 0 & \cos \theta & \sin \theta
\end{array}\right]\left[\begin{array}{l}
u_{1} \\
v_{1} \\
u_{2} \\
v_{2}
\end{array}\right\}
$$

Let $C=\cos \theta, S=\sin \theta$

$$
\left\{\begin{array}{l}
q_{1} \\
q_{2}
\end{array}\right\}=\left[\begin{array}{llll}
C & S & 0 & 0 \\
0 & 0 & C & S
\end{array}\right]\left\{\begin{array}{l}
u_{1} \\
v_{1} \\
u_{2} \\
v_{2}
\end{array}\right\}
$$

$$
\begin{aligned}
& \{q\}=\left\{\begin{array}{l}
q_{1} \\
q_{2}
\end{array}\right\}=\left[\begin{array}{llll}
C & S & 0 & 0 \\
0 & 0 & C & S
\end{array}\right]\left\{\begin{array}{l}
u_{1} \\
v_{1} \\
u_{2} \\
v_{2}
\end{array}\right\} \\
& \{q\}=[L]\left\{\delta^{\prime}\right\}
\end{aligned}
$$

Where,

$$
[L]=\left[\begin{array}{llll}
C & S & 0 & 0 \\
0 & 0 & C & S
\end{array}\right]
$$

Transformation matrix

## Strain Energy

$$
\begin{aligned}
& U=\frac{1}{2}\{q\}^{T}\left[K_{l}\right]\{q\} \\
& U=\frac{1}{2}\left\{\delta^{\prime}\right\}^{T}[L]^{T}\left[K_{l}\right][L]\left\{\delta^{\prime}\right\} \\
& U=\frac{1}{2}\left\{\delta^{\prime}\right\}^{T}[K]\left\{\delta^{\prime}\right\}
\end{aligned}
$$

Truss is a one dimensional element in the local coordinate system. Therefore, element stiffness matrix of a truss element in local co-ordinate system is given by,

$$
\left[K_{l}\right]=\frac{A E}{l}\left[\begin{array}{rr}
1 & -1 \\
-1 & 1
\end{array}\right]
$$

Where

$$
\begin{aligned}
{[K] } & =[L]^{\mathrm{T}}\left[K_{]}\right][L] \\
& =\text { Element stiffness matrix in global co-ordinate } \\
& \text { system }
\end{aligned}
$$

$A$ - Area of cross-section of truss element
$E$ - Young's modulus
$l$ - Length of truss element

$$
\begin{aligned}
& {[K] }=[L]^{T}[K][L] \\
& {[K] }=\left[\begin{array}{ll}
C & 0 \\
S & 0 \\
0 & C \\
0 & S
\end{array}\right] \frac{A E}{l}\left[\begin{array}{rr}
1 & -1 \\
-1 & 1
\end{array}\right]\left[\begin{array}{llll}
C & S & 0 & 0 \\
0 & 0 & C & S
\end{array}\right] \\
&=\frac{A E}{l}\left[\begin{array}{ll}
C & 0 \\
S & 0 \\
0 & C \\
0 & S
\end{array}\right]\left[\begin{array}{llll}
C & S & -C & -S \\
-C & -S & C & S
\end{array}\right] \\
& {[K] }=\frac{A E}{l}\left[\begin{array}{llll}
C^{2} & C S & -C^{2} & -C S \\
C S & S^{2} & -C S & -S^{2} \\
-C^{2} & -C S & C^{2} & C S \\
-C S & -S^{2} & C S & S^{2}
\end{array}\right] \quad l\left(x_{2}, y_{2}\right) \\
& \ell=\cos \theta=\frac{x_{2}-x_{1}}{\ell_{e}} \\
& \ell=\cos \phi=\frac{y_{2}-y_{1}}{\ell_{e}}(=\sin \theta) \\
& \ell=\sqrt{\left(x_{2}-x_{1}\right)^{2}+\left(y_{2}-y_{1}\right)^{2}} \\
& \text { Let } C=\cos \theta, S=\sin \theta
\end{aligned}
$$

| Bar Element |  | Truss Element |  |
| :---: | :---: | :---: | :---: |
| 1. | Displacements of a loaded bar element occurs only in x -direction (horizontal). | 1. | The joint displacements of a loaded truss element are neither horizontal nor vertical, but they are resolved into horizontal and vertical components. |
| 2. | The loads in the bar elements are applied in axial direction. | 2. | The load applied in the truss elements are either compression or tension. |
| 3. | For joining bar elements, various types of weldings are used. | 3. | To join truss elements, pin joints are used. |
| 4. | Stiffness matrix for bar element is given by, $[K]=\frac{E A}{L}\left[\begin{array}{cc} 1 & -1 \\ -1 & 1 \end{array}\right]$ | 4. | Stiffness matrix for truss element is given by, $K=\frac{E A}{L}\left[\begin{array}{cccc} C^{2} & C S & -C^{2} & -C S \\ C S & S^{2} & -C S & -S^{2} \\ -C^{2} & -C S & C^{2} & C S \\ -C S & -S^{2} & C S & S^{2} \end{array}\right]$ <br> Where, $C=\cos \theta$ and $S=\sin \theta$. |
| 5. | Stress in a bar element is, $\sigma=E \varepsilon \text { i.e., } \sigma=E[B]\{\delta\}$ <br> Where, [B]-Strain displacement matrix <br> $\{\delta\}$ - Nodal displacement. | 5. | Stress in a truss element is, $\begin{aligned} \sigma & =E \varepsilon \\ & =\frac{E}{L}\left[\begin{array}{llll} -C & -S & C & S \end{array}\right] u \end{aligned}$ <br> Where, u - Nodal displacement. |

GAX
vesticex $\cos |\Delta| \Delta$










The plane truss shown in figure is composed of members having $0.1 \mathrm{~m}^{2}$ cross-sectional area and modulus of elasticity $\mathrm{E}=70 \mathrm{GPa}$,
(a) Assemble the global stiffness matrix
(b) Compute the nodal displacements in the global coordinate system.


Figure

## Given that,

Area of truss members, $A=0.1 \mathrm{~m}^{2}$
Young's modulus, $E=70 \mathrm{GPa}=70 \times 10^{9} \mathrm{~N} / \mathrm{m}^{2}$


Let $u_{1}, v_{1}, u_{2}, v_{2}, u_{3}$ and $v_{3}$ are the nodal displacements in ' $x$ ' and ' $y$ ' direction respectively.

The stiffness matrix for truss element is given as,

$$
[K]=\frac{A E}{L}\left[\begin{array}{cccc}
C^{2} & C S & -C^{2} & -C S \\
C S & S^{2} & -C S & -S^{2} \\
-C^{2} & -C S & C^{2} & C S \\
-C S & -S^{2} & C S & S^{2}
\end{array}\right]
$$

Where, $C=\cos \theta$ and $S=\sin \theta$

Figure (1): Truss Element

## For Element-1

$$
\begin{aligned}
& \theta_{1}=0^{\circ} \\
& C=\cos \theta_{1}=\cos 0^{\circ}=1 \\
& S=\sin \theta_{1}=\sin 0^{\circ}=0
\end{aligned}
$$

$$
\begin{aligned}
& {\left[K_{1}\right] }=\frac{0.1 \times 70 \times 10^{9}}{1}\left[\begin{array}{cccc}
1 & 0 & -1 & 0 \\
0 & 0 & 0 & 0 \\
-1 & 0 & 1 & 0 \\
0 & 0 & 0 & 0
\end{array}\right] \\
& u_{1} v_{1} u_{2} v_{2} \\
&=7 \times 10^{9} \cdot\left[\begin{array}{cccc}
1 & 0 & -1 & 0 \\
0 & 0 & 0 & 0 \\
-1 & 0 & 1 & 0 \\
0 & 0 & 0 & 0
\end{array}\right] v_{1}^{u_{1}} \\
& v_{1}
\end{aligned}
$$

## For Element-2

$$
\begin{aligned}
\theta_{2} & =90^{\circ} \text { defined from node } 1 \\
C & =\cos \theta_{2}=\cos 90^{\circ}=0 \\
S & =\sin \theta_{2}=\sin 90^{\circ}=1 \\
{\left[K_{2}\right] } & \left.\left.=\frac{0.1 \times 70 \times 10^{9}\left[\begin{array}{cccc}
0 & 0 & 0 & 0 \\
0 & 1 & 0 & -1 \\
0 & 0 & 0 & 0 \\
0 & -1 & 0 & 1
\end{array}\right]}{} \begin{array}{llll}
u_{1} & v_{1} & u_{3} & v_{3} \\
0 & 0 & 0 & 0 \\
0 & 1 & 0 & -1 \\
0 & 0 & 0 & 0 \\
0 & -1 & 0 & 1
\end{array}\right]\right]_{1}^{u_{1}} v_{1} \\
& =7 \times 10_{3}
\end{aligned}
$$

## For Element-3

$$
\theta_{3}=-45^{\circ}(\mathrm{cw})
$$

$$
=135^{\circ}(\mathrm{ccw})
$$

$$
\theta_{3}=135^{\circ} \text { defined from node } 2 \text { to } 3 .
$$

$$
\begin{aligned}
C & =\cos \theta_{3}=\cos \left(135^{\circ}\right)=\frac{-\sqrt{2}}{2}=-0.707 \\
S & =\sin \theta_{3}=\sin \left(135^{\circ}\right)=\frac{\sqrt{2}}{2}=0.707 \\
{\left[K_{3}\right] } & =\frac{0.1 \times 70 \times 10^{9}}{\sqrt{2}} \times\left[\begin{array}{cccc}
0.5 & -0.5 & -0.5 & 0.5 \\
-0.5 & 0.5 & 0.5 & -0.5 \\
-0.5 & 0.5 & 0.5 & -0.5 \\
0.5 & -0.5 & -0.5 & 0.5
\end{array}\right] \\
& =7 \times 10^{u_{2}}\left[\begin{array}{cccc}
0.354 & -0.354 & -0.354 & 0.354 \\
-0.354 & 0.354 & 0.354 & -0.354 \\
-0.354 & 0.354 & 0.354 & -0.354 \\
0.354 & -0.354 & -0.354 & 0.354
\end{array}\right] \begin{array}{l}
u_{2} \\
v_{2} \\
u_{3} \\
v_{3}
\end{array}
\end{aligned}
$$



Figure (2)

The global stiffness matrix is calculated as,

$$
[K]=\left[K_{1}\right]+\left[K_{2}\right]+\left[K_{3}\right]
$$



Global displacement vector, Global force vector,

$$
\left\{\delta^{\prime}\right\}=\left\{\begin{array}{l}
u_{1} \\
v_{1} \\
u_{2} \\
v_{2} \\
u_{3} \\
v_{3}
\end{array}\right\} \quad\{F\}=\left\{\begin{array}{l}
F_{1 x} \\
F_{1 y} \\
F_{2 x} \\
F_{2 y} \\
F_{3 x} \\
F_{3 y}
\end{array}\right\}
$$

Nodal Displacements: The finite element matrix equation for truss structure can be written as,

$$
\begin{aligned}
& {[K]\left\{\delta^{\prime}\right\}=\{F\}} \\
& 7 \times 10^{9}\left[\begin{array}{ccrrcc}
1 & 0 & -1 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 & -1 \\
-1 & 0 & 1.354 & -0.354 & -0.354 & 0.354 \\
0 & 0 & -0.354 & 0.354 & 0.354 & -0.354 \\
0 & 0 & -0.354 & 0.354 & 0.354 & -0.354 \\
0 & -1 & 0.354 & -0.354 & -0.354 & 1.354
\end{array}\right]\left\{\begin{array}{l}
u_{1} \\
v_{1} \\
u_{2} \\
v_{2} \\
u_{3} \\
v_{3}
\end{array}\right\}=\left\{\begin{array}{l}
\mathbf{F}_{1 \mathrm{x}} \\
\mathrm{~F}_{1 \mathrm{y}} \\
\mathrm{~F}_{2 \mathrm{x}} \\
\mathrm{~F}_{2 \mathrm{y}} \\
\mathrm{~F}_{3 \mathrm{x}} \\
\mathrm{~F}_{3 \mathrm{y}}
\end{array}\right\}
\end{aligned}
$$



Applying boundary conditions,

$$
\text { i.e., } u_{1}=v_{1}=u_{3}=0
$$

Hence, omit $1^{\text {st }}, 2^{\text {nd }}, 5^{\text {th }}$ row and columns in finite element equations,

$$
7 \times 10^{9}\left[\begin{array}{ccc}
1.354 & -0.354 & 0.354 \\
-0.354 & 0.354 & -0.354 \\
0.354 & -0.354 & 1.354
\end{array}\right]\left\{\begin{array}{c}
u_{2} \\
v_{2} \\
v_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
-1000 \\
0
\end{array}\right\}
$$

$$
\begin{aligned}
& \left\{\begin{array}{l}
u_{2} \\
v_{2} \\
v_{3}
\end{array}\right\}=\left[\begin{array}{ccc}
1.354 & -0.354 & 0.354 \\
-0.354 & 0.354 & -0.354 \\
0.354 & -0.354 & 1.354
\end{array}\right]^{-1}\left\{\begin{array}{c}
0 \\
-1000 \\
0
\end{array}\right\} \times \frac{1}{7 \times 10^{9}} \\
& \left\{\begin{array}{l}
u_{2} \\
v_{2} \\
v_{3}
\end{array}\right\}=\left\{\begin{array}{c}
-1000 \\
-4824.859 \\
-1000
\end{array}\right] \times \frac{1}{7 \times 10^{9}} \mathrm{~m} \\
& \left\{\begin{array}{l}
u_{2} \\
v_{2} \\
v_{3}
\end{array}\right\}=\left\{\begin{array}{c}
142.857 \\
689.266 \\
142.857
\end{array}\right\} \times 10^{-6} \mathrm{~mm}
\end{aligned}
$$

For the two bar truss as shown in figure determine the displacements at node 2 and the stresses in both elements.


Figure

## Given that,

Young's modulus, $E=70 \mathrm{GPa}=70 \times 10^{3} \mathrm{~N} / \mathrm{mm}^{2}$

$$
\text { Area, } A=200 \mathrm{~mm}^{2}
$$

Truss is divided into two elements as shown in figure.


Figure(1)

Let,
$u_{1}, u_{2}, u_{3}$-Displacements along $x$-axis at nodes 1,2 and 3 respectively.
$v_{1}, v_{2}, v_{3}$-Displacements along $y$-axis at nodes 1,2 and 3 respectively.
$F_{1 x}, F_{2 x}, F_{3 x}-$ Forces along $x$-axis at nodes 1,2 and 3 respectively.
$F_{1 y}, F_{2 y}, F_{3 y}$-Forces along $x$-axis at nodes 1,2 and 3 respectively.

The stiffness matrix for truss element is given as,

$$
[K]=\frac{A E}{L}\left[\begin{array}{cccc}
C^{2} & C S & -C^{2} & -C S \\
C S & S^{2} & -C S & -S^{2} \\
-C^{2} & -C S & C^{2} & C S \\
-C S & -S^{2} & C S & S^{2}
\end{array}\right]
$$

Where, $C=\cos \theta$ and $S=\sin \theta$

Stiffness matrix for element (1),

$$
\left[K_{1}\right]=\frac{A_{1} E_{1}}{L_{1}}\left[\begin{array}{cccc}
C^{2} & C S & -C^{2} & -C S \\
C S & S^{2} & -C S & -S^{2} \\
-C^{2} & -C S & C^{2} & C S \\
-C S & -S^{2} & C S & S^{2}
\end{array}\right]
$$

Where,

$$
C=\cos \theta_{1}=\cos 0^{\circ}=1
$$

$\left(\because\right.$ Element $(1)$ is on $x$-axis, $\left.\theta_{1}=0\right)$

$$
S=\sin \theta_{1}=\sin 0^{\circ}=0
$$

Length, $L_{1}=500 \mathrm{~mm}$
Area, $A_{1}=A=200 \mathrm{~mm}^{2}$
Young's modulus, $E_{1}=E=70 \times 10^{3} \mathrm{~N} / \mathrm{mm}^{2}$ Then,

$$
\begin{aligned}
& \therefore \quad\left[K_{1}\right]=\frac{200 \times 70 \times 10^{3}}{500}\left[\begin{array}{cccc}
u_{1} & v_{1} & u_{2} & v_{2} \\
1 & 0 & -1 & 0 \\
0 & 0 & 0 & 0 \\
-1 & 0 & 1 & 0 \\
0 & 0 & 0 & 0
\end{array}\right] \begin{array}{c}
u_{1} \\
v_{1} \\
u_{2} \\
v_{2}
\end{array} \\
& u_{1} v_{1} \quad u_{2} v_{2} \\
& =28 \times 10^{3}\left[\begin{array}{cccc}
1 & 0 & -1 & 0 \\
0 & 0 & 0 & 0 \\
-1 & 0 & 1 & 0 \\
0 & 0 & 0 & 0
\end{array}\right] u_{v_{1}}^{u_{1}} v_{1}
\end{aligned}
$$

Stiffness matrix for element (2),

$$
\left[K_{2}\right]=\frac{A_{2} E_{2}}{L_{2}}\left[\begin{array}{cccc}
C^{2} & C S & -C^{2} & -C S \\
C S & S^{2} & -C S & -S^{2} \\
-C^{2} & -C S & C^{2} & C S \\
-C S & -S^{2} & C S & S^{2}
\end{array}\right]
$$



From figure,

$$
\begin{aligned}
\sin \theta_{2} & =\frac{300}{500} \\
\therefore \quad \theta_{2} & =\sin ^{-1}(0.6) \\
& =36.87^{\circ}
\end{aligned}
$$

Then,

$$
C=\cos \theta_{2}=\cos (-36.87)=0.8
$$

('-ve' sign due to clockwise consideration from positive $x$-axis)
$S=\sin \theta_{2}=\sin (-36.87)=-0.6$
Length, $l_{2}=\sqrt{300^{2}+400^{2}}=500 \mathrm{~mm}$
Area, $A_{2}=A=200 \mathrm{~mm}$
Young's modulus, $E_{2}=E=70 \times 10^{3} \mathrm{~N} / \mathrm{mm}^{2}$

$$
\begin{aligned}
\therefore \quad\left[K_{2}\right] & =\frac{200 \times 70 \times 10^{3}}{500}\left[\begin{array}{cccc}
u_{2} & v_{2} & u_{3} & v_{3} \\
0.8^{2} & (0.8 \times-0.6) & -0.8^{2} & -(0.8 \times-0.6) \\
(0.8 \times-0.6) & (-0.6)^{2} & -(0.8 \times-0.6) & -(-0.6)^{2} \\
-0.8^{2} & -(0.8 \times-0.6) & 0.8^{2} & (0.8 \times-0.6) \\
-(0.8 \times-0.6) & -(-0.6)^{2} & (0.8 \times-0.6) & (-0.6)^{2}
\end{array}\right] \begin{array}{l}
u_{2} \\
v_{2} \\
u_{3} \\
v_{3}
\end{array} \\
& \left.=28 \times 10^{3}\left[\begin{array}{cccc}
0.64 & v_{2} & u_{3} & v_{3} \\
-0.48 & -0.64 & 0.48 \\
-0.64 & 0.48 & 0.48 & -0.36 \\
0.48 & -0.36 & -0.48 & -0.48 \\
0.36
\end{array}\right] \right\rvert\, \begin{array}{l}
u_{2} \\
v_{2} \\
u_{3} \\
v_{3}
\end{array}
\end{aligned}
$$

Global stiffness matrix, $[K]=\left[K_{1}\right]+\left[K_{2}\right]$
$\therefore \quad[K]=28 \times 10^{3}\left[\begin{array}{cccccc}u_{1} & v_{1} & u_{2} & v_{2} & u_{3} & v_{3} \\ \hline 1 & 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ -1 & 0 & 1+0.64 & 0-0.48 & -0.64 & 0.48 \\ 0 & 0 & 0-0.48 & 0+0.36 & 0.48 & -0.36 \\ 0 & 0 & -0.64 & 0.48 & 0.64 & -0.48 \\ 0 & 0 & 0.48 & -0.36 & -0.48 & 0.36\end{array}\right] \begin{aligned} & u_{1} \\ & v_{1} \\ & u_{2} \\ & v_{2} \\ & u_{3} \\ & v_{3}\end{aligned}$

The finite element equation can be written as,

$$
\begin{aligned}
& {[K]\left[\delta^{\prime}\right]=\{F\}} \\
& 28 \times 10^{3}\left[\begin{array}{cccccc}
1 & 0 & -1 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 \\
-1 & 0 & 1.64 & -0.48 & -0.64 & 0.48 \\
0 & 0 & -0.48 & 0.36 & 0.48 & -0.36 \\
0 & 0 & -0.64 & 0.48 & 0.64 & -0.48 \\
0 & 0 & 0.48 & -0.36 & -0.48 & 0.36
\end{array}\right]\left[\begin{array}{l}
u_{1} \\
v_{1} \\
u_{2} \\
v_{2} \\
u_{3} \\
v_{3}
\end{array}\right]=\left[\begin{array}{l}
F_{1 x} \\
F_{1 y} \\
F_{2 x} \\
F_{2 y} \\
F_{3 x} \\
F_{3 y}
\end{array}\right\}
\end{aligned}
$$

## Nodal Displacements

From nodal boundary conditions,

$$
\begin{aligned}
& u_{1}=0 ; v_{1}=0 ; u_{3}=0 ; v_{3}=0 \\
& F_{2 x}=0 ; F_{2 y}=-12000 \mathrm{~N}
\end{aligned}
$$

Eliminating 1, 2, 5, 6 rows and columns,

$$
\begin{aligned}
& 28 \times 10^{3}\left[\begin{array}{cc}
1.64 & -0.48 \\
-0.48 & 0.36
\end{array}\right]\left\{\begin{array}{l}
u_{2} \\
v_{2}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
-12000
\end{array}\right\} \\
& 1.64 \times 28 \times 10^{3} \times u_{2}-0.48 \times 28 \times 10^{3} \times v_{2}=0 \\
& -0.48 \times 28 \times 10^{3} \times u_{2}+0.36 \times 28 \times 10^{3} \times v_{2}=-12000
\end{aligned}
$$

Solving equations (2) and (3),

$$
\begin{aligned}
& u_{2}=-0.571 \mathrm{~mm} \\
& v_{2}=-1.95 \mathrm{~mm}
\end{aligned}
$$

$\therefore$ The displacement at node 2 along ' $X$ ' and ' $Y$ ' directions (i..e, $u_{2}$ and $v_{2}$ ) are -0.571 mm and -1.95 mm respectively.

Stress Induced in Element (2)]

$$
\left.\begin{array}{rl}
\sigma_{2} & =\frac{E}{L_{2}}\left[\begin{array}{llll}
-C & -S & C & S
\end{array}\right]\left[\begin{array}{l}
u_{2} \\
v_{2} \\
u_{3} \\
v_{3}
\end{array}\right] \\
& =\frac{70 \times 10^{3}}{500}\left[\begin{array}{llll}
-0.8 & 0.6 & 0.8 & -0.6
\end{array}\right]\left\{\begin{array}{c}
-0.571 \\
-1.95 \\
0 \\
0
\end{array}\right] \\
& =140[0.4568-1.17+0-0]
\end{array}\right\} \quad \sigma_{2}=-99.848 \mathrm{~N} / \mathrm{mm}^{2} \quad 4 .
$$

$\therefore$ Stress at element ' $1-3$ ' is $-99.848 \mathrm{~N} / \mathrm{mm}^{2}$.
4.3. For the pin-jointed configuration shown in Fig. P4.3, determine the stiffness values $K_{11}, K_{12}$, and $K_{22}$ of the global stiffness matrix.


FIGURE P4.3


1. For the truss shown in the figure, a horizontal load $P(N)$ is applied in the $\mathbf{x}$ direction at node 2.
(a) Write down the element stiffness matrix $\mathbf{k}$ for each element.
(b) Assemble the $\mathbf{K}$ matrix.
(c) Using the elimination approach, solve for $\mathbf{Q}$.
(d) Evaluate the stress in elements 2 and 3.
(e) Determine the reaction force at node 2 in the $y$ direction.
4.4. For the truss in Fig. P4.4, a horizontal load of $P=2500 \mathrm{lb}$ is applied in the $x$ direction at node 2.

(a) Write down the element stiffness matrix $\mathbf{k}$ for each element.
(b) Assemble the $\mathbf{K}$ matrix.
(c) Using the elimination approach, solve for $\mathbf{Q}$.
(d) Evaluate the stress in elements 2 and 3.
(e) Determine the reaction force at node 2 in the $y$ direction.

$$
\mathbf{k}^{(1)}=\frac{30 \times 10^{6} \times 1.5}{30}\left[\begin{array}{lrrr}
1 & 2 & 3 & 4 \\
0 & 0 & -1 & 0 \\
-1 & 0 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0
\end{array}\right]
$$

$$
\begin{aligned}
& \mathrm{E}=30 \times 10^{6} p s i \\
& A=1.5 \mathrm{in}^{2}
\end{aligned}
$$

$$
\mathbf{k}^{(2)}=\frac{30 \times 10^{6} \times 1.5}{40}\left[\begin{array}{lrll}
3 & 4 & 5 & 6 \\
0 & 0 & 0 & 0 \\
0 & 1 & 0 & -1 \\
0 & 0 & 0 & 0 \\
0 & -1 & 0 & 1
\end{array}\right]
$$

$$
\mathbf{k}^{(3)}=\frac{30 \times 10^{6} \times 1.5}{50}\left[\begin{array}{cccc}
7 & 8 & 3 & 4 \\
.36 & -.48 & -.36 & .48 \\
-.48 & .64 & .48 & -.64 \\
-.36 & .48 & .36 & -.48 \\
.48 & -.64 & -.48 & .64
\end{array}\right]
$$

$$
\mathbf{k}^{(4)}=\frac{30 \times 10^{6} \times 1.5}{30}\left[\begin{array}{llrl}
1 & 8 & 5 & 6 \\
0 & 0 & -1 & 0 \\
-1 & 0 & 0 & 0 \\
0 & 0 & 1 & 0 \\
& & & 0
\end{array}\right]
$$

$$
\mathbf{K}=\frac{30 \times 10^{6} \times 1.5}{600}\left[\begin{array}{cccccccc}
1 & 2 & 3 & 4 & 5 & 6 & 7 & 8 \\
20 & 0 & -20 & 0 & 0 & 0 & 0 & 0 \\
& 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
& 24.32 & -5.76 & 0 & 0 & -4.32 & 5.76 \\
& & 22.68 & 0 & -15 & 0 & -7.68 \\
& & & & 20 & 0 & -20 & 0 \\
& & & & 15 & 0 & 0 \\
& \text { Sym } & & & & 24.32 & -5.76 \\
& & & & & & 7.68
\end{array}\right]
$$

Solution of $\mathbf{K} \mathbf{Q}=\mathbf{F}$

$$
\left\{\begin{array}{l}
Q_{1} \\
Q_{2} \\
Q_{3} \\
Q_{4} \\
Q_{5} \\
Q_{6} \\
Q_{7} \\
Q_{8}
\end{array}\right\}
$$

(c) Eliminating dof's 1, 2, 4, 7,,

$$
\mathbf{K} \mathbf{Q}=\mathbf{F} \text { is }
$$

3
$5 \quad 6$

$$
\frac{30 \times 10^{6} \times 1.5}{600}\left[\begin{array}{ccc}
24.32 & 0 & 0 \\
0 & 20 & 0 \\
0 & 0 & 15
\end{array}\right]\left\{\begin{array}{l}
Q_{3} \\
Q_{5} \\
Q_{6}
\end{array}\right\}=\left\{\begin{array}{l}
4000 \\
0 \\
0
\end{array}\right\}
$$

The solution is

$$
\begin{aligned}
& Q_{3}=219.3 \times 10^{-5} i n \\
& Q_{5}=0 \\
& Q_{6}=0
\end{aligned}
$$

(d)

$$
\begin{aligned}
& \sigma=\frac{E}{l_{e}}\left[\begin{array}{llll}
-\ell & -m & \ell & m
\end{array}\right] \mathbf{q} \\
& \sigma_{2}=\frac{30 \times 10^{6}}{40}\left[\begin{array}{llll}
0 & -1 & 0 & 1
\end{array}\right]\left\{\begin{array}{l}
219.3 \times 10^{-5} \\
0 \\
0 \\
0
\end{array}\right\}=0
\end{aligned}
$$

(Check with free body of Node 3)

$$
\sigma_{3}=\frac{30 \times 10^{6}}{50}\left[\begin{array}{llll}
-.6 & .8 & +.6 & -.8
\end{array}\right]\left\{\begin{array}{l}
0 \\
0 \\
219.3 \times 10^{-5} \\
0
\end{array}\right\}=789.5 p s i
$$

$\mathbf{K}=\frac{30 \times 10^{6} \times 1.5}{600}\left[\begin{array}{cccccccc}1 & 2 & 3 & 4 & 5 & 6 & 7 & 8 \\ 20 & 0 & -20 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ & 24.32 & -5.76 & 0 & 0 & -4.32 & 5.76 \\ & & 22.68 & 0 & -15 & 0 & -7.68 \\ \hline & & & 20 & 0 & -20 & 0 \\ \\ & & & & 15 & 0 & 0 \\ \text { Sym } & & & & 24.32 & -5.76 \\ \hline\end{array}\right.$

$$
\begin{aligned}
R_{4} & =\sum_{j=1}^{8} K_{4 j} Q_{j} \\
& \left.=\frac{30 \times 10^{6} \times 1.5}{600}\left[-5.76 \times 219.3 \times 10^{-5}\right]=-947.4 l b \quad \text { (downward pull }\right)
\end{aligned}
$$

## Unit 2

## Beams

| Bar Element |  | Beam Element |  |
| :---: | :---: | :---: | :---: |
| 1. | Bar is a structural element that is subjected to only axial loading. | 1. | Beam is a structural element that is subjected to transverse loading |
| 2. | When a bar element is loaded, it is described by the --axial-- displacements only. | 2. | When a beam element is loaded, it is described by the transverse displacements and rotational (slope) displacements. |
| 3. | Stiffness matrix for a bar element is given by, $[K]=\frac{A E}{L}\left[\begin{array}{cc} 1 & -1 \\ -1 & 1 \end{array}\right]$ | 3. | Stiffness matrix for a beam element is given by, $[K]=\frac{E I}{L^{3}}\left[\begin{array}{cccc} 12 & 6 L & -12 & 6 L \\ 6 L & 4 L^{2} & -6 L & 2 L^{2} \\ -12 & -6 L & 12 & -6 L \\ 6 L & 2 L^{2} & -6 L & 4 L^{2} \end{array}\right]$ |
| 4. | Bar elements are used in model cables, prismatic structural members and ropes. | 4. | Long horizontal members used in buildings, bridges and shafts are some of the examples of beams. |

## Assumptions used in beam elements are:

1. The beam elements are straight and prismatic.
2. The material of beam is linearly-elastic, isotropic and homogeneous.
3. The cross-section of beam is either constant or varies smoothly. Deformation of cross-section does not occur in its plane, but subjected to warping in longitudinal direction.
4. The transverse shear and axial force effects are assumed to be negligible. In case of bending moment deformation, internal strain energy of the beam element is considered.
5. The resultant of stresses (i.e.,internal moments) are determined by Euler-Bernoulli theories for bending stress and Timoshenko theory for torsional stress.
6. The beam elements have larger displacements and smaller strains.
7. External load applied on a beam is static and conservative.

## Derivation of shape functions for beam element



## Figure: Beam Element

The above figure shows a beam of length ' $l$ ' upon loading the beam will have four displacements. $v_{1} v_{2}$ are the Transverse displacements and $\theta_{1}, \theta_{2}$ are rotational displacement.

The polynomial function of a beam element of two nodes and with four displacements is given by,

$$
\begin{equation*}
v(x)=a_{1}+a_{2} x+a_{3} x^{2}+a_{4} x^{3} \tag{1}
\end{equation*}
$$


$q=\left[q_{1}, q_{2}, q_{3}, q_{4}\right]^{T}$
$=\left[v_{1}, v_{1}^{\prime}, v_{2}, v_{2}^{\prime}\right]$
$a_{1}, a_{2}, a_{3}, a_{4}-$ polynomial coefficients

$$
\begin{equation*}
v(x)=a_{1}+a_{2} x+a_{3} x^{2}+a_{4} x^{3} \tag{1}
\end{equation*}
$$

Upon derivating equation (1),

$$
\begin{equation*}
\frac{d v}{d x}=a_{2}+2 a_{3} x+3 a_{4} x^{2} \tag{2}
\end{equation*}
$$

Applying boundary conditions to equation (2),
At

$$
x=0, v=v_{1}, \frac{d v}{d x}=\theta_{1}
$$

At

$$
x=l, v=v_{2}, \frac{d v}{d x}=\theta_{2}
$$

From equation (1) and (2),

$$
\begin{align*}
& v_{1}=a_{1}  \tag{3}\\
& \theta_{1}=a_{2}  \tag{4}\\
& v_{2}=a_{1}+a_{2} l+a_{3} l^{2}+a_{4} l^{3}  \tag{5}\\
& \theta_{2}=a_{2}+2 a_{3} l+3 a_{4} l^{2} \tag{6}
\end{align*}
$$

From equation (5), (4), (3)

$$
\begin{align*}
v_{2}= & a_{1}+a_{2} l+a_{3} l^{2}+a_{4} l^{3} \\
& v_{2}=v_{1}+\theta_{1} l+a_{3} l^{2}+a_{4} l^{3} \\
& a_{3} l^{2}+a_{4} l^{3}=v_{2}-v_{1}-\theta_{1} l \tag{7}
\end{align*}
$$

From equation (6), (4), (3)

$$
\begin{align*}
\theta_{2}= & a_{2}+2 a_{3} l+3 a_{4} l^{2} \\
& \theta_{2}=\theta_{1}+2 a_{3} l+3 a_{4} l^{2} \\
& 2 a_{3} l+3 a_{4} l^{2}=\theta_{2}-\theta_{1} \tag{8}
\end{align*}
$$

Multiplay equation (7) by ' 3 ' and equation (8) by ' $l$ '

$$
\begin{aligned}
& 3 a_{3} l^{2}+3 a_{4} l^{3}=3 v_{2}-3 v_{1}-3 \theta_{1} l \\
& 2 \mathrm{a}_{3} l^{2}+3 \mathrm{a}_{4} l^{3}=\theta_{2} l-\theta_{1} l
\end{aligned}
$$

$$
\begin{gathered}
a_{3} l^{2}=3 v_{2}-3 v_{1}-3 \theta_{1} l-\theta_{2} l+\theta_{1} l \\
a_{3} l^{2}=3 v_{2}-3 v_{1}-2 \theta_{1} l-\theta_{2} l \\
a_{3}=\frac{3}{l^{2}}\left(v_{2}-v_{1}\right)-\frac{1}{l}\left(2 \theta_{1}+\theta_{2}\right)
\end{gathered}
$$

$$
\text { Substitute the ' } a_{3} \text { ' value in equation (7), }
$$

$$
\begin{aligned}
a_{4} l^{3} & =v_{2}-v_{1}-\theta_{1} l-a_{3} l^{2} \\
a_{4} l^{3} & =v_{2}-v_{1}-\theta_{1} l-\left[\frac{3}{l^{2}}\left(v_{2}-v_{1}\right)-\frac{1}{l}\left(2 \theta_{1}+\theta_{2}\right)\right] l^{2} \\
& =v_{1}-v_{1}-\theta_{1} l-3\left(v_{2}-v_{1}\right)+l\left(2 \theta_{1}+\theta_{2}\right) \\
& =v_{2}-v_{1}-\theta_{1} l-3 v_{2}+3 v_{1}+2 l \theta_{1}+\theta_{2} l \\
a_{4} l^{3} & =-2 v_{2}+2 v_{1}+\theta_{1} l+\theta_{2} l \\
a_{4} & =-\frac{2}{l^{3}} v_{2}+\frac{2}{l^{3}} v_{1}+\frac{\theta_{1}}{l^{2}}+\frac{\theta_{2}}{l^{2}}
\end{aligned}
$$

$$
\begin{aligned}
v(x) & =a_{1}+a_{2} x+a_{3} x^{2}+a_{4} x^{3} \\
v(x) & =v_{1}+\theta_{1} x+\left[\frac{3}{l^{3}}\left(v_{2}-v_{1}\right)-\frac{1}{l}\left(2 \theta_{1}+\theta_{2}\right)\right] x^{2}+\left[\frac{-2}{l^{3}} v_{2}+\frac{2}{l^{3}} v_{1}+\frac{\theta_{1}}{l^{2}}+\frac{\theta_{2}}{l^{2}}\right] x^{3} \\
v(x) & =v_{1}+\theta_{1} x+\frac{3}{l^{2}} v_{2} x^{2}-\frac{3}{l^{2}} v_{1} x^{2}-\frac{1}{l} 2 \theta_{1} x^{2}-\frac{1}{l} \theta_{2} x^{2}-\frac{2}{l^{3}} v_{2} x^{3}+\frac{2}{l^{3}} v_{1} x^{3}+\frac{\theta_{1}}{l^{2}} x^{3}+\frac{\theta_{2}}{l^{2}} \cdot x^{3} \\
& =v_{1}\left[1-\frac{3}{l^{2}} x^{2}+\frac{2}{l^{3}} x^{3}\right]+\theta_{2}\left[x-\frac{2}{l} x^{2}+\frac{x^{3}}{l^{2}}\right]+v_{2}\left[\frac{3}{l^{2}} x^{2}-\frac{2 x^{3}}{l^{3}}\right]+\theta_{2}\left[-\frac{x^{2}}{l}+\frac{x^{3}}{l^{2}}\right] \\
v & =N_{1} v_{1}+N_{2} \theta_{1}+N_{3} v_{2}+N_{4} \theta_{2}(\text { or }) N_{1} \delta_{1}+N_{2} \delta_{2}+N_{3} \delta_{3}+N_{4} \delta_{4}
\end{aligned}
$$

$N_{1}, N_{2}, N_{3}, N_{4}$ - Shape functions of a beam element
$v_{1}, \theta_{1} v_{2}, \theta_{2}$ - Nodal displacements of a beam element

$$
v=N_{1} v_{1}+N_{2} \theta_{1}+N_{3} v_{2}+N_{4} \theta_{2} \text { (or) } N_{1} \delta_{1}+N_{2} \delta_{2}+N_{3} \delta_{3}+N_{4} \delta_{4}
$$

$N_{1}, N_{2}, N_{3}, N_{4}$ - Shape functions of a beam element
$v_{1}, \theta_{1} v_{2}, \theta_{2}$-Nodal displacements of a beam element

Where,

$$
\begin{aligned}
& N_{1}=1-\frac{3 x^{3}}{l^{2}}+\frac{2 x^{2}}{l^{3}} \\
& N_{2}=x-\frac{2 x^{2}}{l}+\frac{x^{3}}{l^{2}} \\
& N_{3}=\frac{3 x^{2}}{l^{2}}-\frac{2 x^{3}}{l^{3}} \\
& N_{4}=\frac{-x^{2}}{l}+\frac{x^{3}}{l^{2}}
\end{aligned}
$$



$$
\mathrm{N}_{4}(K)-1-3\left(\frac{x}{L}\right)^{2}+2\left(\frac{x}{L}\right)
$$

$$
\mathrm{N}_{2}(X)-x\left(1-\frac{x}{L}\right)^{2}
$$


$\mathrm{N}_{3}(x)=3\left(\frac{x}{L}\right)^{2}-2\left(\frac{x}{L}\right)^{3}$
$\mathrm{N}_{4}(x)=\frac{x^{2}}{L}\left(\frac{x}{L}-1\right)$

## Element Stiffness Matrix ( K )_4by4

Strain energy in an element of length dx is

$$
\begin{aligned}
& \begin{aligned}
d U & =\frac{1}{2} \int_{A} \sigma \varepsilon d A d x \\
& =\frac{1}{2}\left(\frac{M^{2}}{E I^{2}} \int_{A}^{2} y^{2} d A\right) d x
\end{aligned} \\
& \int_{A} y^{2} d A \text { is the moment of inertia I }
\end{aligned}
$$

The total strain energy for the beam is given by-

$$
U=\frac{1}{2} \int_{0}^{L} E I\left(d^{2} v / d x^{2}\right) d x
$$

$$
d^{2} v / d x^{2}=M / E I
$$

Differentiating twice on both sides,
$v$ - Beam deflection


Figure: Beam Element

$$
U=\frac{E I}{2} \int_{0}^{l}\left[\frac{d^{2} v}{d x^{2}}\right]^{2} \cdot d x
$$

From nodal displacement equation,

$$
\begin{aligned}
& v=N_{1} v_{1}+N_{2} \theta_{1}+N_{3} v_{2}+N_{4} \theta_{4} \\
& \quad \text { (or) } \\
& v=N_{1} \delta_{1}+N_{2} \delta_{2}+N_{3} \delta_{3}+N_{4} \delta_{4}
\end{aligned}
$$

$$
\frac{d^{2} v}{d x^{2}}=\frac{d^{2} N_{1}}{d x^{2}} \delta_{1}+\frac{d^{2} N_{2}}{d x^{2}} \delta_{2}+\frac{d^{2} N_{3}}{d x^{2}} \delta_{3}+\frac{d^{2} N_{4}}{d x^{2}} \delta_{4}
$$

Let,

$$
\begin{aligned}
& B_{1}=\frac{d^{2} N_{1}}{d x^{2}} \\
& B_{2}=\frac{d^{2} N_{2}}{d x^{2}} \\
& B_{3}=\frac{d^{2} N_{3}}{d x^{2}} \\
& B_{4}=\frac{d^{2} N_{4}}{d x^{2}}
\end{aligned}
$$

Where,

$$
\begin{aligned}
& N_{1}=1-\frac{3 x^{3}}{l^{2}}+\frac{2 x^{2}}{l^{3}} \\
& N_{2}=x-\frac{2 x^{2}}{l}+\frac{x^{3}}{l^{2}} \\
& N_{3}=\frac{3 x^{2}}{l^{2}}-\frac{2 x^{3}}{l^{3}} \\
& N_{4}=\frac{-x^{2}}{l}+\frac{x^{3}}{l^{2}}
\end{aligned}
$$

$$
\frac{d^{2} v}{d x^{2}}=B_{1} \delta_{1}+B_{2} \delta_{2}+B_{3} \delta_{3}+B_{4} \delta_{4}
$$

$$
\frac{d^{2} v}{d x^{2}}=B_{1} \delta_{1}+B_{2} \delta_{2}+B_{3} \delta_{3}+B_{4} \delta_{4}
$$

$$
U=\frac{E I}{2} \int_{0}^{l}\{\delta\}^{T}[B]^{T}[B]\{\delta\} d x
$$

Expressing the above equation in matrix form,

$$
\begin{aligned}
& \frac{d^{2} v}{d x^{2}}=\left[\begin{array}{llll}
B_{1} & B_{2} & B_{3} & B_{4}
\end{array}\right]\left\{\begin{array}{l}
\delta_{1} \\
\delta_{2} \\
\delta_{3} \\
\delta_{4}
\end{array}\right\} \\
& \frac{d^{2} v}{d x^{2}}=[B]\{\delta\}
\end{aligned}
$$

Squaring on both sides,

$$
\begin{aligned}
& {\left[\frac{d^{2} v}{d x^{2}}\right]^{2}=[[B]\{\delta\}]^{2}} \\
& {\left[\frac{d^{2} v}{d x^{2}}\right]^{2}=\{\delta\}^{\mathrm{T}}[B]^{T}[B]\{\delta\}}
\end{aligned}
$$

$$
\begin{aligned}
U & =\frac{1}{2}\{\delta\}^{T}\left[E I \int_{0}^{l}[B]^{T}[B] \cdot d x\right]\{\delta\} \\
U & =\frac{1}{2}\{\delta\}^{T}[k]\{\delta\}
\end{aligned}
$$

$$
U=\frac{1}{2}\{\delta\}^{T}[k]\{\delta\}
$$

Where,

$$
\begin{aligned}
& {[k]=E I \int_{0}^{l}[B]^{T}[B] . d x} \\
& {[k]=E I \int_{0}^{t}\left[\begin{array}{l}
B_{1} \\
B_{2} \\
B_{3} \\
B_{4}
\end{array}\right]\left[\begin{array}{llll}
B_{1} & B_{2} & B_{3} & B_{4}
\end{array}\right] d x} \\
& {[k]=E I \int_{0}^{l}\left[\begin{array}{cccc}
B_{1}^{2} & B_{1} B_{2} & B_{1} B_{3} & B_{1} B_{4} \\
B_{1} B_{2} & B_{2}^{2} & B_{2} B_{3} & B_{2} B_{4} \\
B_{1} B_{3} & B_{2} B_{3} & B_{3}^{2} & B_{3} B_{4} \\
B_{1} B_{4} & B_{2} B_{4} & B_{3} B_{4} & B_{4}^{2}
\end{array}\right] d x}
\end{aligned}
$$

From shape functions,

$$
\begin{aligned}
B_{1} & =\frac{d^{2} N_{1}}{d x^{2}} \\
N_{1} & =l-\frac{3 x^{2}}{l^{2}}+\frac{2 x^{3}}{l^{3}} \\
\frac{d N_{1}}{d x} & =\frac{-6 x}{l^{2}}+\frac{6 x^{2}}{l^{3}} \\
\frac{d^{2} N}{d x^{2}} & =\frac{-6}{l^{2}}+\frac{12 x}{l^{3}} \\
\therefore B_{1} & =\frac{-6}{l^{2}}+\frac{12 x}{l^{3}}
\end{aligned}
$$

Similarly,
For,

$$
\begin{aligned}
& N_{2}=x-\frac{2 x^{2}}{l}+\frac{x^{3}}{l^{2}} \\
& B_{2}=\frac{d^{2} N_{2}}{d x^{2}}=\frac{-4}{l}+\frac{6 x}{l^{2}}
\end{aligned}
$$

Now,

$$
\begin{aligned}
\int_{0}^{l} B_{1}^{2} \cdot d x & =\int_{0}^{l}\left(\frac{-6}{l^{2}}+\frac{12 x}{l^{3}}\right)^{2} \cdot d x \\
& =\int_{0}^{l}\left(\frac{36}{l^{4}}+\frac{144 x^{2}}{l^{6}}-\frac{144 x}{l^{5}}\right) \cdot d x \\
& =\left(\frac{36 x}{l^{4}}+\frac{144 x^{3}}{3 l^{6}}-\frac{144 x^{2}}{2 l^{5}}\right)_{0}^{l} \\
& =\frac{36 l}{l^{4}}+\frac{144 l^{3}}{3 l^{6}}-\frac{72 l^{2}}{l^{5}} \\
& =\frac{36}{l^{3}}+\frac{48}{l^{3}}-\frac{72}{l^{3}} \\
\int_{0}^{l} B_{1}^{2} \cdot d x & =\frac{12}{l^{3}}
\end{aligned}
$$

Similarly solve for all the values ' $B$ ' in equation (3), i.e., $B_{1} B_{2}, B_{1} B_{3}, B_{1} B_{4} \ldots B_{3} B_{4}, B_{4}^{2}$

The end matrix will be in the following form,

$$
\left[\begin{array}{llll}
B_{1}^{2} & B_{1} B_{2} & B_{1} B_{3} & B_{1} B_{4} \\
B_{1} B_{2} & B_{2}^{2} & B_{2} B_{3} & B_{2} B_{4} \\
B_{1} B_{3} & B_{2} B_{3} & B_{3}^{2} & B_{3} B_{4} \\
B_{1} B_{4} & B_{2} B_{4} & B_{3} B_{4} & B_{4}^{2}
\end{array}\right]=\left[\begin{array}{cccc}
\frac{12}{l^{3}} & \frac{6}{l^{2}} & \frac{-12}{l^{3}} & \frac{6}{l^{2}} \\
\frac{6}{l^{2}} & \frac{4}{l} & \frac{-6}{l^{2}} & \frac{2}{l} \\
\frac{-12}{l^{3}} & \frac{-6}{l^{2}} & \frac{12}{l^{3}} & \frac{-6}{l^{2}} \\
\frac{6}{l^{2}} & \frac{2}{l} & \frac{-6}{l^{2}} & \frac{4}{l}
\end{array}\right]=\frac{1}{l^{3}}\left[\begin{array}{cccc}
12 & 6 l & -12 & 6 l \\
6 l & 4 l^{2} & -6 l & 2 l^{2} \\
-12 & -6 l & 12 & -6 l \\
6 l & 2 l^{2} & -6 l & 4 l^{2}
\end{array}\right]
$$

We have

$$
[k]=E I \int_{0}^{l}\left[\begin{array}{cccc}
B_{1}^{2} & B_{1} B_{2} & B_{1} B_{3} & B_{1} B_{4} \\
B_{1} B_{2} & B_{2}^{2} & B_{2} B_{3} & B_{2} B_{4} \\
B_{1} B_{3} & B_{2} B_{3} & B_{3}^{2} & B_{3} B_{4} \\
B_{1} B_{4} & B_{2} B_{4} & B_{3} B_{4} & B_{4}^{2}
\end{array}\right] d x
$$

$$
[k]=\frac{E I}{l^{3}}\left[\begin{array}{cccc}
12 & 6 l & -12 & 6 l \\
6 l & 4 l^{2} & -6 l & 2 l^{2} \\
-12 & -6 l & 12 & -6 l \\
6 l & 2 l^{2} & -6 l & 4 l^{2}
\end{array}\right]
$$

## Force Vector



## Figure: Beam Element

$F_{1}, F_{2}$ - Shear forces in upward direction at nodes 1,2 .
$M_{1}, M_{2}$-Bending Moments in counterclockwise at node
1,2 .


Force vector for a beam element is given by,

$$
\begin{aligned}
&\{F\}=[K][\delta] \\
&\{F\}-\text { Force vector } \\
&=\left\{\begin{array}{l}
F_{1} \\
M_{1} \\
F_{2} \\
M_{2}
\end{array}\right\}
\end{aligned}
$$

$\{\delta\}$ - Nodal displacement vector

$$
=\left\{\begin{array}{l}
v_{1} \\
\theta_{1} \\
v_{2} \\
\theta_{2}
\end{array}\right\}
$$

Figure: Beam Element


Figure (2): FE Model

$$
\left\{\begin{array}{l}
F_{1} \\
M_{1} \\
F_{2} \\
M_{2}
\end{array}\right\}=\frac{E I}{l^{3}}\left[\begin{array}{cccc}
12 & 6 l & -12 & 6 l \\
6 l & 4 l^{2} & -6 l & 2 l^{2} \\
-12 & -6 l & 12 & -6 l \\
6 l & 2 l^{2} & -6 l & 4 l^{2}
\end{array}\right]\left\{\begin{array}{l}
v_{1} \\
\theta_{1} \\
v_{2} \\
\theta_{2}
\end{array}\right\}
$$

$$
\begin{aligned}
& \text { Simple Support } \\
& u=v=0 \\
& \text { unn }=0
\end{aligned}
$$

Roller Support
$\mathrm{v}=0$


## Internal Support <br> $$
\mathrm{v}=0
$$




Fixed Support


Guided Support


Internal Hinge
$\mathrm{M}=0$
мй

Figure 4.12. Typical beam boundary conditions


Deflection: $v$
Rotation: $\quad \theta=\frac{\partial v}{\partial x}$
Bending moment: $\quad M=E I \frac{\partial^{2} v}{\partial x^{2}}$
Shear force: $\quad V=\frac{\partial}{\partial x}\left[E I \frac{\partial^{2} v}{\partial x^{2}}\right]$

For the cantilever beam shown in the figure determine the nodal displacements. Construct the shear force and bending moment diagrams. Compare the results. Given $E=210 \mathrm{GPa}$ and $\mathrm{I}=5000 \mathrm{~cm}^{4}$.


Given that,
Young's modulus, $E=210 \mathrm{GPa}=210 \times 10^{9} \mathrm{~N} / \mathrm{m}^{2}$
Moment of inertia, $I=5000 \mathrm{~cm}^{4}=5000 \times 10^{-8} \mathrm{~m}^{4}$


Figure (1): Discretization of Beam

Let $v_{1}, \theta_{1}, v_{2}, \theta_{2}, v_{3}, \theta_{2}$ are the nodal displacements.
The element stiffness matrix for beam element is given by,

$$
[K]=\frac{E I}{l^{3}}\left[\begin{array}{cccc}
12 & 6 l & -12 & 6 l \\
6 l & 4 l^{2} & -6 l & 2 l^{2} \\
-12 & -6 l & 12 & -6 l \\
6 l & 2 l^{2} & -6 l & 4 l^{2}
\end{array}\right]
$$

## Element-1

Stiffiness matrix for element (1) is given by,

$$
\left[\mathrm{K}_{1}\right]=\frac{E_{1} I}{l_{1}^{3}}\left[\begin{array}{cccc}
12 & 6 l_{1} & -12 & 6 l_{1} \\
6 l_{1} & 4 l^{2} & -6 l_{1} & 2 l^{2} \\
-12 & -6 l_{1} & 12 & -6 l_{1} \\
6 l_{1} & 2 l_{1}^{2} & -6 l_{1} & 4 l_{1}^{2}
\end{array}\right]
$$

Where, $I_{1}=2 I$

$$
\begin{aligned}
{\left[K_{1}\right] } & =\frac{210 \times 10^{9} \times 5000 \times 10^{-8} \times 2}{(1.5)^{3}}\left[\begin{array}{cccc}
12 & 6(1.5) & -12 & 6(1.5) \\
6(1.5) & 4(1.5)^{2} & -6(1.5) & 2(1.5)^{2} \\
-12 & -6(1.5) & 12 & -6(1.5) \\
6(1.5) & 2(1.5)^{2} & -6(1.5) & 4(1.5)^{2}
\end{array}\right] \\
& =6.22 \times 10^{6} \cdot\left[\begin{array}{cccc}
12 & 9 & -12 & 9 \\
9 & 9 & -9 & 4.5 \\
-12 & -9 & 12 & -9 \\
9 & 4.5 & -9 & 9
\end{array}\right] \\
& =0.622 \times 10^{7} \cdot\left[\begin{array}{cccc}
12 & 9 & -12 & 9 \\
9 & 9 & -9 & 4.5 \\
-12 & -9 & 12 & -9 \\
9 & 4.5 & -9 & 9
\end{array}\right] \quad\left[K_{1}\right]=10^{7} \cdot\left[\begin{array}{ccccc}
v_{1} & \theta_{1} & v_{2} & \theta_{2} \\
7.464 & 5.598 & -7.464 & 5.598 \\
5.598 & 5.598 & -5.598 & 2.799 \\
-7.464 & -5.598 & 7.464 & -5.598 \\
5.598 & 2.799 & -5.598 & 5.598
\end{array}\right] v_{1}^{v_{1}} \theta_{1}
\end{aligned}
$$

Element-2
Stiffness matrix for element (2) is given by,

$$
\left[K_{2}\right]=\frac{E_{2} I}{l_{2}^{3}}\left[\begin{array}{cccc}
12 & 6 l_{2} & -12 & 6 l_{2} \\
6 l_{2} & 4 l_{2}^{2} & -6 l_{2} & 2 l_{2}^{2} \\
-12 & -6 l_{2} & 12 & -6 l_{2} \\
6 l_{2} & 2 l_{2}^{2} & 6 l_{2} & 4 l_{2}^{2}
\end{array}\right]
$$

Where, $I_{2}=I$

$$
\begin{array}{r}
{\left[K_{2}\right]=\frac{210 \times 10^{9} \times 5000 \times 10^{-8}}{(1.5)^{3}}\left[\begin{array}{cccc}
12 & 6(1.5) & -12 & 6(1.5) \\
6(1.5) & 4(1.5)^{2} & -6(1.5) & 2(1.5)^{2} \\
-12 & -6(1.5) & 12 & -6(1.5) \\
6(1.5) & 2(1.5)^{2} & -6(1.5) & 4(1.5)^{2}
\end{array}\right]=} \\
{\left[K_{2}\right]=10^{7} \cdot\left[\begin{array}{cccc}
3.732 & 2.799 & -3.732 & 2.799 \\
2.799 & 2.799 & -2.799 & 1.4 \\
-3.732 & -2.799 & 3.732 & -2.799 \\
2.799 & 1.4 & -2.799 & 2.799
\end{array}\right] \begin{array}{l}
v_{2} \\
\theta_{2} \\
v_{3} \\
\theta_{3}
\end{array}}
\end{array}
$$

$$
v_{1}
$$

$$
\left[\begin{array}{cccc}
7.464 & 5.598 & -7.464 & 5.598 \\
5.598 & 5.598 & -5.598 & 2.799 \\
-7.464 & -5.598 & \left.\begin{array}{cc}
7.464 & -5.598 \\
5.598 & 2.799
\end{array} \begin{array}{ccc}
-5.598 & 5.598
\end{array}\right] & v_{1} \\
\theta_{1} \\
v_{2} \\
\theta_{2}
\end{array} \quad\left[K_{2}\right]=10^{7} \cdot\left[\begin{array}{cccc}
\begin{array}{ccc}
3.732 & 2.799 \\
2.799 & 2.799
\end{array} & -3.732 & 2.799 \\
-2.799 & 1.4 \\
-3.732 & -2.799 & 3.732 & -2.799 \\
2.799 & 1.4 & -2.799 & 2.799
\end{array}\right]\left[\begin{array}{c}
v_{2} \\
\theta_{2} \\
v_{3} \\
\theta_{3}
\end{array}\right.\right.
$$

$\square$
$v_{2}$

Global stiffness matrix,

$$
\begin{gathered}
{[K]=\left[K_{1}\right]+\left[K_{2}\right]} \\
v_{1} \\
\theta_{1}
\end{gathered} v_{2} \quad \theta_{2} \quad v_{3} \quad \theta_{3} \begin{aligned}
& {[K]=10^{7}\left[\begin{array}{cccccc}
7.464 & 5.598 & -7.464 & 5.598 & 0 & 0 \\
5.598 & 5.598 & -5.598 & 2.799 & 0 & 0 \\
-7.464 & -5.598 & \begin{array}{ccc}
11.196 & -2.799 \\
5.598 & 2.799 & -2.799
\end{array} & -3.397 & -2.732 & 2.799 \\
0 & 0 & -3.732 & -2.799 & 3.732 & -2.799 \\
0 & 0 & 2.799 & 1.4 & -2.799 & 2.799
\end{array}\right] \begin{array}{l}
v_{1} \\
\theta_{1} \\
v_{2} \\
\theta_{2} \\
v_{3} \\
\theta_{3}
\end{array}}
\end{aligned}
$$

The finite element equation is given by,

$$
[K]\{\delta\}=\{F\}
$$

$$
10^{7}\left[\begin{array}{cccccc}
7.464 & 5.598 & -7.464 & 5.598 & 0 & 0 \\
5.598 & 5.598 & -5.598 & 2.799 & 0 & 0 \\
-7.464 & -5.598 & 11.196 & -2.799 & -3.732 & 2.799 \\
5.598 & 2.799 & -2.799 & 8.397 & -2.799 & 1.4 \\
0 & 0 & -3.732 & -2.799 & 3.732 & -2.799 \\
0 & 0 & 2.799 & 1.4 & -2.799 & 2.799
\end{array}\right]\left\{\begin{array}{c}
v_{1} \\
\theta_{1} \\
v_{2} \\
\theta_{2} \\
v_{3} \\
\theta_{3}
\end{array}\right]=\left\{\begin{array}{c}
F_{1} \\
M_{1} \\
F_{2} \\
M_{2} \\
F_{3} \\
M_{3}
\end{array}\right\}=\left\{\begin{array}{c}
F_{1} \\
M_{1} \\
-20 \times 10^{3} \\
M_{2} \\
-20 \times 10^{3} \\
M_{3}
\end{array}\right\}
$$

Applying boundary conditions,

$$
v_{1}=\theta_{1}=0
$$

Deleteing $1^{\text {st }}, 2^{\text {nd }}$ row and column from the above equation,

$$
10^{7}\left[\begin{array}{cccc}
11.196 & -2.799 & -3.732 & 2.799 \\
-2.799 & 8.397 & -2.799 & 1.4 \\
-3.732 & -2.799 & 3.732 & -2.799 \\
2.799 & 1.4 & -2.799 & 2.799
\end{array}\right]\left[\begin{array}{c}
v_{2} \\
\theta_{2} \\
v_{3} \\
\dot{\theta}_{3}
\end{array}\right]=\left\{\begin{array}{c}
-20 \times 10^{3} \\
0 \\
-20 \times 10^{3} \\
0
\end{array}\right\}
$$

On solving above equation,

$$
\left\{\begin{array}{l}
v_{2} \\
\theta_{2} \\
v_{3} \\
\theta_{4}^{3}
\end{array}\right\}=\left\{\begin{array}{c}
-0.00375 \mathrm{~m} \\
-0.00428 \mathrm{rad} \\
-0.01232 \mathrm{~m} \\
-0.00642 \mathrm{rad}
\end{array}\right\}
$$

$\therefore$ Nodal displacement vectors,

$$
\{\delta\}=\left\{\begin{array}{llllll}
0 & 0 & -0.00375 & -0.00428 & -0.01232 & -0.00642
\end{array}\right\}^{T}
$$

## Gauss Elimination Method

## $G$ to eliminate or remove variables from?

## System of Linear Equations

$\left.a_{1} x+b_{1} y+c_{1} z=d_{1}\right) \rightarrow$ Remiove $x$
$a_{2} x+b_{2} y+c_{2} z=d_{2} \rightarrow$ Reniove $a_{3} x+b_{3} y+c_{3} z=d_{3}$
Second Method Then find ${ }^{L}$
Form augmented matrix [ $\mathrm{A} \mid \mathrm{B}$ ]


Solution of Augmented Matrix

A is upper triangular matrix

$$
\begin{aligned}
& x+4 y-z=-5 \\
& x+y-6 z=-12 \\
& 3 x-y-z=4
\end{aligned}
$$

$$
\text { :We have }\left[\begin{array}{rrr}
1 & 4 & -1 \\
1 & 1 & -6 \\
3 & -1 & -1
\end{array}\right]\left[\begin{array}{l}
x \\
y \\
z
\end{array}\right]=\left[\begin{array}{c}
-5 \\
-12 \\
4
\end{array}\right]
$$

$$
\text { Operate } R_{2}-R_{1} \text { and } R_{3}-3 R_{1},\left[\begin{array}{rrr}
1 & 4 & -1 \\
0 & -3 & -5 \\
0 & -13 & 2
\end{array}\right]\left[\begin{array}{c}
x \\
y \\
z
\end{array}\right]=\left[\begin{array}{c}
-5 \\
-7 \\
19
\end{array}\right]
$$

$$
\text { Operate } R_{3}-\frac{13}{3} R_{2},\left[\begin{array}{rrr}
1 & 4 & -1 \\
0 & -3 & -5 \\
0 & 0 & 71 / 3
\end{array}\right]\left[\begin{array}{l}
x \\
y \\
z
\end{array}\right]=\left[\begin{array}{c}
-5 \\
-7 \\
148 / 3
\end{array}\right]
$$

Thus, we have $z=148 / 71=2.0845$,
$3 y=7-5 z=7-10.4225=-3.4225$ i.e., $y=-1.1408$
$x=-5-4 y+z=-5+4(1.1408)+2.0845=1.6479$
Hence $x=1.6479, y=-1.1408, z=2.0845$.

$$
\begin{aligned}
& x+4 y-z=-5 \\
& x+y-6 z=-12 \\
& 3 x-y-z=4
\end{aligned}
$$

$$
\text { : We have }\left[\begin{array}{rrr}
1 & 4 & -1 \\
1 & 1 & -6 \\
3 & -1 & -1
\end{array}\right]\left[\begin{array}{l}
x \\
y \\
z
\end{array}\right]=\left[\begin{array}{c}
-5 \\
-12 \\
4
\end{array}\right]
$$

Operate $R_{2}-R_{1}$ and $R_{3}-3 R_{1},\left[\begin{array}{rrr}1 & 4 & -1 \\ 0 & -3 & -5 \\ 0 & -13 & 2\end{array}\right]\left[\begin{array}{l}x \\ y \\ z\end{array}\right]=\left[\begin{array}{c}-5 \\ -7 \\ 19\end{array}\right]$
Operate $R_{3}-\frac{13}{3} R_{2},\left[\begin{array}{rrr}1 & 4 & -1 \\ 0 & -3 & -5 \\ 0 & 0 & 71 / 3\end{array}\right]\left[\begin{array}{l}x \\ y \\ z\end{array}\right]=\left[\begin{array}{c}-5 \\ -7 \\ 148 / 3\end{array}\right]$
Thus, we have $z=148 / 71=2.0845$,

$$
\begin{gathered}
3 y=7-5 z=7-10.4225=-3.4225 \text { i.e., } y=-1.1408 \\
x=-5-4 y+z=-5+4(1.1408)+2.0845=1.6479
\end{gathered}
$$

Hence $x=1.6479, y=-1.1408, z=2.0845$.

Deleteing $1^{\text {st }}, 2^{\text {nd }}$ row and column from the above equation,

$$
10^{7}\left[\begin{array}{cccc}
11.196 & -2.799 & -3.732 & 2.799 \\
-2.799 & 8.397 & -2.799 & 1.4 \\
-3.732 & -2.799 & 3.732 & -2.799 \\
2.799 & 1.4 & -2.799 & 2.799
\end{array}\right]\left[\begin{array}{c}
v_{2} \\
\theta_{2} \\
v_{3} \\
\theta_{4}
\end{array}\right]=\left\{\begin{array}{c}
-20 \times 10^{3} \\
0 \\
-20 \times 10^{3} \\
0
\end{array}\right\}
$$

Rounding off the values
$\left.\qquad \begin{array}{rrrr|r}11 & -3 & -4 & 3 & -2 \\ -3 & 8 & -3 & 1 & 0 \\ -4 & -3 & 4 & -3 & -2 \\ 3 & 1 & -3 & 3 & 0\end{array}\right)$

| $\mathrm{R} 3=11 * \mathrm{R} 3+4 * \mathrm{R} 1 \int_{11}$ | -3 | -4 | 3 | -2 |
| :---: | :---: | :---: | :---: | :---: |
| 0 | 9 | -6 | 4 | 0 |
| 0 | -5 | 0 | 3 | -6 |
| 3 | 1 | -3 | 3 | 0 |
| $\mathrm{R} 4=1 \mathrm{l}^{*} \mathrm{R} 4-3 * \mathrm{R} 1 \times 11$ | -3 | -4 | 3 | -2 |
| 0 | 9 | -6 | 4 | 0 |
| 0 | -5 | 0 | 3 | -6 |
| 0 | 20 | -21 | 24 | 6 |
| $R 3=9 * R 3+5 * R 2 \bigcirc 11$ | -3 | -4 | 3 | -2 |
|  | 9 | -6 | 4 | 0 |
|  | 0 | -30 | 47 | -54 |
| 0 | 20 | -21 | 24 | 6 |


| R4=9*R4-20*R2 | [ 11 | -3 | -4 | 3 | $-2$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | 0 | 9 | -6 | 4 | 0 |
|  | 0 | 0 | -30 | 47 | -54 |
|  | - 0 | 0 | -69 | 136 | 54 |
| $\begin{aligned} & R 3=-R 3 \\ & R 4=-R 4 \end{aligned}$ | [ 11 | -3 | -4 | 3 | -2 |
|  | 0 | 9 | -6 | 4 | 0 |
|  | 0 | 0 | 30 | -47 | 54 |
|  | $\bigcirc$ | 0 | 69 | -136 | $-54$ |
| $\mathrm{R} 4=30 * R 4-69 * R 3$ | [ 11 | -3 | -4 | 3 |  |
|  | 0 | 9 | -6 | 4 | 0 |
|  | 0 | 0 | 30 | -47 | 54 |
|  | 0 | 0 | 0 | -837 | -5346 |

$$
10^{3}\left[\begin{array}{rrrr}
11 & -3 & -4 & 3 \\
0 & 9 & -6 & 4 \\
0 & 0 & 30 & -47 \\
0 & 0 & 0 & -837
\end{array}\right]\left\{\begin{array}{l}
v_{2} \\
\theta_{2} \\
v_{3} \\
\theta_{4} \\
u_{4}
\end{array}\right\}=\left\{\begin{array}{r}
-2 \\
0 \\
54 \\
-5346
\end{array}\right\}
$$

On writing each equation separately one can calculate the unknowns

$$
[K]\{\delta\}=\{F\}
$$

$$
10^{7}\left[\begin{array}{cccccc}
7.464 & 5.598 & -7.464 & 5.598 & 0 & 0 \\
5.598 & 5.598 & -5.598 & 2.799 & 0 & 0 \\
\hline-7.464 & -5.598 & 11.196 & -2.799 & -3.732 & 2.799 \\
5.598 & 2.799 & -2.799 & 8.397 & -2.799 & 1.4 \\
0 & 0 & -3.732 & -2.799 & 3.732 & -2.799 \\
0 & 0 & 2.799 & 1.4 & -2.799 & 2.799
\end{array}\right]\left\{\begin{array}{c}
v_{1} \\
\theta_{1} \\
v_{2} \\
\theta_{2} \\
v_{3} \\
\theta_{3}
\end{array}\right\}=\left\{\begin{array}{c}
F_{1} \\
M_{1} \\
F_{2} \\
M_{2} \\
F_{3} \\
M_{3}
\end{array}\right\}=\left\{\begin{array}{c}
F_{1} \\
M_{1} \\
-20 \times 10^{3} \\
M_{2} \\
-20 \times 10^{3} \\
M_{3}
\end{array}\right\}
$$

M2, M3 are zero
From the diagram

$$
\begin{aligned}
F_{1} & =10^{7}\left[-7.464 v_{2}+5.598 \times \theta_{2}\right] \\
& =10^{7}[(7.464 \times 0.00375)-(5.598 \times 0.00428)] \\
F_{1} & =40305.6 \mathrm{~N} \\
M_{1} & =10^{7}\left[-5.598 \times v_{2}+2.799 \times \theta_{2}\right] \\
& =10^{7}[(5.598 \times 0.00375)-(2.799 \times 0.00428)] \\
& =10^{7} \times 0.00901 \mathrm{~N}-\mathrm{m}=90127.8 \mathrm{~N}-\mathrm{m}
\end{aligned}
$$



Determine the maximum deflection and slope for the simple supported beam subjected to uniformly load ' $q$ ' as shown in Figure.


Figure
Q. For the beam shown in figure calculate the deflection under the load for the beam.


Figure

Given that,
Young's modulus of the beam material, $E=20 \times 10^{6} \mathrm{~N} / \mathrm{cm}^{2}$
Moment of inertia, $I=2500 \mathrm{~cm}^{4}$
Point load, $\mathrm{W}=30 \mathrm{kN}=30000 \mathrm{~N}$
Length of each element, $l_{1}=l_{2}=2 \mathrm{~m}=200 \mathrm{~cm}$.


Figure: Discretization of Beam

Let $v_{1}, \theta_{1}, v_{2}, \theta_{2}, v_{3}, \theta_{3}$ are the nodal displacements.

Element stiffness matrix,

$$
[K]=\frac{E I}{l^{3}}\left[\begin{array}{cccc}
12 & 6 l & -12 & 6 l \\
6 l & 4 l^{2} & -6 l & 2 l^{2} \\
-12 & -6 l & 12 & -6 l \\
6 l & 2 l^{2} & -6 l & 4 l^{2}
\end{array}\right]
$$

For element - 1
Stiffness matrix for element (1) is given by,

$$
\begin{aligned}
& {\left[K_{1}\right]=\frac{E_{1} I_{1}}{l_{1}^{3}}\left[\begin{array}{cccc}
12 & 6 l_{1} & -12 & 6 l_{1} \\
6 l_{1} & 4 l^{2} & -6 l_{1} & 2 l^{2} \\
-12 & -6 l_{1} & 12 & -6 l_{1} \\
6 l_{1} & 2 l_{1}^{2} & -6 l_{1} & 4 l_{1}^{2}
\end{array}\right]} \\
& {\left[K_{1}\right]=\frac{20 \times 10^{6} \times 2500}{(200)^{3}}\left[\begin{array}{cccc}
12 & 1200 & -12 & 1200 \\
1200 & 160000 & -1200 & 80000 \\
-12 & -1200 & 12 & -1200 \\
1200 & 80000 & -1200 & 160000
\end{array}\right] \quad=6250\left[\begin{array}{cccc}
v_{1} & \theta_{1} & v_{2} & \theta_{2} \\
12 & 1200 & -12 & 1200 \\
1200 & 160000 & -1200 & 80000 \\
-12 & -1200 & 12 & -1200 \\
1200 & 80000 & -1200 & 160000
\end{array}\right] v_{1}} \\
& v_{1} \\
& v_{2} \\
& \theta_{2}
\end{aligned}
$$

## For element - 2

Stiffness matrix for element (2) is given by,

$$
\begin{aligned}
& {\left[K_{2}\right]=\frac{E_{2} I_{2}}{l_{2}^{3}}\left[\begin{array}{cccc}
12 & 6 l_{2} & -12 & 6 l_{2} \\
6 l_{2} & 4 l_{2}^{2} & -6 l_{2} & 2 l_{2}^{2} \\
-12 & -6 l_{2} & 12 & -6 l_{2} \\
6 l_{2} & 2 l_{2}^{2} & 6 l_{2} & 4 l_{2}^{2}
\end{array}\right]} \\
& {\left[K_{2}\right]=6250\left[\begin{array}{cccc}
v_{2} & \theta_{2} & v_{3} & \theta_{3} \\
12 & 1200 & -12 & 1200 \\
1200 & 160000 & -1200 & 80000 \\
-12 & -1200 & 12 & -1200 \\
1200 & 80000 & -1200 & 160000
\end{array}\right] \begin{array}{l}
v_{2} \\
\theta_{2} \\
v_{3} \\
\theta_{3}
\end{array}}
\end{aligned}
$$

$$
\left[K_{1}\right]=6250\left[\begin{array}{cccc}
v_{1} & \theta_{1} & v_{2} & \theta_{2} \\
12 & 1200 & -12 & 1200 \\
1200 & 160000 & -1200 & 80000 \\
-12 & -1200 & 12 & -1200 \\
1200 & 80000 & -1200 & 160000
\end{array}\right] \begin{aligned}
& v_{1} \\
& \theta_{1} \\
& v_{2} \\
& \theta_{2}
\end{aligned}
$$

## Global Stiffness Matrix

$$
\begin{aligned}
& {[K]=\left[K_{1}\right]+\left[K_{2}\right]} \\
& {[K]=6250 \times\left[\begin{array}{cccccc}
v_{1} & \theta_{1} & v_{2} & \theta_{2} & v_{3} & \theta_{3} \\
12 & 1200 & -12 & 1200 & 0 & 0 \\
1200 & 160000 & -1200 & 80000 & 0 & 0 \\
-12 & -1200 & 24 & 0 & -12 & 1200 \\
1200 & 80000 & 0 & 320000 & -1200 & 80000 \\
0 & 0 & -12 & -1200 & 12 & -1200 \\
0 & 0 & 1200 & 80000 & -1200 & 160000
\end{array}\right] \begin{array}{l}
v_{1} \\
\theta_{1} \\
v_{2} \\
\theta_{2} \\
v_{3} \\
\theta_{3}
\end{array}}
\end{aligned}
$$

The finite element equation is given by,

$$
[K]\{\delta\}=\{F\}
$$



$$
v_{1}=\theta_{1}=v_{3}=0
$$

The rows and columns related to degree of freedom 1,2 and 5 are deleted from $[K]$ matrix.

$$
\therefore 6250 \times\left[\begin{array}{ccc}
24 & 0 & 1200 \\
0 & 320000 & 80000 \\
1200 & 80000 & 160000
\end{array}\right]\left\{\begin{array}{l}
v_{2} \\
\theta_{2} \\
\theta_{3}
\end{array}\right\}=\left\{\begin{array}{c}
-30000 \\
0 \\
0
\end{array}\right\}
$$

On solving above equations,

$$
\begin{aligned}
& 24 v_{2}+1200 \theta_{3}=-4.8 \Rightarrow v_{2}=\left(\frac{-4.8-1200 \theta_{3}}{24}\right) \\
& 320000 \theta_{2}+80000 \theta_{3}=0 \Rightarrow \theta_{2}=-\frac{\theta_{3}}{4} \\
& 1200 v_{2}+80000 \theta_{2}+160000 \theta_{3}=0 \\
& 1200\left[\frac{-4.8-1200 \theta_{3}}{24}\right]+80000\left[\frac{-\theta_{3}}{4}\right]+160000 \theta_{3}=0 \\
& -240-60000 \theta_{3}-20000 \theta_{3}+160000 \theta_{3}=0 \\
& \therefore \theta_{3}=0.003 \\
& v_{2}=\frac{-4.8-1200(0.003)}{24}=-0.35 \mathrm{~cm} \\
& \theta_{2}=\frac{-0.003}{4}=-0.00075 \mathrm{rad}
\end{aligned}
$$

The nodal displacement vector is given by,

$$
\{\delta\}=\left[\begin{array}{llllll}
0 & 0 & -0.35 \mathrm{~cm} & -0.00075 \mathrm{rad} & 0 & 0.003 \mathrm{~cm}
\end{array}\right]^{T}
$$

## Gravity loading

Gravity loading is a typical body force and is given by ( $\rho g$ ) per unit volume or ( $\rho A g$ ) per unit length, where $\rho$ is the mass density of the material. The equivalent nodal force vector for the distributed body force can be obtained as

$$
\left\{f^{e}=\int_{v}[N]^{T}(\rho g) d v=\int_{0}^{L}[N]^{T}(\rho g) A d x=(\rho A g)\left\{\begin{array}{r}
\ell / 2  \tag{4.124}\\
L^{2} / 12 \\
\ell / 2 \\
-L^{2} / 12
\end{array}\right\}\right.
$$

$$
W_{q}=\int_{0}^{L} q v d s
$$

$$
\{f\}^{e}=\int_{0}^{L}[N]^{T} q_{0} d x=\left\{\begin{array}{c}
q_{0} L / 2 \\
q_{0} L^{2} / 12 \\
q_{0} L / 2 \\
-q_{0} L^{2} / 12
\end{array}\right\}
$$



Q. For the beam and loading shown in the figure determine,
(i) The slopes at 2 and 3
(ii) The vertical deflection at the midpoint of the distributed load.



Figure

Given that,
Young's modulus, $E=200 \mathrm{GPa}=200 \times 10^{9} \mathrm{~N} / \mathrm{m}^{2}$
Moment of inertia, $I=4 \times 10^{6} \mathrm{~mm}^{4}=4 \times 10^{-6} \mathrm{~m}^{4}$
Length of elements, $l_{1}=l_{2}=1 \mathrm{~m}$

Let, $v_{1}, \theta_{1}, v_{2}, \theta_{2}, v_{3}, \theta_{3}$ are the nodal displacements.
Nodal displacement vector $\{\delta\}$ is given as,

$$
\{\delta\}=\left[\begin{array}{llllll}
v_{1} & \theta_{1} & v_{2} & \theta_{2} & v_{3} & \theta_{3}
\end{array}\right]
$$



Figure (1): Discretization of Beam

Element stiffness matrix is given by,

## For Element (1)

Stiffness matrix,

$$
[K]=\frac{E I}{l^{3}}\left[\begin{array}{cccc}
12 & 6 l & -12 & 6 l \\
6 l & 4 l^{2} & -6 l & 2 l^{2} \\
-12 & -6 l & 12 & -6 l \\
6 l & 2 l^{2} & -6 l & 4 l^{2}
\end{array}\right]
$$

$$
\left.\begin{array}{l}
{\left[K_{1}\right]=\frac{200 \times 10^{9} \times 4 \times 10^{-6}}{1^{3}}\left[\begin{array}{cccc}
12 & 6 & -12 & 6 \\
6 & 4 & -6 & 2 \\
-12 & -6 & 12 & -6 \\
6 & 2 & -6 & 4
\end{array}\right]} \\
v_{1} \quad \theta_{1}
\end{array} v_{2} \theta_{2},\left[\begin{array}{cccc}
12 & 6 & -12 & 6 \\
6 & 4 & -6 & 2 \\
-12 & -6 & 12 & -6 \\
6 & 2 & -6 & 4
\end{array}\right] \begin{array}{l}
v_{1} \\
\theta_{1} \\
v_{2}
\end{array}\right]=8 \times 10^{5} .
$$

Element stiffness matrix is given by,

## For Element (2)

## Stiffness matrix,

$$
\left[K_{2}\right]=\frac{200 \times 10^{9} \times 4 \times 10^{-6}}{1^{2}}\left[\begin{array}{cccc}
12 & 6 & -12 & 6 \\
6 & 4 & -6 & 2 \\
-12 & -6 & 12 & -6 \\
6 & 2 & -6 & 4
\end{array}\right]
$$

$$
\left[K_{2}\right]=8 \times 10^{5}\left[\begin{array}{cccc}
v_{2} & \theta_{2} & v_{3} & \theta_{3} \\
{\left[\begin{array}{ccc}
12 & 6 & -12
\end{array}\right.} & 6 \\
6 & 4 & -6 & 2 \\
-12 & -6 & 12 & -6 \\
6 & 2 & -6 & 4
\end{array}\right] \begin{aligned}
& v_{2} \\
& \theta_{2} \\
& v_{3} \\
& \theta_{3}
\end{aligned}
$$

$$
[K]=\frac{E I}{l^{3}}\left[\begin{array}{cccc}
12 & 6 l & -12 & 6 l \\
6 l & 4 l^{2} & -6 l & 2 l^{2} \\
-12 & -6 l & 12 & -6 l \\
6 l & 2 l^{2} & -6 l & 4 l^{2}
\end{array}\right]
$$

Then, the global stiffness matrix ' $K$ ' is given by,

$$
\left.\begin{array}{rl}
{[K]} & =\left[K_{1}\right]+\left[K_{2}\right] \\
& =8 \times 10^{5}\left[\begin{array}{cccccc}
12 & 6 & -12 & 6 & 0 & 0 \\
6 & 4 & -6 & 2 & 0 & 0 \\
-12 & -6 & (12+12) & (-6+6) & -12 & 6 \\
6 & 2 & (-6+6) & (4+4) & -6 & 2 \\
0 & 0 & -12 & -6 & 12 & -6 \\
0 & 0 & 6 & 2 & -6 & 4
\end{array}\right] \\
{[K]} & =8 \times 10^{5}\left[\begin{array}{cccccc}
1 & \theta_{1} & v_{2} & \theta_{2} & v_{3} & \theta_{3} \\
12 & 6 & -12 & 6 & 0 & 0 \\
6 & 4 & -6 & 2 & 0 & 0 \\
-12 & -6 & 24 & 0 & -12 & 6 \\
6 & 2 & 0 & 8 & -6 & 2 \\
0 & 0 & -12 & -6 & 12 & -6 \\
0 & 0 & 6 & 2 & -6 & 4
\end{array}\right] \theta_{1}^{v_{1}} \\
\theta_{1}
\end{array}\right] v_{2} \begin{aligned}
& \theta_{2}
\end{aligned}
$$

Finite element equation is given by,

$$
[\mathrm{K}]\{\delta\}=\{\mathrm{F}\}
$$

$$
8 \times 10^{5}\left[\begin{array}{cccccc}
12 & 6 & -12 & 6 & 0 & 0 \\
6 & 4 & -6 & 2 & 0 & 0 \\
-12 & -6 & 24 & 0 & -12 & 6 \\
6 & 2 & 0 & 8 & -6 & 2 \\
0 & 0 & -12 & -6 & 12 & -6 \\
0 & 0 & 6 & 2 & -6 & 4
\end{array}\right]\left\{\begin{array}{c}
v_{1} \\
\theta_{1} \\
v_{2} \\
\theta_{2} \\
v_{3} \\
\theta_{3}
\end{array}\right\}=\left\{\begin{array}{c}
F_{1}+F_{1 d} \\
M_{1}+M_{1 d} \\
F_{2}+F_{2 d} \\
M_{2}+M_{2 d} \\
F_{3}+F_{3 d} \\
M_{3}+M_{3 d}
\end{array}\right\}
$$

Following are due to UDL $F_{1 a+} F_{2 d}, F_{3 d} \rightarrow$ Nodal forces $M_{1 a}, M_{2 a}, M_{3 d} \rightarrow$ Nodal bending moments


Applying the boundary conditions,

$$
v_{1}=0, \theta_{1}=0, v_{2}=0, v_{3}=0, F_{1 d}=0, M_{1 d}=0
$$

$$
F_{2 d}=-6000 \mathrm{~N}, M_{2 d}=-1000 \mathrm{~N}-\mathrm{m}, F_{3 d}=-6000 \mathrm{~N},
$$

$$
M_{3 d}=+1000 \mathrm{~N}-\mathrm{m}, M_{2}=0, M_{3}=0
$$



Figure

$$
\begin{aligned}
& F_{2 d}=F_{3 d}=\frac{W l}{2}=\frac{12000 \times 1}{2}=6000 \mathrm{~N} \\
& M_{2 d}=M_{3 d}=\frac{W l^{2}}{12}=\frac{12000 \times 1^{2}}{12}=1000 \mathrm{~N}-\mathrm{m}
\end{aligned}
$$

$$
8 \times 10^{5}\left[\begin{array}{cccccc}
12 & 6 & -12 & 6 & 0 & 0 \\
6 & 4 & -6 & 2 & 0 & 0 \\
-12 & -6 & 24 & 0 & -12 & 6 \\
6 & 2 & 0 & 8 & -6 & 2 \\
0 & 0 & -12 & -6 & 12 & -6 \\
0 & 0 & 6 & 2 & -6 & 4
\end{array}\right]\left[\begin{array}{c}
v_{1} \\
\theta_{1} \\
v_{2} \\
\theta_{2} \\
v_{3} \\
\theta_{3}
\end{array}\right]=\left\{\begin{array}{c}
F_{1} \\
M_{1} \\
F_{2}-6000 \\
-1000 \\
F_{3}-6000 \\
1000
\end{array}\right\}
$$

$$
\begin{aligned}
& 8 \times 10^{5}\left[\begin{array}{ll}
8 & 2 \\
2 & 4
\end{array}\right]\left\{\begin{array}{l}
\theta_{2} \\
\theta_{3}
\end{array}\right\}=\left\{\begin{array}{c}
-1000 \\
1000
\end{array}\right\} \\
& 8 \times 10^{5}\left[8 \theta_{2}+2 \theta_{3}\right]=-1000 \\
& 8 \times 10^{5}\left[2 \theta_{3}+4 \theta_{3}\right]=1000
\end{aligned}
$$

On solving the above equations,

$$
\begin{aligned}
& \theta_{2}=-2.679 \times 10^{-4} \mathrm{rad} \\
& \theta_{3}=4.464 \times 10^{-4} \mathrm{rad}
\end{aligned}
$$

$\therefore \quad$ Nodal displacement vector,

$$
[\delta]=\left[\begin{array}{llllll}
0 & 0 & 0 & -2.679 \times 10^{-4} & 0 & 4.464 \times 10^{-4}
\end{array}\right]^{T}
$$

## Vertical Deflection at the Midpoint of the Distributed

## Load

Consider element (2)


Figure (3): Element (2)
For a beam, element, the vertical deflection is given by,

$$
\begin{equation*}
V=N_{1} v_{1}+N_{2} \theta_{1}+N_{3} v_{2}+N_{4} \theta_{2} \tag{1}
\end{equation*}
$$

Where,

$$
N_{1}, N_{2}, N_{3}, N_{4} \text { - Shape functions }
$$

Consider element (2) as separate element and mark as node 1 and 2 as shown below,


Figure (4)

Now,

$$
\begin{array}{lr}
v_{1}=0 \\
\theta_{1}=-2.679 \times 10^{-4} \mathrm{rad} & N_{4} \theta_{2}=\left(\frac{-x^{2}}{l}+\frac{x^{3}}{l^{2}}\right) \theta_{2} \\
v_{2}=0 & =\left(\frac{-(0.5)^{2}}{1}+\frac{(0.5)^{3}}{1^{2}}\right)\left(4.464 \times 10^{-4}\right) \\
\theta_{2}=4.464 \times 10^{-4} \mathrm{rad} & N_{4} \theta_{2}=-5.58 \times 10^{-5} \\
N_{1} v_{1}=0 & V=N_{1} v_{1}+N_{2} \theta_{1}+N_{3} v_{2}+N_{4} \theta_{2} \\
N_{2} \theta_{1}=\left[x-\frac{2 x^{2}}{l}+\frac{x^{3}}{l^{2}}\right]_{1} & V=0-3.348 \times 10^{-5}+0-5.58 \times 10^{-5} \\
& =\left[0.5-\frac{2(0.5)^{2}}{1}+\frac{(0.5)^{3}}{1^{3}}\right]\left(-2.679 \times 10^{-4}\right)
\end{array} r=8028 \times 10^{-5} \mathrm{~m} .
$$

w kN/m

(i) Load Diagram


Force vector,

$$
\{F\}=\left[\left.\begin{array}{c}
-\frac{w L}{4}+R_{1} \\
-\frac{w L^{2}}{48} \\
-\frac{w L}{2} \\
0 \\
-\frac{w L}{4}+R_{3} \\
\frac{w L^{2}}{48}
\end{array} \right\rvert\,\right.
$$



Determine the deflection in the beam, loaded as shown in figure, at the mid-span and at a length of 0.5 m from left support. Determine also the reactions at the fixed ends.
$E=200 \mathrm{GPa} . \mathrm{I}_{1}=20 \times 10^{-6} \mathrm{~m}^{4} . \mathrm{I}_{2}=10 \times 10^{-6} \mathrm{~m}^{4}$.

5.2. A three-span beam is shown in Fig. P5.2. Determine the deflection curve of the beam and evaluate the reactions at the supports.

2. Determine the deflection in the beam, loaded as shown in figure, at the mid-span and at a length of 0.5 m from left support. Determine also the reactions at the fixed ends.
$E=200 \mathrm{GPa} . \mathrm{I}_{1}=20 \times 10^{-6} \mathrm{~m}^{4} . \mathrm{I}_{2}=10 \times 10^{-6} \mathrm{~m}^{4}$.


## Unit 3

## Plane Problems (Two Dimensional Problems)


(i) Triangle

(ii) Rectangle

(iii) Quadrilateral

(iv) Parallelogram Figure: Two Dimensional Elements


Figure (1): Constant Strain Triangle (CST)


FIGURE 6.1 Two-dimensional problem.

$$
\mathbf{f}=\left[\begin{array}{ll}
f_{x} & f_{y}
\end{array}\right]^{\mathrm{T}} \quad \mathbf{T}=\left[\begin{array}{ll}
T_{x} & T_{y}
\end{array}\right]^{\mathrm{T}} \quad \text { and } \quad d V=t d A
$$

## Plane (2-D) Problems

- Plane stress:

$$
\begin{equation*}
\sigma_{z}=\tau_{y z}=\tau_{z x}=0 \quad\left(\varepsilon_{z} \neq 0\right) \tag{1}
\end{equation*}
$$

A thin planar structure with constant thickness and loading within the plane of the structure ( $x y$-plane).


- Plane strain:

$$
\begin{equation*}
\varepsilon_{z}=\gamma_{y z}=\gamma_{z x}=0 \quad\left(\sigma_{z} \neq 0\right) \tag{2}
\end{equation*}
$$

A long structure with a uniform cross section and transverse loading along its length ( $z$-direction).


## Stress-Strain-Temperature (Constitutive) Relations

For elastic and isotropic materials, we have,

$$
\left\{\begin{array}{c}
\varepsilon_{4}  \tag{3}\\
\varepsilon_{j} \\
\gamma_{*}
\end{array}\right\}=\left[\begin{array}{ccc}
1 / E & -v / E & 0 \\
-v / E & 1 / E & 0 \\
0 & 0 & 1 / G
\end{array}\right]\left\{\begin{array}{l}
\sigma_{v} \\
\sigma_{i} \\
\tau_{n}
\end{array}\right\}+\left\{\begin{array}{c}
\varepsilon_{+0} \\
\varepsilon_{, 0} \\
\gamma_{* b}
\end{array}\right\}
$$

where $\varepsilon_{\mathrm{s}}$ is the initial strain, $E$ the Young's modulus, $v$ the Poisson's ratio and $G$ the shear modulus. Note that,

$$
\begin{equation*}
G=\frac{E}{2(1+v)} \tag{4}
\end{equation*}
$$

which means that there are only two independent materials constants for homogeneous and isotropic materials.

We can also express stresses in terms of strains by solving the above equation.

$$
\left\{\begin{array}{l}
\sigma_{n}  \tag{5}\\
\sigma_{,} \\
\tau_{0}
\end{array}\right\}=\frac{E}{1-v^{2}}\left[\begin{array}{ccc}
1 & v & 0 \\
v & 1 & 0 \\
0 & 0 & (1-v) / 2
\end{array}\right]\left\{\left\{\begin{array}{l}
e_{v} \\
\varepsilon_{v} \\
\gamma_{v}
\end{array}\right\}-\left\{\begin{array}{l}
c_{v 0} \\
\varepsilon_{v 0} \\
\gamma_{v 0}
\end{array}\right\}\right)
$$

The above relations are valid for plane stress case. For plane strain case, we need to replace the material constants in the above equations in the following fashion,

$$
\begin{align*}
& E \rightarrow \frac{E}{1-v^{2}} \\
& v \rightarrow \frac{v}{1-v}  \tag{6}\\
& G \rightarrow G
\end{align*}
$$

For example, the stress is related to strain by
in the plame strain case.

Initial strains due to temperature change (thermal loading) is given by.

$$
\left\{\begin{array}{c}
\varepsilon_{v 0}  \tag{7}\\
\varepsilon_{v n} \\
y_{v ⿻}
\end{array}\right\}=\left\{\begin{array}{c}
a \Delta T \\
a \Delta T \\
0
\end{array}\right\}
$$

where $\alpha$ is the coefficient of themal expansion, $\Delta T$ the change of temperature. Note that if the structure is free to deform under thermal loading, there will be no (elastic) stresses in the structure.

$$
\boldsymbol{\sigma}=\mathbf{D} \boldsymbol{\epsilon}
$$

$$
\boldsymbol{\epsilon}=\left[\frac{\partial u}{\partial x} \frac{\partial v}{\partial y}\left(\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}\right)\right]^{\mathrm{T}}
$$

## Strain and Displacement Relations

Plane stress:
$\mathbf{D}=\frac{E}{1-v^{2}}\left(\begin{array}{ccc}1 & v & 0 \\ v & 1 & 0 \\ 0 & 0 & \frac{1-v}{2}\end{array}\right)$

Plane strain:

$$
\mathbf{D}=\frac{E}{(1+v)(1-2 v)}\left(\begin{array}{ccc}
1-v & v & 0 \\
v & 1-v & 0 \\
0 & 0 & \frac{1-2 v}{2}
\end{array}\right)
$$

For small strains and small rotations, we have,

$$
\varepsilon_{x}=\frac{\partial u}{\partial x}, \quad \varepsilon_{y}=\frac{\partial v}{\partial y}, \quad \gamma_{x y}=\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}
$$

In matrix form,

$$
\left\{\begin{array}{c}
\varepsilon_{x}  \tag{8}\\
\varepsilon_{y} \\
\gamma_{x y}
\end{array}\right\}=\left[\begin{array}{cc}
\partial / \partial x & 0 \\
0 & \partial / \partial y \\
\partial / \partial y & \partial / \partial x
\end{array}\right]\left\{\begin{array}{l}
u \\
v
\end{array}\right\}
$$

From this relation, we know that the strains (and thus stresses) are one order lower than the displacements, if the displacements are represented by polynomials.


FIGURE 6.2 Finite element discretization.

Constant term (1) -

Linear terms (2) -

Quadratic terms (6) -

Cubic terms (10) -
x

$$
x^{2} \quad x y
$$

$$
x^{3} \quad x^{2} y \quad x y^{2}
$$

## Pascal triangle

In order to develop a polynomial with three terms,
Expression to be selected is,

$$
u=a_{1}+a_{2} x+a_{3} y
$$

In order to develop a polynomial with four terms,
Expression to be selected is,

$$
u=a_{1}+a_{2} x+a_{3} y+a_{4} x y
$$

In order to develop a polynomial with six terms,
Expression to be selected is,

$$
u=a_{1}+a_{2} x+a_{3} y+a_{4} x^{2}+a_{5} x y+a_{6} y^{2}
$$

## LINEAR Displacement TRIANGLE or (CST)

A linear triangle is a plane triangle whose field quantity varies linearly with Cartesian coordinates $x$ and $y$. In stress analysis, a linear displacement field produces a constant strain field, so the element may be called a constant-strain triangle (CST).

1. Constant Strain Triangle (CST): It consists of three nodes and six unknown nodal displacements. Its field varies linearly with coordinates $x$ and $y$, giving rise to a linear displacement and a constant strain field.

Let, at a particular node,
$u$-Displacement along x -axis
$v$-Displacement along y-axis.
Then, components of displacement for CST element are given by,

$$
\begin{aligned}
& u=a_{1}+a_{2} x+a_{3} y \\
& v=a_{4}+a_{5} x+a_{6} y
\end{aligned}
$$


$x^{2} \quad x y \quad y^{2}$



Figure (1): Constant Strain Triangle (CST)

$$
\begin{aligned}
& u=a_{1}+a_{2} x+a_{3} y \\
& v=a_{4}+a_{5} x+a_{6} y
\end{aligned} \quad u=\left\lfloor\begin{array}{lll}
1 & x & y
\end{array}\right\rfloor\left\{\begin{array}{l}
a_{1} \\
a_{2} \\
a_{3}
\end{array}\right\} \quad v=\left\lfloor\begin{array}{lll}
1 & x & y
\end{array}\right]\left\{\begin{array}{l}
a_{4} \\
a_{5} \\
a_{6}
\end{array}\right\}
$$

$$
\begin{array}{ccc}
\varepsilon_{x}=\frac{\partial u}{\partial x} & \varepsilon_{y}=\frac{\partial v}{\partial y} & \gamma_{x y}=\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x} \\
\varepsilon_{x}=a_{2} & \varepsilon_{y}=a_{6} & \gamma_{x y}=a_{3}+a_{5}
\end{array}
$$

$$
\begin{gathered}
\varepsilon_{x}=\frac{\partial u}{\partial x} \quad \varepsilon_{y}=\frac{\partial v}{\partial y} \quad \gamma_{x y}=\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x} \\
\left\{\begin{array}{l}
\varepsilon_{x} \\
\varepsilon_{y} \\
\gamma_{x y}
\end{array}\right\}=\left[\begin{array}{cc}
\frac{\partial}{\partial x} & 0 \\
0 & \frac{\partial}{\partial y} \\
\frac{\partial}{\partial y} & \frac{\partial}{\partial x}
\end{array}\right]\left[\begin{array}{l}
u \\
v
\end{array}\right\}
\end{gathered}
$$

$\qquad$ displacement field varies quadratically
2. Linear Strain Triangle (LST): It consists of three primary nodes and three secondary nodes at the mid-points of the sides of the triangle. Each node possess two degrees of freedom (DOF). Therefore each element has 12 DOF. Displacement function for this element is a quadratic equation and strain field varies linearly
Components of displacement for LST element are given by,

$$
\begin{aligned}
& u=a_{1}+a_{2} x+a_{3} y+a_{4} x^{2}+a_{5} x y+a_{6} y^{2} \\
& v=a_{7}+a_{8} x+a_{9} y+a_{10} x^{2}+a_{11} x y+a_{12} y^{2}
\end{aligned}
$$

$$
\varepsilon_{x}=\frac{\partial u}{\partial x} \quad \varepsilon_{y}=\frac{\partial v}{\partial y} \quad \gamma_{x y}=\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}
$$

## CST element:(constant strains, linear displacements)

shape functions $N_{1}, N_{2}$ and $N_{3}$. Variation of these shape functions occurs linearly in CST element. These shape functions have a value of unity at their corresponding nodes and reduce to the value of zero at other nodes. Only two of the shape functions are independent and are represented by $\xi$ and $\eta$ in natural coordinate system.

$$
\begin{aligned}
& N_{1}=\xi, N_{2}=\eta, N_{3}=1-\xi-\eta \\
\therefore \quad & N_{1}+N_{2}+N_{3}=1
\end{aligned}
$$

In two dimensional problem, the $x-y$ coordinates are

| $\underset{\text { alpha }}{\alpha}$ | $\underset{\text { beta }}{\beta}$ | $\underset{\text { gamma }}{\gamma}$ | $\oint_{\text {delta }}$ |
| :---: | :---: | :---: | :---: |
| $\underset{\text { epsilon }}{\mathcal{E}}$ | $\underbrace{\zeta}_{\text {zeta }}$ | $\eta_{\text {eta }}$ | $\underset{\text { theta }}{\theta}$ |
| $\underset{\text { iota }}{\mathbf{l}}$ | $\underset{\text { kappa }}{\mathcal{K}}$ | $\underset{\text { lambda }}{\lambda}$ | $\underset{\text { mu }}{\mu}$ |
| $\begin{gathered} V \\ \text { nu } \end{gathered}$ | ${\underset{x i}{ }}_{\underline{x}}$ | $\mathrm{O}$ omicron | $\begin{aligned} & \pi \\ & \mathrm{pi} \end{aligned}$ |
| $\underset{\text { tho }}{\rho}$ | $\underset{\text { sigma }}{\sigma}$ | $\underset{\text { tau }}{\tau}$ | U upsilon |
| $\underset{\text { phi }}{\oint}$ | $\underset{\text { chi }}{\chi}$ | $\underset{\text { psi }}{\Psi}$ | $\underset{\text { omega }}{\omega}$ |


(c)

FIGURE 6.4 Shape functions.

Fig. (a)
i.e, at node-1,

$$
\begin{aligned}
& N_{1}=1 \\
& N_{2}=N_{3}=0
\end{aligned}
$$

Fig. (b) At node-2,

$$
\begin{aligned}
& N_{2}=1 \\
& N_{1}=N_{3}=0
\end{aligned}
$$

Fig. (c) At node-3,

$$
\begin{aligned}
& N_{3}=1 \\
& N_{2}=N_{1}=0
\end{aligned}
$$

The shape functions can be physically represented by area coordinates. A point $(x, y)$ in a triangle divides it into three areas, $A_{1}, A_{2}$, and $A_{3}$, as shown in Fig. 6.5. The shape functions $N_{1}, N_{2}$, and $N_{3}$ are precisely represented by

$$
\begin{equation*}
N_{1}=\frac{A_{1}}{A} \quad N_{2}=\frac{A_{2}}{A} \quad N_{3}=\frac{A_{3}}{A} \tag{6.11}
\end{equation*}
$$

where $A$ is the area of the element. Clearly, $N_{1}+N_{2}+N_{3}=1$ at all points inside the triangle.


The displacements inside the element are now written using the shape functions and the nodal values of the unknown displacement field.

The displacements can be written as,

$$
\begin{aligned}
& u=N_{1} u_{1}+N_{2} u_{2}+N_{3} u_{3} \\
& v=N_{1} v_{1}+N_{2} v_{2}+N_{3} v_{3}
\end{aligned}
$$

Figure
This is in the form of,

$$
u=N \delta
$$

Where,

$$
\begin{aligned}
& N=\left[\begin{array}{lccccc}
N_{1} & 0 & N_{2} & 0 & N_{3} & 0 \\
0 & N_{1} & 0 & N_{2} & 0 & N_{3}
\end{array}\right] \\
& u=\left\{\begin{array}{l}
u \\
v
\end{array}\right\} \text { and } \\
& \delta=\left\{\begin{array}{l}
u_{1} \\
v_{1} \\
u_{2} \\
v_{2} \\
u_{3} \\
v_{3}
\end{array}\right\}
\end{aligned}
$$

## Isoparametric Representation

It is possible to represent the coordinates $(x, y)$ of any point ' $P$ ' within the linear triangular element in terms of nodal coordinates by employing the same shape functions used to represent displacements $u$ and $v$. Such a method of representation is known as 'isoparametric representation '.

Thus, coordinates $(x, y)$ of point ' $P$ ',

$$
\begin{aligned}
& N_{1}=\xi \\
& N_{2}=\eta \\
& N_{3}=1-\xi-\eta
\end{aligned}
$$

$$
\text { i.e., } \begin{aligned}
x & =N_{1} x_{1}+N_{2} x_{2}+N_{3} x_{3} \\
& =\xi x_{1}+\eta x_{2}+(1-\xi-\eta) x_{3} \\
& =\left(x_{1}-x_{3}\right) \xi+\left(x_{2}-x_{3}\right) \eta+x_{3} \\
& \therefore x=x_{13} \xi+x_{23} \eta+x_{3} \\
y & =N_{1} y_{1}+N_{2} y_{2}+N_{3} y_{3} \\
& =\xi y_{1}+\eta y_{2}+(1-\xi-\eta) y_{3} \\
& =\left(y_{1}-y_{3}\right) \xi+\left(y_{2}-y_{3}\right) \eta+y_{3} \\
\therefore \quad y & =y_{13} \xi+y_{23} \eta+y_{3}
\end{aligned}
$$



Figure

$$
\begin{array}{ll}
x=N_{1} x_{1}+N_{2} x_{2}+N_{3} x_{3} & x=x_{13} \xi+x_{23} \eta+x_{3} \\
y=N_{1} y_{1}+N_{2} y_{2}+N_{3} y_{3} & y=y_{13} \xi+y_{23} \eta+y_{3}
\end{array}
$$

## Example 6.1

Evaluate the shape functions $N_{1}, N_{2}$, and $N_{3}$ at the interior point $P$ for the triangular element shown in Fig. E6.1.



Figure

FIGURE E6.1 Examples 6.1 and 6.2.

$$
\begin{array}{cc}
\begin{array}{r}
x=N_{1} x_{1}+N_{2} x_{2}+N_{3} x_{3} \\
y=N_{1} y_{1}+N_{2} y_{2}+N_{3} y_{3}
\end{array} & \begin{array}{r}
N_{1}=\xi \\
N_{2}=\eta \\
N_{3}=1
\end{array} \\
\begin{array}{c}
3.85=1.5 N_{1}+7 N_{2}+4 N_{3} \\
4.8=2 N_{1}+3.5 N_{2}+7 N_{3}
\end{array} & \\
3.85=-2.5 \xi+3 \eta+4 & \\
4.8=-5 \xi-3.5 \eta+7 &
\end{array}
$$

$2.5 \xi-3 \eta=0.15$
$5 \xi+3.5 \eta=2.2$
Solving the equations, we obtain $\xi=0.3$ and $\eta=0.2$,

$$
N_{1}=0.3 \quad N_{2}=0.2 \quad N_{3}=0.5
$$

Q. Estimate the shape functions of a triangular element at the point $P(22,44)$ of a CST with the coordinates $1(0,0), 2(46,8)$ and $3(18,62)$. All dimensions are in mm .

Ans: $\mathrm{N} 1=0.11, \mathrm{~N} 2=0.21, \mathrm{~N} 3=0.68$

The displacements can be written as,

$$
\begin{aligned}
& u=N_{1} u_{1}+N_{2} u_{2}+N_{3} u_{3} \\
& v=N_{1} v_{1}+N_{2} v_{2}+N_{3} v_{3}
\end{aligned}
$$

Similarly we can write the coordinates:

$$
\begin{aligned}
& x=N_{1} x_{1}+N_{2} x_{2}+N_{3} x_{3} \\
& y=N_{1} y_{1}+N_{2} y_{2}+N_{3} y_{3}
\end{aligned}
$$



$$
N_{1}=\xi
$$

$$
N_{2}=\eta
$$

$$
N_{3}=1-\xi-\eta
$$



Figure

$$
\begin{aligned}
& x=x_{13} \xi+x_{23} \eta+x_{3} \\
& y=y_{13} \xi+y_{23} \eta+y_{3}
\end{aligned}
$$

similarly

$$
\begin{aligned}
& \mathbf{u}=\mathbf{u}_{13} \xi+\mathbf{u}_{23} \eta+\mathbf{u}_{3} \\
& \mathbf{v}=\mathbf{v}_{13} \xi+\mathbf{v}_{23} \eta+\mathbf{v}_{3}
\end{aligned}
$$

(Recall the strains definitions)


Figure 3.1-1. An infinitesimal rectangle, subjected to (a) $x$-direction normal strain, (b) $y$-direction normal strain, and (c) shear strain.

$$
\varepsilon_{x}=\frac{\partial u}{\partial x} \quad \varepsilon_{y}=\frac{\partial v}{\partial y} \quad \gamma_{x y}=\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}
$$

: In fluid mechanics, velocity V=u i + v j + w k
The velocity of a fluid is not only a function of time but also of space:

$$
u=f(x, y, z, t)
$$

By the chain rule of differentiation,

$$
\frac{d u}{d t}=\frac{\partial u}{\partial t}+\frac{\partial u}{\partial x} \frac{\partial x}{\partial t}+\frac{\partial u}{\partial y} \frac{\partial y}{\partial t}+\frac{\partial u}{\partial z} \frac{\partial z}{\partial t}
$$

Using the chain rule for partial derivatives of $u$, ( $u$ is the $x$-displacement of nodes)

$$
\begin{aligned}
& \frac{\partial u}{\partial \xi}=\frac{\partial u \partial x}{\partial x \partial \xi}+\frac{\partial u \partial y}{\partial y \partial \xi} \\
& \frac{\partial u}{\partial \eta}=\frac{\partial u \partial x}{\partial x \partial \eta}+\frac{\partial u \partial y}{\partial y \partial \eta}
\end{aligned}
$$

$$
\begin{aligned}
& \frac{\partial u}{\partial \xi}=\frac{\partial u \partial x}{\partial x \partial \xi}+\frac{\partial u \partial y}{\partial y \partial \xi} \\
& \frac{\partial u}{\partial \eta}=\frac{\partial u \partial x}{\partial x \partial \eta}+\frac{\partial u \partial y}{\partial y \partial \eta}
\end{aligned}
$$

which can be written in matrix notation as

$$
\left\{\begin{array}{l}
\frac{\partial u}{\partial \xi} \\
\frac{\partial u}{\partial \eta}
\end{array}\right\}=\left[\begin{array}{ll}
\frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} \\
\frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta}
\end{array}\right]\left\{\begin{array}{l}
\frac{\partial u}{\partial x} \\
\frac{\partial u}{\partial y}
\end{array}\right\}
$$


(2 * 2 ) square matrix is denoted as the Jacobian of the transformation, J :

$$
\mathbf{J}=\left[\begin{array}{ll}
\frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} \\
\frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta}
\end{array}\right]=\left[\begin{array}{ll}
x_{13} & y_{13} \\
x_{23} & y_{23}
\end{array}\right]
$$

$$
\begin{aligned}
& x=x_{13} \xi+x_{23} \eta+x_{3} \\
& y=y_{13} \xi+y_{23} \eta+y_{3}
\end{aligned}
$$

$$
\left\{\begin{array}{l}
\frac{\partial u}{\partial \xi} \\
\frac{\partial u}{\partial \eta}
\end{array}\right\}=\left[\begin{array}{ll}
\frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} \\
\frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta}
\end{array}\right]\left\{\begin{array}{l}
\frac{\partial u}{\partial x} \\
\frac{\partial u}{\partial y}
\end{array}\right\} \longmapsto\left\{\begin{array}{l}
\frac{\partial u}{\partial x} \\
\frac{\partial u}{\partial y}
\end{array}\right\}=\mathbf{J}^{-1}\left\{\begin{array}{l}
\frac{\partial u}{\partial \xi} \\
\frac{\partial u}{\partial \eta}
\end{array}\right\}
$$

## where $\mathbf{J}^{-1}$ is the inverse of the Jacobian $\mathbf{J}$, given by

$$
\mathbf{J}=\left[\begin{array}{ll}
x_{13} & y_{13} \\
x_{23} & y_{23}
\end{array}\right]
$$

$$
\begin{aligned}
\mathbf{J}^{-1} & =\frac{1}{\operatorname{det} \mathbf{J}}\left[\begin{array}{rr}
y_{23} & -y_{13} \\
-x_{23} & x_{13}
\end{array}\right] \\
\operatorname{det} \mathbf{J} & =x_{13} y_{23}-x_{23} y_{13}
\end{aligned}
$$

From the knowledge of the area of the triangle, it can be seen that the magnitude of det $J$ is twice the area of the triangle

$$
A=\frac{1}{2}|\operatorname{det} \mathbf{J}|
$$

$$
\begin{aligned}
& \text { When Matrix } A=\left[\begin{array}{ll}
a & b \\
c & d
\end{array}\right] \\
& A^{-1}=\frac{1}{a d-b c}\left[\begin{array}{cc}
d & -b \\
-c & a
\end{array}\right]
\end{aligned}
$$



## Example 6.2

Determine the Jacobian of the transformation $\mathbf{J}$ for the triangular element shown in Fig. E6.1.


$$
\mathbf{J}=\left[\begin{array}{ll}
x_{13} & y_{13} \\
x_{23} & y_{23}
\end{array}\right]=\left[\begin{array}{rr}
-2.5 & -5.0 \\
3.0 & -3.5
\end{array}\right]
$$

Thus, $\operatorname{det} \mathbf{J}=23.75$ units. This is twice the area of the triangle.

$$
\left\{\begin{array}{l}
\frac{\partial u}{\partial x} \\
\frac{\partial u}{\partial y}
\end{array}\right\}=\mathbf{J}^{-1}\left\{\begin{array}{l}
\frac{\partial u}{\partial \xi} \\
\frac{\partial u}{\partial \eta}
\end{array}\right\} \quad \begin{array}{r}
\mathbf{J}^{-1}=\frac{1}{\operatorname{det} \mathbf{J}}\left[\begin{array}{cc}
y_{23} & -y_{13} \\
-x_{23} & x_{13}
\end{array}\right] \\
\operatorname{det} \mathbf{J}=x_{13} y_{23}-x_{23} y_{13}
\end{array}
$$

$$
\left\{\begin{array}{l}
\frac{\partial u}{\partial x} \\
\frac{\partial u}{\partial y}
\end{array}\right\}=\frac{1}{\operatorname{det} \mathbf{J}}\left\{\begin{array}{c}
y_{23} \frac{\partial u}{\partial \xi}-y_{13} \frac{\partial u}{\partial \eta} \\
-x_{23} \frac{\partial u}{\partial \xi}+x_{13} \frac{\partial u}{\partial \eta}
\end{array}\right\}
$$

Replacing $u$ by the displacement $v$, we get a similar expression

$$
\left\{\begin{array}{l}
\frac{\partial v}{\partial x} \\
\frac{\partial v}{\partial y}
\end{array}\right\}=\frac{1}{\operatorname{det} \mathbf{J}}\left\{\begin{array}{c}
y_{23} \frac{\partial v}{\partial \xi}-y_{13} \frac{\partial v}{\partial \eta} \\
-x_{23} \frac{\partial v}{\partial \xi}+x_{13} \frac{\partial v}{\partial \eta}
\end{array}\right\} \quad \begin{aligned}
& \mathbf{u}=\mathbf{u}_{13} \xi+\mathbf{u}_{23} \eta+\mathbf{u}_{3} \\
& \mathbf{v}=\mathbf{v}_{13} \xi+\mathbf{v}_{23} \eta+\mathbf{v}_{3}
\end{aligned}
$$

Using the strain-displacement relations

$$
\begin{aligned}
\boldsymbol{\epsilon}=\left\{\begin{array}{l}
\left.\begin{array}{l}
\frac{\partial u}{\partial x} \\
\frac{\partial v}{\partial y} \\
\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}
\end{array}\right\} \\
\\
\end{array}\right\}=\frac{1}{\operatorname{det} \mathbf{J}}\left\{\begin{array}{l}
y_{23}(\mathrm{u} 1-\mathrm{u} 3)-y_{13}(\mathrm{u} 2-\mathrm{u} 3) \\
-x_{23}(\mathrm{v} 1-\mathrm{v} 3)+x_{13}(\mathrm{v} 2-\mathrm{v} 3) \\
-x_{23}(\mathrm{u} 1-\mathrm{u} 3)+x_{13}(\mathrm{u} 2-\mathrm{u} 3)+y_{23}(\mathrm{v} 1-\mathrm{v} 3)-y_{13}(\mathrm{v} 2-\mathrm{v} 3)
\end{array}\right\}
\end{aligned}
$$

$$
\text { we can write } y_{31}=-y_{13} \text { and } y_{12}=y_{13}-y_{23}
$$


where $\mathbf{B}$ is a $(3 \times 6)$ element strain-displacement matrix relating the three strains to the six nodal displacements and is given by

$$
\boldsymbol{\sigma}=\mathbf{D} \boldsymbol{\epsilon}
$$

## Example 6.3

Find the strain-nodal displacement matrices $\mathbf{B}^{e}$ for the elements shown in Fig. E6.3. Use local numbers given at the corners.


FIGURE E6.3

Solution We have

## Remember,

$$
\mathbf{B}^{1}=\frac{1}{\operatorname{det} \mathbf{J}}\left[\begin{array}{cccccc}
y_{23} & 0 & y_{31} & 0 & y_{12} & 0 \\
0 & x_{32} & 0 & x_{13} & 0 & x_{21} \\
x_{32} & y_{23} & x_{13} & y_{31} & x_{21} & y_{12}
\end{array}\right]
$$

$$
y 23=y 2-y 3
$$

$$
x 13=x 1-x 3
$$

Taking origin at 1
$\operatorname{det} \mathbf{J}$ is obtained from $x_{13} y_{23}-x_{23} y_{13}=(3)(2)-(3)(0)=6$.

$$
\begin{aligned}
&=\frac{1}{6}\left[\begin{array}{rrrrrr}
2 & 0 & 0 & 0 & -2 & 0 \\
0 & -3 & 0 & 3 & 0 & 0 \\
-3 & 2 & 3 & 0 & 0 & -2
\end{array}\right] \\
& \mathbf{B}^{2}=\frac{1}{6}\left[\begin{array}{rrrrrr}
-2 & 0 & 0 & 0 & 2 & 0 \\
0 & 3 & 0 & -3 & 0 & 0 \\
3 & -2 & -3 & 0 & 0 & 2
\end{array}\right]
\end{aligned}
$$

## Potential Energy Approach

General expression for total potential energy in an elastically loaded structure is,

$$
\pi=\frac{1}{2} \int_{V} \sigma^{T} \varepsilon d V-\int_{V} u^{T} f d v-\int_{A} u^{T} \mathbf{T} d A-\sum_{i} u_{i}^{T} P_{i}
$$

$$
\text { TPE }(T) \text { = Strain Energy + External Work done }
$$

$$
\boldsymbol{\sigma}=\mathbf{D} \boldsymbol{\epsilon}
$$

For plane problems,


$$
\Pi=\frac{1}{2} \int_{A} \boldsymbol{\epsilon}^{\mathrm{T}} \mathbf{D} \boldsymbol{\epsilon} t d A-\int_{A} \mathbf{u}^{\mathrm{T}} \mathbf{f} t d A-\int_{L} \mathbf{u}^{\mathrm{T}} \mathbf{T} t d \ell-\sum_{i} \mathbf{u}_{i}^{\mathrm{T}} \mathbf{P}_{i}
$$

$$
\begin{aligned}
& \Pi=\frac{1}{2} \int_{A} \epsilon^{\mathrm{T}} \mathbf{D} \boldsymbol{\epsilon} t d A-\int_{A} \mathbf{u}^{\mathrm{T}} \mathbf{f} t d A-\int_{L} \mathbf{u}^{\mathrm{T}} \mathbf{T} t d \ell-\sum_{i} \mathbf{u}_{i}^{\mathrm{T}} \mathbf{P}_{i} \\
& \Pi=\sum_{e} \frac{1}{2} \int_{e} \boldsymbol{\epsilon}^{\mathrm{T}} \mathbf{D} \boldsymbol{\epsilon} t d A-\sum_{e} \int_{e} \mathbf{u}^{\mathrm{T}} \mathbf{f} t d A-\int_{L} \mathbf{u}^{\mathrm{T}} \mathbf{T} t d \ell-\sum_{i} \mathbf{u}_{i}^{\mathrm{T}} \mathbf{P}_{i} \\
& \Pi=\sum_{e} U_{e}-\sum_{e} \int_{e} \mathbf{u}^{\mathrm{T}} \mathbf{f} t d A-\int_{L} \mathbf{u}^{\mathrm{T}} \mathbf{T} t d \ell-\sum_{i} \mathbf{u}_{i}^{T} \mathbf{P}_{i}
\end{aligned}
$$

where $U_{e}=\frac{1}{2} \int_{e} \epsilon^{\mathrm{T}} \mathbf{D} \boldsymbol{\epsilon} t d A$ is the element strain energy.

$$
\begin{aligned}
U_{e} & =\frac{1}{2} \int_{e} \mathbf{\epsilon}^{\mathrm{T}} \mathbf{D} \boldsymbol{\epsilon} t d A \quad \boldsymbol{\epsilon}=\mathbf{B} \mathbf{q}=\mathbf{B} \mathbf{U} \\
& =\frac{1}{2} \int_{e} \mathbf{q}^{\mathrm{T}} \mathbf{B}^{\mathrm{T}} \mathbf{D B} \mathbf{q} t d A
\end{aligned}
$$

## Element Stiffness, Ke

$$
\begin{aligned}
U_{e} & =\frac{1}{2} \int_{e} \boldsymbol{\epsilon}^{\mathrm{T}} \mathbf{D} \boldsymbol{\epsilon} t d A \\
& =\frac{1}{2} \int_{e} \mathbf{q}^{\mathrm{T}} \mathbf{B}^{\mathrm{T}} \mathbf{D B} \mathbf{q} t d A
\end{aligned}
$$

Taking the element thickness ( te ) as constant over the element and remembering that all terms in the $\mathbf{D}$ and $\mathbf{B}$ matrices are constants, we get

$$
\begin{aligned}
& U_{e}=\frac{1}{2} \mathbf{q}^{\mathrm{T}} \mathbf{B}^{\mathrm{T}} \mathbf{D B} t_{e}\left(\int_{e} d A\right) \mathbf{q} \\
& U_{e}=\frac{1}{2} \mathbf{q}^{\mathrm{T}} t_{e} A_{e} \mathbf{B}^{\mathrm{T}} \mathbf{D B} \mathbf{q}
\end{aligned}
$$

$$
\int_{c} d A=A_{c} \text { where } A_{e} \text { is the area of the element. }
$$

$$
\begin{array}{r}
U_{e}=\frac{1}{2} \mathbf{q}^{\mathrm{T}} \mathbf{k}^{e} \mathbf{q} \quad \text { where } \mathbf{k}^{e} \text { is the element stiffness matrix given by } \\
\qquad \mathbf{k}^{e}=t_{e} A_{e} \mathbf{B}^{\mathrm{T}} \mathbf{D B}
\end{array}
$$

$$
\begin{aligned}
& U_{e}=\frac{1}{2} \mathbf{q}^{\mathrm{T}} \mathbf{k}^{e} \mathbf{q} \\
& \qquad U=\sum_{e} \frac{1}{2} \mathbf{q}^{\mathrm{T}} \mathbf{k}^{e} \mathbf{q}=\frac{1}{2} \mathbf{Q}^{\mathrm{T}} \mathbf{K} \mathbf{Q}=\frac{1}{2} U^{T} K U
\end{aligned}
$$

## Force Terms

The body force term $\int_{e} \mathbf{u}^{\mathrm{T}} \mathbf{f} t d A$

$$
\int_{e} \mathbf{u}^{\mathrm{T}} \mathbf{f} t d A=t_{e} \int_{e}\left(u f_{x}+v f_{y}\right) d A
$$

Using the interpolation relations
Substitute,

$$
\begin{aligned}
& u=N_{1} u_{1}+N_{2} u_{2}+N_{3} u_{3} \\
& v=N_{1} v_{1}+N_{2} v_{2}+N_{3} v_{3}
\end{aligned}
$$

$$
\begin{aligned}
\int_{e} \mathbf{u}^{\mathrm{T}} \mathbf{f} t d A= & \mathbf{u} 1\left(t_{e} f_{x} \int_{e} N_{1} d A\right)+\mathbf{v} 1\left(t_{e} f_{y} \int_{e} N_{1} d A\right) \\
& +\mathbf{u} \mathbf{2}\left(t_{e} f_{x} \int_{e} N_{2} d A\right)+\mathbf{v} \mathbf{2}\left(t_{e} f_{y} \int_{e} N_{2} d A\right) \\
& +\mathbf{u} 3\left(t_{e} f_{x} \int_{e} N_{3} d A\right)+\mathbf{v} \mathbf{3}\left(t_{e} f_{y} \int_{e} N_{3} d A\right)
\end{aligned}
$$

$$
\int_{e} \mathbf{u}^{\mathrm{T}} \mathbf{f} t d A=\mathbf{q}^{\mathrm{T}} \mathbf{f}^{e}
$$

$$
\int_{e} N_{1} d A=\frac{1}{3} A_{e} h=\frac{1}{3} A_{e}
$$

$$
\mathbf{f}^{e}=\frac{t_{e} A_{e}}{3}\left[\begin{array}{llllll}
f_{x} & f_{y} & f_{x} & f_{y} & f_{x} & f_{y}
\end{array}\right]^{\mathrm{T}}
$$

$$
\int_{e} N_{2} d A=\int_{e} N_{3} d A=1 / 3 A_{e}
$$

Traction force vector

$$
\mathbf{T}^{e}=\frac{t_{e} \ell_{1-2}}{6}\left[2 T_{x_{1}}+T_{x_{2}} \quad 2 T_{y_{1}}+T_{y_{2}} \quad T_{x_{1}}+2 T_{x_{2}} \quad T_{y_{1}}+2 T_{y_{2}}\right]^{\mathrm{T}}
$$

$$
\Pi=\frac{1}{2} \int_{A} \epsilon^{\mathrm{T}} \mathbf{D} \boldsymbol{\epsilon} t d A-\int_{A} \mathbf{u}^{\mathrm{T}} \mathbf{f} t d A-\int_{L} \mathbf{u}^{\mathrm{T}} \mathbf{T} t d \ell-\sum_{i} \mathbf{u}_{i}^{\mathrm{T}} \mathbf{P}_{i}
$$

$$
\Pi=\frac{1}{2} \mathbf{Q}^{\mathrm{T}} \mathbf{K} \mathbf{Q}-\mathbf{Q}^{\mathrm{T} \mathbf{f}}
$$

$$
K \mathbf{K}=\mathbf{F} \quad \text { (Or) } \quad \mathbf{K} \mathbf{U}=\mathbf{F}
$$

Q. Assuming plane stress condition, evaluate stiffness matrix for the element shown in figure. Assume E $=200 \mathrm{GPa}$, Poisson's ratio 0.3.


Coordinates are in mm
Figure
Q. Assuming plane stress condition, evaluate stiffness matrix for the element shown in figure. Assume E $=200 \mathrm{GPa}$, Poisson's ratio 0.3.

Coordinates are in mm


Figure
$\mathbf{B}=\frac{1}{\operatorname{det} \mathbf{J}}\left[\begin{array}{cccccc}y_{23} & 0 & y_{31} & 0 & y_{12} & 0 \\ 0 & x_{32} & 0 & x_{13} & 0 & x_{21} \\ x_{32} & y_{23} & x_{13} & y_{31} & x_{21} & y_{12}\end{array}\right]$
Plane stress: $\quad \mathbf{D}=\frac{E}{1-v^{2}}\left(\begin{array}{ccc}1 & v & 0 \\ v & 1 & 0 \\ 0 & 0 & \frac{1-v}{2}\end{array}\right)$

Plane strain:

$$
\mathbf{D}=\frac{E}{(1+v)(1-2 v)}\left(\begin{array}{ccc}
1-v & v & 0 \\
v & 1-v & 0 \\
0 & 0 & \frac{1-2 v}{2}
\end{array}\right)
$$

$\mathbf{k}^{e}$ is the element stiffness matrix given by

$$
\mathbf{k}^{e}=t_{e} A_{e} \mathbf{B}^{\mathrm{T}} \mathbf{D B}
$$

Given that,
Young's modulus, $E=200 \mathrm{GPa}=200 \times 10^{3} \mathrm{~N} / \mathrm{mm}^{2}$
Poisson's ratio, $\mu=0.3$
$(0,1)$


Global coordinates of,
Node-1, $\left(x_{1}, y_{1}\right)=(0,-1)$
Node-2, $\left(x_{2}, y_{2}\right)=(2,0)$
Node-3, $\left(x_{3}, y_{3}\right)=(0,1)$

## Assume,

Thickness, $t=1 \mathrm{~mm}$
And, given coordinate are in mm .

Stiffness matrix for linear triangular element,

$$
[K]=[B]^{T}[D][B] \text { A.t }
$$

Stress-strain relationship matrix, considering plane stress condition,

$$
[D]=\frac{E}{1-\mu^{2}}\left[\begin{array}{ccc}
1 & \mu & 0 \\
\mu & 1 & 0 \\
0 & 0 & \frac{1-\mu}{2}
\end{array}\right]=\frac{200 \times 10^{3}}{1-0.3^{2}}\left[\begin{array}{ccc}
1 & 0.3 & 0 \\
0.3 & 1 & 0 \\
0 & 0 & \frac{1-0.3}{2}
\end{array}\right]
$$

$$
[D]=219.78 \times 10^{3}\left[\begin{array}{ccc}
1 & 0.3 & 0 \\
0.3 & 1 & 0 \\
0 & 0 & 0.35
\end{array}\right]
$$

$$
\begin{aligned}
& \mathbf{J}=\left[\begin{array}{ll}
x_{13} & y_{13} \\
x_{23} & y_{23}
\end{array}\right] \\
& \operatorname{det} \mathbf{J}=x_{13} y_{23}-x_{23} y_{13}
\end{aligned}
$$

## $A=\frac{1}{2}|\operatorname{det} \mathbf{J}|$

$$
\begin{aligned}
& x_{3}-x_{2}=0-2=-2 \\
& x_{1}-x_{3}=0-0=0 \\
& x_{2}-x_{1}=2 \times 0=2 \\
& y_{2}-y_{3}=0-1=-1 \\
& y_{3}-y_{1}=1-(-1)=2 \\
& y_{1}-y_{2}=(-1)-0=-1
\end{aligned}
$$

$$
\mathbf{B}=\frac{1}{\operatorname{det} \mathbf{J}}\left[\begin{array}{cccccc}
y_{23} & 0 & y_{31} & 0 & y_{12} & 0 \\
0 & x_{32} & 0 & x_{13} & 0 & x_{21} \\
x_{32} & y_{23} & x_{13} & y_{31} & x_{21} & y_{12}
\end{array}\right]
$$

$$
[B]=\frac{1}{4}\left[\begin{array}{cccccc}
-1 & 0 & 2 & 0 & -1 & 0 \\
0 & -2 & 0 & 0 & 0 & 2 \\
-2 & -1 & 0 & 2 & 2 & -1
\end{array}\right] \quad[B]^{T}=\frac{1}{4}\left[\begin{array}{ccc}
-1 & 0 & -2 \\
0 & -2 & -1 \\
2 & 0 & 0 \\
0 & 0 & 2 \\
-1 & 0 & 2 \\
0 & 2 & -1
\end{array}\right]
$$

Stiffness matrix for linear triangular element,

$$
[K]=[B]^{T}[D][B] \text { A.t }
$$

$$
[K]=\frac{1}{4}\left[\begin{array}{ccc}
-1 & 0 & -2 \\
0 & -2 & -1 \\
2 & 0 & 0 \\
0 & 0 & 2 \\
-1 & 0 & 2 \\
0 & 2 & -1
\end{array}\right] \times 219.78 \times 10^{3}\left[\begin{array}{ccc}
1 & 0.3 & 0 \\
0.3 & 1 & 0 \\
0 & 0 & 0.35
\end{array}\right] \times \frac{1}{4}\left[\begin{array}{cccccc}
-1 & 0 & 2 & 0 & -1 & 0 \\
0 & -2 & 0 & 0 & 0 & 2 \\
-2 & -1 & 0 & 2 & 2 & -1
\end{array}\right] \times 2 \times 1
$$

$$
\begin{aligned}
& =27472.5\left[\begin{array}{ccc}
-1 & 0 & -2 \\
0 & -2 & -1 \\
2 & 0 & 0 \\
0 & 0 & 2 \\
-1 & 0 & 2 \\
0 & 2 & -1
\end{array} \left\lvert\,\left[\begin{array}{ccc}
1 & 0.3 & 0 \\
0.3 & 1 & 0 \\
0 & 0 & 0.35
\end{array}\right]\left[\begin{array}{cccccc}
-1 & 0 & 2 & 0 & -1 & 0 \\
0 & -2 & 0 & 0 & 0 & 2 \\
-2 & -1 & 0 & 2 & 2 & -1
\end{array}\right]\right.\right. \\
& =27472.5\left[\begin{array}{ccc}
-1 & -0.3 & -0.7 \\
-0.6 & -2 & -0.35 \\
2 & 0.6 & 0 \\
0 & 0 & 0.7 \\
-1 & -0.3 & 0.7 \\
0.6 & 2 & -0.35
\end{array}\right]\left[\begin{array}{llllll}
-1 & 0 & 2 & 0 & -1 & 0 \\
0 & -2 & 0 & 0 & 0 & 2 \\
-2 & -1 & 0 & 2 & 2 & -1
\end{array}\right] \\
& {[K]=27472.5\left[\begin{array}{cccccc}
2.4 & 1.3 & -2 & -1.4 & -0.4 & 0.1 \\
1.3 & 4.35 & -1.2 & -0.7 & -0.1 & -3.65 \\
-2 & -1.2 & 4 & 0 & -2 & 1.2 \\
-1.4 & -0.7 & 0 & 1.4 & 1.4 & -0.7 \\
-0.4 & -0.1 & -2 & 1.4 & 2.4 & -1.3 \\
0.1 & -3.65 & 1.2 & -0.7 & -1.3 & 4.35
\end{array}\right]}
\end{aligned}
$$

Q. It is required to determine the transverse displacement and the stresses induced in the plate shown in figure using a one-element idealization. Determine the constitutive matrix and the strain-displacement matrix and hence the stiffness matrix and the load vector. Assume $\mathrm{E}=205 \mathrm{GPa}, \mu=0.33$, and $\mathrm{t}=10 \mathrm{~mm}$.


Figure

## Given that,

Young's modulus, $E=205 \mathrm{GPa}=205 \times 10^{3} \mathrm{~N} / \mathrm{mm}^{2}$
Poisson's ratio, $\mu=0.33$
Thickness, $t=10 \mathrm{~mm}$


Coordinates of,

$$
\begin{aligned}
& \text { Node-1, }\left(x_{1}, y_{1}\right)=(0,0) \\
& \text { Node-2, }\left(x_{2}, y_{2}\right)=(50,20) \\
& \text { Node-3, }\left(x_{3}, y_{3}\right)=(0,40)
\end{aligned}
$$

## Transverse Displacement

Nodal Displacement can be obtained by using the relation,

$$
[K]\{\delta\}=\{F\}
$$

Stiffness matrix for linear triangular element,

$$
[K]=[B]^{T}[D][B] \text { A.t }
$$

$$
\mathbf{J}=\left[\begin{array}{ll}
x_{13} & y_{13} \\
x_{23} & y_{23}
\end{array}\right]
$$

Area of triangle,

$$
A=\frac{1}{2}\left|\begin{array}{lll}
1 & x_{1} & y_{1} \\
1 & x_{2} & y_{2} \\
1 & x_{3} & y_{3}
\end{array}\right| \text { or }
$$

$$
\operatorname{det} \mathbf{J}=x_{13} y_{23}-x_{23} y_{13}
$$

$$
\begin{aligned}
A & =\frac{1}{2}|\operatorname{det} \mathbf{J}| \\
& =\frac{1}{2}[2000]
\end{aligned}
$$

$$
\begin{aligned}
{[D] } & =\frac{E}{1-\mu^{2}}\left[\begin{array}{ccc}
1 & \mu & 0 \\
\mu & 1 & 0 \\
0 & 0 & \frac{1-\mu}{2}
\end{array}\right] \\
& =\frac{205 \times 10^{3}}{1-0.33^{2}}\left[\begin{array}{ccc}
1 & 0.33 & 0 \\
0.33 & 1 & 0 \\
0 & 0 & \frac{1-0.33}{2}
\end{array}\right]
\end{aligned}
$$

$[D]=230.052 \times 10^{3}\left[\begin{array}{ccc}1 & 0.33 & 0 \\ 0.33 & 1 & 0 \\ 0 & 0 & 0.335\end{array}\right]$
$\mathbf{B}=\frac{1}{\operatorname{det} \mathbf{J}}\left[\begin{array}{cccccc}y_{23} & 0 & y_{31} & 0 & y_{12} & 0 \\ 0 & x_{32} & 0 & x_{13} & 0 & x_{21} \\ x_{32} & y_{23} & x_{13} & y_{31} & x_{21} & y_{12}\end{array}\right]$

$$
\begin{aligned}
& y_{2}-y_{3}=20-40=-20 \\
& y_{3}-y_{1}=40-0=40 \\
& y_{1}-y_{2}=0-20=-20
\end{aligned}
$$

Then,

$$
x_{3}-x_{2}=0-50=-50
$$

$$
\begin{array}{ll}
{[B]=\frac{1}{2 \times 1000}\left[\begin{array}{cccccc}
-20 & 0 & 40 & 0 & -20 & 0 \\
0 & -50 & 0 & 0 & 0 & 50 \\
-50 & -20 & 0 & 40 & 50 & -20
\end{array}\right]} & \left.\begin{array}{l}
x_{1}-x_{3}=0-0=0 \\
x_{2}-x_{1}=50-0=50
\end{array}\right] \\
{[B]=\frac{1}{2000}\left[\begin{array}{cccccc}
-20 & 0 & 40 & 0 & -20 & 0 \\
0 & -50 & 0 & 0 & 0 & 50 \\
-50 & -20 & 0 & 40 & 50 & -20
\end{array}\right]} & \text { And, } \\
{[B]^{T}=\frac{1}{2000}\left[\begin{array}{ccc}
-20 & 0 & -50 \\
0 & -50 & -20 \\
40 & 0 & 0 \\
0 & 0 & 40 \\
-20 & 0 & 50 \\
0 & 50 & -20
\end{array}\right]}
\end{array}
$$

$$
\begin{aligned}
{[K] } & =\frac{1}{2000}\left[\begin{array}{ccc}
-20 & 0 & -50 \\
0 & -50 & -50 \\
40 & 0 & 0 \\
0 & 0 & 40 \\
-20 & 0 & 50 \\
0 & 50 & -20
\end{array}\right] \times 230.052 \times 10^{3}\left[\begin{array}{ccccc}
1 & 0.33 & 0 \\
0.33 & 1 & 0 \\
0 & 0 & 0.335
\end{array}\right] \\
& \times \frac{1}{2000}\left[\begin{array}{ccccc}
-20 & 0 & 40 & 0 & -20 \\
0 & -50 & 0 & 0 & 0 \\
-50 & -20 & 0 & 40 & 50 \\
-20
\end{array}\right] \times 1000 \times 10 \\
& =575.13\left[\begin{array}{ccc}
-20 & 0 & -50 \\
0 & -50 & -20 \\
40 & 0 & 0 \\
0 & 0 & 40 \\
-20 & 0 & 50 \\
0 & 50 & -20
\end{array}\right]
\end{aligned}
$$

$$
\begin{aligned}
& =575.13\left[\begin{array}{ccc}
-20 & 0 & -50 \\
0 & -50 & -20 \\
40 & 0 & 0 \\
0 & 0 & 40 \\
-20 & 0 & 50 \\
0 & 50 & -20
\end{array}\right]\left[\begin{array}{ccccc}
-20 & -16.5 & 40 & 0 & -20 \\
16.5 \\
-6.6 & -50 & 13.2 & 0 & -6.6 \\
-16.75 & -6.7 & 0 & 13.4 & 16.75 \\
-6.7
\end{array}\right] \\
{[K] } & =575.13\left[\begin{array}{cccccc}
1237.5 & 665 & -800 & -670 & -437.5 & 5 \\
665 & 2634 & -660 & -268 & -5 & -2366 \\
-800 & -660 & 1600 & 0 & -800 & 660 \\
-670 & -268 & 0 & 536 & 670 & -268 \\
-437.5 & -5 & -800 & 670 & 1237.5 & -665 \\
5 & -2366 & 660 & -268 & -665 & -2634
\end{array}\right]
\end{aligned}
$$

Verification: Evaluated stiffness matrix is said to be correct if it satisfies two conditions such as it should be symmetric and sum of values of any row or column should be zero. As these conditions are satisfied by obtained stiffness matrix, it is said to be correct.

$575.13\left[\begin{array}{cccccc}1237.5 & 665 & -800 & -670 & -437.5 & 5 \\ 665 & 2634 & -660 & -268 & -5 & -2366 \\ -800 & -660 & 1600 & 0 & -800 & 660 \\ -670 & -268 & 0 & 536 & 670 & -268 \\ -437.5 & -5 & -800 & 670 & 1237.5 & -665 \\ 5 & -2366 & 660 & -268 & -665 & 2364\end{array}\right]\left\{\begin{array}{l}u_{1} \\ v_{1} \\ u_{2} \\ v_{2} \\ u_{3} \\ u_{3}\end{array}\right\}=\left\{\begin{array}{l}\mathbf{0} \\ \mathbf{0} \\ \mathbf{0} \\ \mathbf{- 1 0 0 0} \\ \mathbf{0} \\ \mathbf{0}\end{array}\right\}$
Applying boundary conditions,

$$
u_{1}=v_{1}=u_{3}=v_{3}=0
$$

$575.13\left[\begin{array}{cc}1600 & 0 \\ 0 & 536\end{array}\right]\left[\begin{array}{l}u_{2} \\ v_{2}\end{array}\right\}=\left\{\begin{array}{c}0 \\ -1000\end{array}\right\}$
On writing above matrix in equation form,

$$
\begin{aligned}
575.13 \times 1600 u_{2} & =0 \\
u_{2} & =0 \mathrm{~mm}
\end{aligned}
$$

Also,

$$
\begin{aligned}
575.13 \times 536 v_{2} & =-1000 \\
v_{2} & =\frac{-1000}{575.13 \times 536} \\
v_{2} & =-3.244 \times 10^{-3} \mathrm{~mm} \quad \text { (' }- \text { ve' sign indicates downward displacement) }
\end{aligned}
$$

## Stresses Induced in the plate

Stress vector for linear triangular element,

$$
\{\sigma\}=[D][B]\{\delta\}
$$

$$
\{\sigma\}=230.052 \times 10^{3}\left[\begin{array}{ccc}
1 & 0.33 & 0 \\
0.33 & 1 & 0 \\
0 & 0 & 0.335
\end{array}\right] \times \frac{1}{2000}\left[\begin{array}{cccccc}
-20 & 0 & 40 & 0 & -20 & 0 \\
0 & -50 & 0 & 0 & 0 & 50 \\
-50 & -20 & 0 & 40 & 50 & -20
\end{array}\right] \times 10^{-3}\left[\begin{array}{c}
0 \\
0 \\
0 \\
-3.244 \\
0 \\
0
\end{array}\right\}
$$

$$
\begin{aligned}
& \left\{\begin{array}{c}
\sigma_{x} \\
\sigma_{y} \\
\tau_{x y}
\end{array}\right\}=0.115\left[\begin{array}{ccc}
1 & 0.33 & 0 \\
0.33 & 1 & 0 \\
0 & 0 & 0.335
\end{array}\right]\left[\begin{array}{cccccc}
-20 & 0 & 40 & 0 & -20 & 0 \\
0 & -50 & 0 & 0 & 0 & 50 \\
-50 & -20 & 0 & 40 & 50 & -20
\end{array}\right]\left[\begin{array}{c}
0 \\
0 \\
0 \\
-3.244 \\
0 \\
0
\end{array}\right\} \\
& \left\{\begin{array}{c}
\sigma_{x} \\
\sigma_{y} \\
\tau_{x y}
\end{array}\right\}=0.115\left[\begin{array}{cccccc}
-20 & -16.5 & 40 & 0 & -20 & 16.5 \\
-6.6 & -50 & 13.2 & 0 & -6.6 & 50 \\
-16.75 & -6.7 & 0 & 13.4 & 16.75 & -6.7
\end{array}\right]\left\{\begin{array}{c}
0 \\
0 \\
0 \\
-3.244 \\
0 \\
0
\end{array}\right\} \\
& \left\{\begin{array}{c}
\sigma_{x} \\
\sigma_{y} \\
\tau_{x y}
\end{array}\right\}=0.115\left\{\begin{array}{c}
0 \\
0 \\
-43.47
\end{array}\right\}
\end{aligned}
$$

$$
\left\{\begin{array}{c}
\sigma_{x} \\
\sigma_{y} \\
\tau_{x y}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
0 \\
-4.999
\end{array}\right\} \mathrm{N} / \mathrm{mm}^{2}
$$

Therefore,

Normal stress in $x$-direction, $\sigma_{x}=0 \mathrm{~N} / \mathrm{mm}^{2}$
Normal stress in $y$-direction, $\sigma_{y}=0 \mathrm{~N} / \mathrm{mm}^{2}$
Shear stress in $x-y$ plane, $\tau_{x y}=-4.999 \mathrm{~N} / \mathrm{mm}^{2}$


Figure: Nodal Displacements
Figure


Figure



Figure: Nodal Forces

## Axisymmetric Solids Subjected to Axisymmetric Loading

The problem is said to be axisymmetric type, if the object has an axis of symmetry and parameters such as geometry boundary conditions, loading and materials are symmetric about this axis. Axisymmetric solids are also known as solids of revolutions. Analysis of such problems is termed as axisymmetric analysis.



Rigid shaft



Figure (3): Hollow Cylinder Subjected to Internal Pressure


Figure (2)
coordinate system are $r, \theta, \mathrm{z}$ and $u, v, w$ are their respective displacement functions.

Strain occurring in the element,

In axisymmetric problem, parameters such as surface loading and geometry are independent of the circumferential direction ' $\theta$ '. Thereby, displacement in circumferential direction ' $v$ ' will be zero. Thus, only displacements corresponding to direction ' $r$ ' and ' $z$ ' remains.

$$
\text { i.e., } \begin{aligned}
u & =f(r, z) \\
v & =0 \\
w & =f(r, z)
\end{aligned}
$$

$$
\begin{aligned}
& \{\varepsilon\}=\left\{\begin{array}{l}
\varepsilon_{r} \\
\varepsilon_{\theta} \\
\varepsilon_{z} \\
\gamma_{r z}
\end{array}\right\} \\
& \{\varepsilon\}=\left\{\begin{array}{c}
\frac{\partial u}{\partial r} \\
\frac{u}{r} \\
\frac{\partial w}{\partial z} \\
\frac{\partial u}{\partial z}+\frac{\partial w}{\partial r}
\end{array}\right\}
\end{aligned}
$$

$\{\sigma\}=\left\{\begin{array}{c}\sigma_{r} \\ \sigma_{\theta} \\ \sigma_{z} \\ \tau_{r r}\end{array}\right\}$
Where,
$\sigma_{r}-$ Radial stress
$\sigma_{\theta}$ - Circumferential or tangential stress
$\sigma_{z}$-Longitudinal or axial stress

$$
\left\{\begin{array}{c}
\sigma_{r} \\
\sigma_{\theta} \\
\sigma_{z} \\
\tau_{r z}
\end{array}\right\}=\frac{E}{(1+\mu)(1-2 \mu)}\left[\begin{array}{cccc}
1-\mu & \mu & \mu & 0 \\
\mu & 1-\mu & \mu & 0 \\
\mu & \mu & 1-\mu & 0 \\
0 & 0 & 0 & \frac{1-2 \mu}{2}
\end{array}\right]\left\{\begin{array}{c}
\varepsilon_{r} \\
\varepsilon_{\theta} \\
\varepsilon_{z} \\
\gamma_{r z}
\end{array}\right\}
$$

$\tau_{r z}$ - Shear stress.

$$
\{\sigma\}=[D]\{\varepsilon\}
$$

Using Hooke's law, for stress-strain relationship,

$$
\{\sigma\}=[D]\{\varepsilon\}
$$

Where,
[D] - Stress- strain relationship matrix.

Axisymmetric element:

$$
[D]=\frac{E}{(1+\mu)(1-2 \mu)}
$$

$$
\left[\begin{array}{cccc}
1-\mu & \mu & \mu & 0 \\
\mu & 1-\mu & \mu & 0 \\
\mu & \mu & 1-\mu & 0 \\
0 & 0 & 0 & \frac{1-2 \mu}{2}
\end{array}\right]
$$

Plane stress:

$$
\mathbf{D}=\frac{E}{1-v^{2}}\left(\begin{array}{ccc}
1 & v & 0 \\
v & 1 & 0 \\
0 & 0 & \frac{1-v}{2}
\end{array}\right)
$$

Plane strain:

$$
\mathbf{D}=\frac{E}{(1+v)(1-2 v)}\left(\begin{array}{ccc}
1-v & v & 0 \\
v & 1-v & 0 \\
0 & 0 & \frac{1-2 v}{2}
\end{array}\right)
$$



Figure
Q.. Calculate the stiffness matrix for the axisymmetric element Shown in the figure

Given that,
Coordinates of,

$$
\begin{aligned}
& \text { Node-1, }\left(r_{1}, z_{1}\right)=(10,10) \\
& \text { Node-2, }\left(r_{2}, z_{2}\right)=(30,10) \\
& \text { Node-3, }\left(r_{3}, z_{3}\right)=(20,30)
\end{aligned}
$$

Assume, given coordinates are in mm.

## Constitutive Matrix

Assume,
Young's modulus, $E=2 \times 10^{5} \mathrm{~N} / \mathrm{mm}^{2}$
Poisson's ratio, $\mu=0.3$
Constitutive matrix or stress-strain relationship matrix is given by,

$$
\begin{gathered}
{[D]=\frac{E}{(1+\mu)(1-2 \mu)}\left[\begin{array}{cccc}
1-\mu & \mu & \mu & 0 \\
\mu & 1-\mu & \mu & 0 \\
\mu & \mu & 1-\mu & 0 \\
0 & 0 & 0 & \left(\frac{1-2 \mu}{2}\right)
\end{array}\right]=\frac{2 \times 10^{5}}{(1+0.3)(1-0.6)}\left[\begin{array}{cccc}
1-0.3 & 0.3 & 0.3 & 0 \\
0.3 & 1-0.3 & 0.3 & 0 \\
0.3 & 0.3 & 1-0.3 & 0 \\
0 & 0 & 0 & \left(\frac{1-0.6}{2}\right)
\end{array}\right]} \\
\therefore[D]=38.4 \times 10^{3}\left[\begin{array}{cccc}
7 & 3 & 3 & 0 \\
3 & 7 & 3 & 0 \\
3 & 3 & 7 & 0 \\
0 & 0 & 0 & 2
\end{array}\right]
\end{gathered}
$$

## Higher Order elements


(a) Linear Element $(\mathrm{n}=1)$

Figure(1): Single Order Element

(b) Non-linear Quadratic Element

## Figure(2): Higher Order Element

If the order of interpolation polynomial of an element is two or more than two such an element is called higher order element. They are complex or multiplex elements, whose order is greater than one. These elements consist both primary as well as secondary nodes. Primary nodes include corner nodes while secondary nodes include intemal or mid-point nodes.

## Element Name

Spring, Damper Beam, Truss

2D Elements Surface Element
1D Elements Line Element


Shell, Plane2D

## Element Shape

 First Order Second Order

3D Elements
Volume element


Hexahedral

Tetrahedral


## Isoparametric Element

If the shape or geometry and field displacement variables of the elements are described by the same shape fumctions of the same order, then the elements are known as isoparametric elements. In isoparametric elements, the number of nodes for defining both geometry and displacements are equal $(i=j)$.


Figure: Isoparametric Element

- $\rightarrow$ Nodes for defining displacements
$\square \rightarrow$ Nodes for defining geometry
Two dimensional and three dimensional elasticity problems can be solved using the isoparametric elements.


## The displacements can be written as,

$$
\begin{aligned}
& u=N_{1} u_{1}+N_{2} u_{2}+N_{3} u_{3} \\
& v=N_{1} v_{1}+N_{2} v_{2}+N_{3} v_{3}
\end{aligned}
$$

Similarly we can write the geometric coordinates:

$$
\begin{aligned}
& x=N_{1} x_{1}+N_{2} x_{2}+N_{3} x_{3} \\
& y=N_{1} y_{1}+N_{2} y_{2}+N_{3} y_{3}
\end{aligned}
$$

## Significance of Jacobian Transformation

1. A Jacobian transformation gives the relation between the derivatives in the global and local coordinate systems.
2. It evaluates the deviation of given element from the standard element.
3. It is used to compute the strain displacement matrix.

Most of the problems for which isoparametric elements are employed involves curved boundaries. The finite elements of such problems involves curved sides. These curve sided actual elements are approximated into simple shapes possessing flat surfaces or straight edges. Triangular and quadrilateral elements are the mostly used isoparametric elements.

## 1. Triangular Elements




Figure(1): Two Dimensional Triangular Element
The actual finite triangular element with curved sides, represented in global coordinate system is transformed into straight edged triangular element, represented in natural coordinate system.

## 2. Quadrilateral Elements




Figure(2): Two Dimensional Quadrilateral Element
The actual finite Quadrilateral element with distorted sides, represented in global coordinate system is transformed into straight edged element, represented in natural coordinate system.

Natural co-ordinates are used to specify a point with in the element with the help of dimensionless numbers whose magnitude do not exceed unity.

Difference between isoparametric, Subparametric, Super parametric elements

| Isoparametric Element |  | Subparametric Element |  | Superparametric Element |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 1. | If the shape or geometry and field displacement variables of the elements are described by same shape functions of the same order, then the elements are known as isoparametric elements. | 1. | If the geometry of the elements are described by lower order shape functions compared to field variables (displacements), then the elements are known as subparametric element | 1. | If the geometry of the element is described by higher order shape functions compared to field variables (displacements) then the elements are known as super parametric elements. |
| 2. | The number of nodes for defining both geometry and displacements are equal. | 2. | The number of nodes for defining the displacements is more than the number of nodes for defining geometry | 2. | The number of nodes for defining the displacement is less than the number of nodes for defining geometry |
| 3. | Figure: Isoparametric element | 3. | Figure: Subparametric element | 3. | Figure: Superparametric element |

Consider the general quadrilateral element as shown in figure. The local nodes are numbered as 1,2,3 and 4


Figure (1) : Four-Noded Element (Global Co-ordinate System) Four-node quadrilateral element.

Let $\left(x_{1}, y_{1}\right),\left(x_{2}, y_{2}\right),\left(x_{3}, y_{3}\right),\left(x_{4}, y_{4}\right)$ are the co-ordinates at nodes $1,2,3,4$. $u_{1}, u_{2}, u_{3}, u_{4}$ - Displacements along $x$-axis at nodes $1,2,3,4$.
$v_{1}, v_{2}, v_{3}, v_{4}$ - Displacements along $y$-axis at nodes $1,2,3,4$.,

## The quadrilateral element in $\xi, \eta$ space (the master element).

$$
\begin{aligned}
& u(\xi, \eta)=a_{1}+a_{2} \xi+a_{3} \eta+a_{4} \xi \eta \\
& v(\xi, \eta)=a_{5}+a_{6} \xi+a_{7} \eta+a_{8} \xi \eta
\end{aligned}
$$


$a_{1}, a_{2}, a_{3}, \ldots a_{8}$-Polynomial coefficients.

Figure (2): Isoparametric Quadrilateral Element (Natural Co-ordinate System)


(a) Slave (distorted) element

$$
\begin{array}{ll}
x=x(\xi, \eta) & \xi=\xi(x, y) \\
y=y(\xi, \eta) & \eta=\eta(x, y)
\end{array}
$$

Coordinate

(b) Master (parent) element

Isoparametric coordinate transformation.

$$
u=a_{1}+a_{2} \xi+a_{3} \eta+a_{4} \xi \eta
$$

After substituting the coordinates, we will get the following form

$$
u=\frac{1}{4}(1-\xi)(1-\eta) u_{1}+\frac{1}{4}(1+\xi)(1-\eta) u_{2}+\frac{1}{4}(1+\xi)(1+\eta) u_{3}+\frac{1}{4}(1-\xi)(1+\eta) u_{4}
$$

The above equation can be written as,

$$
u=N_{1} u_{1}+N_{2} u_{2}+N_{3} u_{3}+N_{4} u_{4}
$$

$$
N_{i}=\left\{\begin{array}{l}
1 \text { at node } i \\
0 \text { at all other nodes }
\end{array}\right.
$$

Where,

$$
N_{1}, N_{2}, N_{3}, N_{4} \text { - Shape functions of the isoparametric element }
$$

$$
\begin{array}{l|l|}
N_{1}=\frac{1}{4}(1-\xi)(1-\eta) \\
N_{2}=\frac{1}{4}(1+\xi)(1-\eta) \\
N_{3}=\frac{1}{4}(1+\xi)(1+\eta) & \begin{array}{c}
\text { The Four Node Bilinear Quad: } \\
\text { Shape Function Plot }
\end{array} \\
N_{4}=\frac{1}{4}(1-\xi)(1+\eta) & \begin{array}{l}
N_{1}=1 \\
=0
\end{array} \quad \begin{array}{l}
\text { at node } 1 \\
\text { at nodes 2, 3, and 4 }
\end{array} \\
\hline
\end{array}
$$

$$
\text { Similarly } v=a_{5}+a_{6} \xi+a_{7} \eta+a_{8} \xi \eta \text { becomes } \quad v=N_{1} v_{1}+N_{2} v_{2}+N_{3} v_{3}+N_{4} v_{4}
$$

The displacements are

$$
\begin{aligned}
& u=N_{1} u_{1}+N_{2} u_{2}+N_{3} u_{3}+N_{4} u_{4} \\
& v=N_{1} v_{1}+N_{2} v_{2}+N_{3} v_{3}+N_{4} v_{4}
\end{aligned}
$$



$$
=[N]\{\delta\}
$$

$[N]$ - Shape function matrix
$\{\delta\}$ - Nodal displacement vector

Displacements of any point $P$, inside the quadratic element

$$
\left\{\begin{array}{l}
u \\
v
\end{array}\right\}=\left[\begin{array}{cccccccc}
N_{1} & 0 & N_{2} & 0 & N_{3} & 0 & N_{4} & 0 \\
0 & N_{1} & 0 & N_{2} & 0 & N_{3} & 0 & N_{4}
\end{array}\right]\left\{\begin{array}{l}
u_{2} \\
v_{2} \\
u_{3} \\
v_{3} \\
u_{4} \\
v_{4}
\end{array}\right\}
$$



Using Isoparametric representation, we can write the geometric coordinates of Point "P"

$$
\left\{\begin{array}{l}
x \\
y
\end{array}\right\}=\left[\begin{array}{cccccccc}
N_{1} & 0 & N_{2} & 0 & N_{3} & 0 & N_{4} & 0 \\
0 & N_{1} & 0 & N_{2} & 0 & N_{3} & 0 & N_{4}
\end{array}\right]\left\{\begin{array}{l}
y_{1} \\
x_{2} \\
y_{2} \\
x_{3} \\
y_{3} \\
x_{4} \\
y_{4}
\end{array}\right\}
$$

Let,

$$
\begin{aligned}
& f=f(x, y) \\
& f=f[x(\xi, \eta), y(\xi, \eta)]
\end{aligned}
$$

By chain rule of partial differentiation,

$$
\begin{aligned}
& \frac{\partial f}{\partial \xi}=\frac{\partial f}{\partial x} \times \frac{\partial x}{\partial \xi}+\frac{\partial f}{\partial y} \times \frac{\partial y}{\partial \xi} \\
& \frac{\partial f}{\partial \eta}=\frac{\partial f}{\partial x} \times \frac{\partial x}{\partial \eta}+\frac{\partial f}{\partial y} \times \frac{\partial y}{\partial \eta}
\end{aligned}
$$

Above equations in matrix form can be expressed as,

$$
\left\{\begin{array}{l}
\frac{\partial f}{\partial \xi} \\
\frac{\partial f}{\partial \eta}
\end{array}\right\}=\left[\begin{array}{ll}
\frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} \\
\frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta}
\end{array}\right\}\left[\begin{array}{l}
\frac{\partial f}{\partial x} \\
\frac{\partial f}{\partial y}
\end{array}\right\}
$$

This is known as the
Jacobian matrix ( J ) for the mapping $(\xi, \eta) \rightarrow(x, y)$

## From equations,

$$
\begin{aligned}
& x=N_{1} x_{1}+N_{2} x_{2}+N_{3} x_{3}+N_{4} x_{4} \\
& y=N_{1} y_{1}+N_{2} y_{2}+N_{3} y_{3}+N_{4} y_{4} \\
& {[J]=\left[\begin{array}{ll}
\frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} \\
\frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta}
\end{array}\right]=\left[\begin{array}{ll}
J_{11} & J_{12} \\
J_{21} & J_{22}
\end{array}\right]}
\end{aligned}\left\{\begin{array}{l}
J_{11}=\frac{\partial x}{\partial \xi}=\frac{\partial N_{1}}{\partial \xi} x_{1}+\frac{\partial N_{2}}{\partial \xi} x_{2}+\frac{\partial N_{3}}{\partial \xi} x_{3}+\frac{\partial N_{4}}{\partial \xi} x_{4} \\
J_{12}=\frac{\partial y}{\partial \xi}=\frac{\partial N_{1}}{\partial \xi} y_{1}+\frac{\partial N_{2}}{\partial \xi} y_{2}+\frac{\partial N_{3}}{\partial \xi} y_{3}+\frac{\partial N_{4}}{\partial \xi} y_{4} \\
J_{21}=\frac{\partial x}{\partial \eta}=\frac{\partial N_{1}}{\partial \eta} x_{1}+\frac{\partial N_{2}}{\partial \eta} x_{2}+\frac{\partial N_{3}}{\partial \eta} x_{3}+\frac{\partial N_{4}}{\partial \eta} x_{4} \\
J_{22}=\frac{\partial y}{\partial \eta}=\frac{\partial N_{1}}{\partial \eta} y_{1}+\frac{\partial N_{2}}{\partial \eta} y_{2}+\frac{\partial N_{3}}{\partial \eta} y_{3}+\frac{\partial N_{4}}{\partial \eta} y_{4}
\end{array}\right.
$$

Shape functions for an isoparametric Quadrilateral element is given by,

$$
\begin{aligned}
& N_{1}=\frac{1}{4}(1-\xi)(1-\eta) \\
& N_{2}=\frac{1}{4}(1+\xi)(1-\eta) \\
& N_{3}=\frac{1}{4}(1+\xi)(1+\eta) \\
& N_{4}=\frac{1}{4}(1-\xi)(1+\eta)
\end{aligned}
$$

$$
\begin{array}{rlrl}
J_{11}=\frac{\partial x}{\partial \xi}=\frac{\partial N_{1}}{\partial \xi} x_{1}+\frac{\partial N_{2}}{\partial \xi} x_{2}+\frac{\partial N_{3}}{\partial \xi} x_{3}+\frac{\partial N_{4}}{\partial \xi} x_{4} & \frac{\partial N_{1}}{\partial \xi} & =\frac{1}{4}(-1)(1-\eta)=\frac{-1}{4}(1-\eta) \\
J_{12}=\frac{\partial y}{\partial \xi}=\frac{\partial N_{1}}{\partial \xi} y_{1}+\frac{\partial N_{2}}{\partial \xi} y_{2}+\frac{\partial N_{3}}{\partial \xi} y_{3}+\frac{\partial N_{4}}{\partial \xi} y_{4} & \frac{\partial N_{2}}{\partial \xi} & =\frac{1}{4}(1)(1-\eta)=\frac{1}{4}(1-\eta) \\
J_{21} & =\frac{\partial x}{\partial \eta}=\frac{\partial N_{1}}{\partial \eta} x_{1}+\frac{\partial N_{2}}{\partial \eta} x_{2}+\frac{\partial N_{3}}{\partial \eta} x_{3}+\frac{\partial N_{4}}{\partial \eta} x_{4} & \frac{\partial N_{3}}{\partial \xi} & =\frac{1}{4}(1)(1+\eta)=\frac{1}{4}(1+\eta) \\
J_{22}=\frac{\partial y}{\partial \eta}=\frac{\partial N_{1}}{\partial \eta} y_{1}+\frac{\partial N_{2}}{\partial \eta} y_{2}+\frac{\partial N_{3}}{\partial \eta} y_{3}+\frac{\partial N_{4}}{\partial \eta} y_{4} & \frac{\partial N_{4}}{\partial \xi} & =\frac{1}{4}(-1)(1+\eta)=\frac{-1}{4}(1+\eta) \\
\frac{\partial N_{1}}{\partial \eta} & =\frac{1}{4}(1-\xi)(-1)=\frac{-1}{4}(1-\xi) \\
\frac{\partial N_{2}}{\partial \eta} & =\frac{1}{4}(1+\xi)(-1)=\frac{-1}{4}(1+\xi) \\
\frac{\partial N_{3}}{\partial \eta} & =\frac{1}{4}(1+\xi)(1)=\frac{1}{4}(1+\xi) \\
\frac{\partial N_{4}}{\partial \eta} & =\frac{1}{4}(1-\xi)(1)=\frac{1}{4}(1-\xi)
\end{array}
$$

$\left\{\begin{array}{l}\frac{\partial f}{\partial \xi} \\ \frac{\partial f}{\partial \eta}\end{array}\right\}=[\mathrm{J}]\left\{\begin{array}{l}\frac{\partial f}{\partial x} \\ \frac{\partial f}{\partial y}\end{array}\right\}$
$\left\{\begin{array}{l}\frac{\partial f}{\partial \xi} \\ \frac{\partial f}{\partial \eta}\end{array}\right\}=\left[\begin{array}{ll}J_{11} & J_{12} \\ J_{21} & J_{22}\end{array}\right]\left\{\begin{array}{l}\frac{\partial f}{\partial x} \\ \frac{\partial f}{\partial y}\end{array}\right\}$
$\left\{\begin{array}{l}\frac{\partial f}{\partial x} \\ \frac{\partial f}{\partial y}\end{array}\right\}=\left[\begin{array}{ll}J_{11} & J_{12} \\ J_{21} & J_{22}\end{array}\right]^{-1}\left\{\begin{array}{l}\frac{\partial f}{\partial \xi} \\ \frac{\partial f}{\partial \eta}\end{array}\right\}$
$\left\{\begin{array}{l}\frac{\partial f}{\partial x} \\ \frac{\partial f}{\partial y}\end{array}\right\}=\frac{1}{|J|}\left[\begin{array}{cc}J_{22} & -J_{12} \\ -J_{21} & J_{11}\end{array}\right]\left[\begin{array}{l}\frac{\partial f}{\partial \xi} \\ \frac{\partial f}{\partial \eta}\end{array}\right\}$

$$
\left\{\begin{array}{l}
\frac{\partial f}{\partial x} \\
\frac{\partial f}{\partial y}
\end{array}\right\}=\frac{1}{|J|}\left[\begin{array}{cc}
J_{22} & -J_{12} \\
-J_{21} & J_{11}
\end{array}\right]\left\{\begin{array}{l}
\frac{\partial f}{\partial \xi} \\
\frac{\partial f}{\partial \eta}
\end{array}\right\}
$$

Replacing $f$ with $u$ and $v$ separately we get,

$$
\begin{aligned}
& \left\{\begin{array}{l}
\frac{\partial u}{\partial x} \\
\frac{\partial u}{\partial y}
\end{array}\right\}=\frac{1}{|J|}\left[\begin{array}{cc}
J_{22} & -J_{12} \\
-J_{21} & J_{11}
\end{array}\right]\left\{\begin{array}{l}
\frac{\partial u}{\partial \xi} \\
\frac{\partial u}{\partial \eta}
\end{array}\right\} \\
& \left\{\begin{array}{l}
\frac{\partial v}{\partial x} \\
\frac{\partial v}{\partial y}
\end{array}\right\}=\frac{1}{|J|}\left[\begin{array}{cc}
J_{22} & -J_{12} \\
-J_{21} & J_{11}
\end{array}\right]\left\{\begin{array}{l}
\frac{\partial v}{\partial \xi} \\
\frac{\partial v}{\partial \eta}
\end{array}\right\}
\end{aligned}
$$

Using the strain-displacement relations

$$
\boldsymbol{\epsilon}=\left\{\begin{array}{l}
\frac{\partial u}{\partial x} \\
\frac{\partial v}{\partial y} \\
\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}
\end{array}\right\}=\frac{1}{|J|}\left[\begin{array}{cccc}
J_{22} & -J_{12} & 0 & 0 \\
0 & 0 & -J_{21} & J_{11} \\
-J_{22} & J_{11} & J_{22} & -J_{12}
\end{array}\right\}\left\{\begin{array}{l}
\frac{\partial u}{\partial \xi} \\
\frac{\partial u}{\partial \eta} \\
\frac{\partial v}{\partial \xi} \\
\frac{\partial v}{\partial \eta}
\end{array}\right\}
$$

$$
\boldsymbol{\epsilon}=\left\{\begin{array}{c}
\boldsymbol{\epsilon}_{x} \\
\boldsymbol{\epsilon}_{y} \\
\gamma_{x y}
\end{array}\right\}=\left\{\begin{array}{c}
\frac{\partial u}{\partial x} \\
\frac{\partial v}{\partial y} \\
\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}
\end{array}\right\}
$$

From equations,

$$
\begin{aligned}
& u=N_{1} u_{1}+N_{2} u_{2}+N_{3} u_{3}+N_{4} u_{4} \\
& v=N_{1} v_{1}+N_{2} v_{2}+N_{3} v_{3}+N_{4} v_{4}
\end{aligned}
$$

Differentiating the above equations w.r. to ' $\xi$ ' and ' $\eta$ ',

$$
\begin{aligned}
& \frac{\partial u}{\partial \xi}=\frac{\partial N_{1}}{\partial \xi} u_{1}+\frac{\partial N_{2}}{\partial \xi} u_{2}+\frac{\partial N_{3}}{\partial \xi} u_{3}+\frac{\partial N_{4}}{\partial \xi} u_{4} \\
& \frac{\partial u}{\partial \eta}=\frac{\partial N_{1}}{\partial \eta} u_{1}+\frac{\partial N_{2}}{\partial \eta} u_{2}+\frac{\partial N_{3}}{\partial \eta} u_{3}+\frac{\partial N_{4}}{\partial \eta} u_{4} \\
& \frac{\partial v}{\partial \xi}=\frac{\partial N_{1}}{\partial \xi} v_{1}+\frac{\partial N_{2}}{\partial \xi} v_{2}+\frac{\partial N_{3}}{\partial \xi} v_{3}+\frac{\partial N_{4}}{\partial \xi} v_{4} \\
& \frac{\partial v}{\partial \eta}=\frac{\partial N_{1}}{\partial \eta} v_{1}+\frac{\partial N_{2}}{\partial \eta} v_{2}+\frac{\partial N_{3}}{\partial \eta} v_{3}+\frac{\partial N_{4}}{\partial \eta} v_{4}
\end{aligned} \quad\left[\begin{array}{l}
\frac{\partial u}{\partial \xi} \\
\frac{\partial u}{\partial \eta} \\
\frac{\partial v}{\partial \xi} \\
\frac{\partial v}{\partial \eta}
\end{array}\right\}=\left[\begin{array}{cccccccc}
\frac{\partial N_{1}}{\partial \xi} & 0 & \frac{\partial N_{2}}{\partial \xi} & 0 & \frac{\partial N_{3}}{\partial \xi} & 0 & \frac{\partial N_{4}}{\partial \xi} & 0 \\
\frac{\partial N_{1}}{\partial \eta} & 0 & \frac{\partial N_{2}}{\partial \eta} & 0 & \frac{\partial N_{3}}{\partial \eta} & 0 & \frac{\partial N_{4}}{\partial \eta} & 0 \\
0 & \frac{\partial N_{1}}{\partial \xi} & 0 & \frac{\partial N_{2}}{\partial \xi} & 0 & \frac{\partial N_{3}}{\partial \xi} & 0 & \frac{\partial N_{4}}{\partial \xi} \\
0 & \frac{\partial N_{1}}{\partial \eta} & 0 & \frac{\partial N_{2}}{\partial \eta} & 0 & \frac{\partial N_{3}}{\partial \eta} & 0 & \frac{\partial N_{4}}{\partial \eta}
\end{array}\right]\left\{\begin{array}{l}
u_{1} \\
v_{1} \\
u_{2} \\
v_{2} \\
u_{3} \\
v_{3} \\
u_{4} \\
v_{4}
\end{array}\right\}
$$

$\boldsymbol{\epsilon}=\left\{\begin{array}{c}\boldsymbol{\epsilon}_{\boldsymbol{x}} \\ \boldsymbol{\epsilon}_{\boldsymbol{y}} \\ \gamma_{x y}\end{array}\right\}=\frac{1}{|J|}\left[\begin{array}{ccccccccc}J_{22} & -J_{12} & 0 & 0 \\ 0 & 0 & -J_{21} & J_{11} \\ -J_{21} & J_{11} & J_{22} & -J_{12}\end{array}\right] \times\left[\begin{array}{ccccccc}\frac{\partial N_{1}}{\partial \xi} & 0 & \frac{\partial N_{2}}{\partial \xi} & 0 & \frac{\partial N_{3}}{\partial \xi} & 0 & \frac{\partial N_{4}}{\partial \xi} \\ 0 \\ \frac{\partial N_{1}}{\partial \eta} & 0 & \frac{\partial N_{2}}{\partial \eta} & 0 & \frac{\partial N_{3}}{\partial \eta} & 0 & \frac{\partial N_{4}}{\partial \eta} \\ 0 & 0 \\ 0 & \frac{\partial N_{1}}{\partial \xi} & 0 & \frac{\partial N_{2}}{\partial \xi} & 0 & \frac{\partial N_{3}}{\partial \xi} & 0 \\ \frac{\partial N_{4}}{\partial \xi} \\ 0 & \frac{\partial N_{1}}{\partial \eta} & 0 & \frac{\partial N_{2}}{\partial \eta} & 0 & \frac{\partial N_{3}}{\partial \eta} & 0 \\ \frac{\partial N_{4}}{\partial \eta}\end{array}\right]\left[\begin{array}{l}u_{1} \\ v_{1} \\ u_{2} \\ v_{2} \\ u_{3} \\ v_{3} \\ u_{4} \\ v_{4}\end{array}\right\}$
[B]
Strain Displacement Matrix

$$
[B]=\frac{1}{|J|}\left[\begin{array}{cccc}
J_{22} & -J_{12} & 0 & 0 \\
0 & 0 & -J_{21} & J_{11} \\
-J_{21} & J_{11} & J_{22} & -J_{12}
\end{array}\right] \times \frac{1}{4}\left[\begin{array}{cccccccc}
-(1-\eta) & 0 & (1-\eta) & 0 & (1+\eta) & 0 & -(1+\eta) & 0 \\
-(1-\xi) & 0 & -(1+\xi) & 0 & (1+\xi) & 0 & (1-\xi) & 0 \\
0 & -(1-\eta) & 0 & (1-\eta) & 0 & (1+\eta) & 0 & -(1+\eta) \\
0 & -(1-\xi) & 0 & -(1+\xi) & 0 & (1+\xi) & 0 & (1-\xi)
\end{array}\right]
$$

## Element stiffness matrix is given by,

$$
\begin{aligned}
{[K] } & =\oint[B]^{T}[D][B] d v \\
{[k] } & =\iint_{A}[B]^{T}[D][B] t d x d y
\end{aligned}
$$

Where,
[B]- Strain-displacement matrix
[D]- Stress-strain relationship matrix
$t$ - Thickness of the element

$$
[k]=\iint_{A}[B]^{T}[D][B] t d x d y
$$

The above equation is in global co-ordinate system.

$$
\int_{A}^{\infty} f(x, y) d x d y=\iint_{A} f\left(\sum_{A}^{5} \eta\right)|d| d \sum_{5} d r
$$

In natural co-ordinate system, stiffness matrix is given by,

$$
[k]=\int_{-1}^{1} \int_{-1}^{1}[B]^{T}[D][B] t|J| d \xi d \eta
$$

Element Stress $\{\sigma\}$
$\{\sigma\}=[D][B]\{\delta\}$
[D] - Stress-strain relationship matrix
Q. Figure shows a four-node quadrilateral. The ( $\mathrm{x}, \mathrm{y}$ ) co-ordinates of each node are given in the figure. The element displacement vector $\{\delta\}$ is given as,

$$
\{\delta\}=[0,0,0.20,0,0.15,0.10,0,0.05]^{\top}
$$

Find the following,
(a) The $x$, $y$-coordinates of a point $P$ whose location in the master element is given by $\xi=0.5$ and $\eta=0.5$ and
(b) The $u$, $v$ displacements of the point $P$.

Given that,
Displacement vector,

$$
\begin{aligned}
& \left.\{\delta\}=\begin{array}{cccccccc}
u_{1} & v_{1} & u_{2} & v_{2} & u_{3} & v_{3} & u_{4} & v_{4} \\
0, & 0, & 0.20, & 0, & 0.15 & 0.10, & 0, & 0.05
\end{array}\right]^{\mathrm{T}} \\
& \eta=\xi=0.5 \Rightarrow \mathbf{P}(\xi, \eta)
\end{aligned}
$$

The coordinates of the four noded quadratic element are,

$$
\begin{array}{ll}
x_{1}=1 ; & y_{1}=1 \\
x_{2}=5 ; & y_{2}=1 \\
x_{3}=6 ; & y_{3}=6 \\
x_{4}=1 ; & y_{4}=4
\end{array}
$$



Figure

## (a) Coordinates of Point ' P '

The coordinates of point $P(x, y)$ are given by,

$$
\begin{aligned}
& x=N_{1} x_{1}+N_{2} x_{2}+N_{3} x_{3}+N_{4} x_{4} \\
& y=N_{1} y_{1}+N_{2} y_{2}+N_{3} y_{3}+N_{4} y_{4}
\end{aligned}
$$

$$
N_{4}=\frac{1}{4}(1-\xi)(1+\eta) \quad{ }^{+} \quad N_{3}=\frac{1}{4}(1+\xi)(1+\eta)
$$



$$
N_{1}=\frac{1}{4}(1-\xi)(1-\eta) \quad N_{2}=\frac{1}{4}(1+\xi)(1-\eta)
$$

Shape functions of a four noded quadratic element are,

$$
\begin{aligned}
N_{1}=\frac{(1-\xi)(1-\eta)}{4} & =\frac{(1-0.5)(1-0.5)}{4}=0.0625 \\
N_{2}=\frac{(1+\xi)(1-\eta)}{4} & =\frac{(1+0.5)(1-0.5)}{4}=0.1875 \\
N_{3}=\frac{(1+\xi)(1+\eta)}{4} & =\frac{(1+0.5)(1+0.5)}{4}=0.5625 \\
N_{4}=\frac{(1-\xi)(1+\eta)}{4} & =\frac{(1-0.5)(1+0.5)}{4}=0.1875 \\
\therefore \quad x & =N_{1} x_{1}+N_{2} x_{2}+N_{3} x_{3}+N_{4} x_{4} \\
& =0.0625(1)+0.1875(5)+0.5625(6)+0.1875(1) \\
x & =4.5625 \\
y & =N_{1} y_{1}+N_{2} y_{2}+N_{3} y_{3}+N_{4} y_{4} \\
& =0.0625(1)+0.1875(1)+0.5625(6)+0.1875 \\
& =4.375
\end{aligned}
$$

$$
\therefore P(x, y)=(4.5625,4.375)
$$

## (b) Displacements of Point ' P '

$u, v$ displacements of point ' $P$ ' are given by,

$$
\begin{aligned}
& u= N_{1} u_{1}+N_{2} u_{2}+N_{3} u_{3}+N_{4} u_{4} \\
&= 0.0625(0)+0.1875(0.20)+0.5625(0.15) \\
& \quad \quad+0.1875(0) \\
&= 0.121875 \\
& v= N_{1} v_{1}+N_{2} v_{2}+N_{3} v_{3}+N_{4} v_{4} \\
&= 0.0625(0)+0.1875(0)+0.5625(0.10)+0.1875 \\
&(0.05) \\
& v= 0.065625
\end{aligned}
$$

$$
\therefore(u, v)=(0.121875,0.065625)
$$

Q. Consider a rectangular element as shown in figure. Assume plane stress condition, $\mathbf{E}=206850 \mathrm{MPa}$, $v=0.3,\{\delta\}=[0,0,0.05,0.075,0.15,0.8,0,0]^{\top} \mathrm{cm}$. Evaluate Jacobian J, B and $\sigma$ at $\xi=0$ and $\eta=0$.


## Given that,

Young's modulus, $E=206850 \mathrm{MPa}$

$$
\begin{aligned}
& =206850 \mathrm{~N} / \mathrm{mm}^{2} \\
E & =20.685 \times 10^{6} \mathrm{~N} / \mathrm{cm}^{2}
\end{aligned}
$$

Poisons ratio, $v=0.3$
Displacement vector,

$$
\{\delta\}=\begin{array}{cccccccc}
u_{1} & v_{1} & u_{2} & v_{2} & u_{3} & v_{3} & u_{4} & v_{4} \\
{[0,} & 0, & 0.05, & 0.075, & 0.15 & 0.8, & 0, & 0]^{\mathrm{T}} \mathrm{~cm}
\end{array}
$$



Local coordinates, $\xi=0, \eta=0$
Co-ordinates of the Quadratic element is

For a plane stress condition,

$$
[D]=\frac{E}{1-v^{2}}\left[\begin{array}{ccc}
1 & v & 0 \\
v & 1 & 0 \\
0 & 0 & \frac{1-v}{2}
\end{array}\right]
$$

$$
\begin{aligned}
& x_{1}=0, y_{1}=0 \\
& x_{2}=5, y_{2}=0 \\
& x_{3}=5, y_{3}=2.5 \\
& x_{4}=0, y_{4}=2.5
\end{aligned}
$$

Jacobian J, B and $\sigma$

$$
\begin{aligned}
& \text { [J]- Jacobian matrix } \\
& {[J]=\left[\begin{array}{ll}
\frac{\partial x}{\partial \xi} & \frac{\partial y}{\partial \xi} \\
\frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta}
\end{array}\right]=\left[\begin{array}{ll}
J_{11} & J_{12} \\
J_{21} & J_{22}
\end{array}\right]} \\
& \text { Where, } \\
& \begin{array}{l}
x=N_{1} x_{1}+N_{2} x_{2}+N_{3} x_{3}+N_{4} x_{4} \\
y=N_{1} y_{1}+N_{2} y_{2}+N_{3} y_{3}+N_{4} y_{4}
\end{array} \\
& \frac{\partial N_{1}}{\partial \xi}=\frac{1}{4}(-1)(1-\eta)=\frac{-1}{4}(1-\eta) \\
& \frac{\partial N_{2}}{\partial \xi}=\frac{1}{4}(1)(1-\eta)=\frac{1}{4}(1-\eta) \\
& \frac{\partial N_{3}}{\partial \xi}=\frac{1}{4}(1)(1+\eta)=\frac{1}{4}(1+\eta) \\
& \frac{\partial N_{4}}{\partial \xi}=\frac{1}{4}(-1)(1+\eta)=\frac{-1}{4}(1+\eta) \\
& \frac{\partial N_{1}}{\partial \eta}=\frac{1}{4}(1-\xi)(-1)=\frac{-1}{4}(1-\xi) \\
& \frac{\partial N_{2}}{\partial \eta}=\frac{1}{4}(1+\xi)(-1)=\frac{-1}{4}(1+\xi) \\
& \frac{\partial N_{3}}{\partial \eta}=\frac{1}{4}(1+\xi)(1)=\frac{1}{4}(1+\xi) \\
& N_{4}=\frac{1}{4}(1-\xi)(1+\eta) \\
& N_{1}=\frac{1}{4}(1-\xi)(1-\eta) \\
& N_{2}=\frac{1}{4}(1+\xi)(1-\eta) \\
& N_{3}=\frac{1}{4}(1+\xi)(1+\eta) \\
& \frac{\partial N_{4}}{\partial \eta}=\frac{1}{4}(1-\xi)(1)=\frac{1}{4}(1-\xi)
\end{aligned}
$$

$$
\begin{aligned}
J_{11} & =\frac{1}{4}\left[-(1-\eta) x_{1}+(1-\eta) x_{2}+(1+\eta) x_{3}-(1+\eta) x_{4}\right] & J_{21} & =\frac{1}{4}\left[-(1-\xi) x_{1}-(1+\xi) x_{2}+(1+\xi) x_{3}+(1-\xi) x_{4}\right] \\
& =\frac{1}{4}[(1-0) 0+(1-0) 5+(1+0) 5-(1+0) 0] & & \frac{1}{4}[-(1-0) 0-(1+0) 5+(1+0) 5+(1-0) 0] \\
& =\frac{1}{4}(5+5) & J_{21} & =0 \\
J_{11} & =2.5 & J_{22} & =\frac{1}{4}\left[-(1-\xi) y_{1}-(1+\xi) y_{2}+(1+\xi) y_{3}+(1-\xi) y_{4}\right] \\
J_{12} & =\frac{1}{4}\left[-(1-\eta) y_{1}+(1-\eta) y_{2}+(1+\eta) y_{3}-(1+\eta) y_{4}\right] & & =\frac{1}{4}[-(1-0) 0-(1+0) 0+(1+0) 2.5+(1-0) 2.5] \\
& =\frac{1}{4}[-(1-0) 0+(1-0) 0+(1+0) 2.5-(1+0) 2.5] & & =\frac{1}{4}(5) \\
J_{12} & =0 & J_{22} & =1.25
\end{aligned}
$$

Jacobian Matrix,

$$
[J]=\left[\begin{array}{ll}
J_{11} & J_{12} \\
J_{21} & J_{22}
\end{array}\right]=\left[\begin{array}{cc}
2.5 & 0 \\
0 & 1.25
\end{array}\right]
$$

$$
\begin{aligned}
& \left\{\begin{array}{l}
\frac{\partial f}{\partial x} \\
\frac{\partial f}{\partial y}
\end{array}\right\}=\left[\begin{array}{ll}
J_{11} & J_{12} \\
J_{21} & J_{22}
\end{array}\right]^{-1}\left\{\begin{array}{l}
\frac{\partial f}{\partial \xi} \\
\frac{\partial f}{\partial \eta}
\end{array}\right\} \\
& \left\{\begin{array}{l}
\frac{\partial f}{\partial x} \\
\frac{\partial f}{\partial y}
\end{array}\right\}=\frac{1}{|J|}\left[\begin{array}{cc}
J_{22} & -J_{12} \mid \\
-J_{21} & J_{11}
\end{array}\right\}\left\{\begin{array}{l}
\frac{\partial f}{\partial \xi} \\
\frac{\partial f}{\partial \eta}
\end{array}\right\} \\
& \text { Replacing } f \text { with } u \text { and } v \text { separately we get, } \\
& \left\{\begin{array}{l}
\frac{\partial u}{\partial x} \\
\frac{\partial u}{\partial y}
\end{array}\right\}=\frac{1}{|J|}\left[\begin{array}{cc}
J_{22} & -J_{12} \\
-J_{21} & J_{11}
\end{array}\right]\left[\begin{array}{c}
\frac{\partial u}{\partial \xi} \\
\frac{\partial u}{\partial \eta}
\end{array}\right\} \\
& \left\{\begin{array}{l}
\frac{\partial v}{\partial x} \\
\frac{\partial v}{\partial y}
\end{array}\right\}=\frac{1}{|J|}\left[\begin{array}{cc}
J_{22} & -J_{12} \\
-J_{21} & J_{11}
\end{array}\right]\left\{\begin{array}{l}
\frac{\partial v}{\partial \dot{\xi}} \\
\frac{\partial v}{\partial \eta}
\end{array}\right\} \\
& \boldsymbol{\epsilon}=\left\{\begin{array}{c}
\boldsymbol{\epsilon}_{x} \\
\boldsymbol{\epsilon}_{y} \\
\gamma_{x y}
\end{array}\right\}=\left\{\begin{array}{c}
\frac{\partial u}{\partial x} \\
\frac{\partial v}{\partial y} \\
\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}
\end{array}\right\} \\
& \text { determinant } \\
& =\frac{1}{|J|}\left[\begin{array}{cccc}
J_{22} & -J_{12} & 0 & 0 \\
0 & 0 & -J_{21} & J_{11} \\
-J_{22} & J_{11} & J_{22} & -J_{12}
\end{array}\right]\left\{\begin{array}{l}
\frac{\partial u}{\partial \eta} \\
\frac{\partial v}{\partial \xi} \\
\frac{\partial v}{\partial \eta}
\end{array}\right\}
\end{aligned}
$$

Using the strain-displacement relations

$$
\boldsymbol{\epsilon}=\left\{\begin{array}{l}
\frac{\partial u}{\partial x} \\
\frac{\partial v}{\partial y} \\
\frac{\partial u}{\partial y}+\frac{\partial v}{\partial x}
\end{array}\right\}=\frac{1}{|J|}\left[\begin{array}{cccc}
J_{22} & -J_{12} & 0 & 0 \\
0 & 0 & -J_{21} & J_{11} \\
-J_{22} & J_{11} & J_{22} & -J_{12}
\end{array}\right]\left\{\begin{array}{l}
\frac{\partial u}{\partial \xi} \\
\frac{\partial u}{\partial \eta} \\
\frac{\partial v}{\partial \xi} \\
\frac{\partial v}{\partial \eta}
\end{array}\right\}, \quad \begin{aligned}
u & =N_{1} u_{1}+N_{2} u_{2}+N_{3} u_{3}+N_{4} u_{4} \\
v & =N_{1} v_{1}+N_{2} v_{2}+N_{3} v_{3}+N_{4} v_{4} \\
\frac{\partial u}{\partial \xi} & =\frac{\partial N_{1}}{\partial \xi} u_{1}+\frac{\partial N_{2}}{\partial \xi} u_{2}+\frac{\partial N_{3}}{\partial \xi} u_{3}+\frac{\partial N_{4}}{\partial \xi} u_{4}
\end{aligned}
$$

$$
\left\{\begin{array}{l}
\frac{\partial u}{\partial \xi} \\
\frac{\partial u}{\partial \eta} \\
\frac{\partial v}{\partial \xi} \\
\frac{\partial v}{\partial \eta}
\end{array}\right\}=\left[\begin{array}{cccccccc}
\frac{\partial N_{1}}{\partial \xi} & 0 & \frac{\partial N_{2}}{\partial \xi} & 0 & \frac{\partial N_{3}}{\partial \xi} & 0 & \frac{\partial N_{4}}{\partial \xi} & 0 \\
\frac{\partial N_{1}}{\partial \eta} & 0 & \frac{\partial N_{2}}{\partial \eta} & 0 & \frac{\partial N_{3}}{\partial \eta} & 0 & \frac{\partial N_{4}}{\partial \eta} & 0 \\
0 & \frac{\partial N_{1}}{\partial \xi} & 0 & \frac{\partial N_{2}}{\partial \xi} & 0 & \frac{\partial N_{3}}{\partial \xi} & 0 & \frac{\partial N_{4}}{\partial \xi} \\
0 & \frac{\partial N_{1}}{\partial \eta} & 0 & \frac{\partial N_{2}}{\partial \eta} & 0 & \frac{\partial N_{3}}{\partial \eta} & 0 & \frac{\partial N_{4}}{\partial \eta}
\end{array}\right]\left[\begin{array}{l}
u_{1} \\
v_{1} \\
u_{2} \\
v_{2} \\
u_{3} \\
v_{3} \\
u_{4} \\
v_{4}
\end{array}\right\} \quad \begin{aligned}
\frac{\partial N_{1}}{\partial \xi} & =\frac{1}{4}(-1)(1-\eta)=\frac{-1}{4}(1-\eta) \\
\frac{\partial N_{2}}{\partial \xi} & =\frac{1}{4}(1)(1-\eta)=\frac{1}{4}(1-\eta) \\
\frac{\partial N_{3}}{\partial \xi} & =\frac{1}{4}(1)(1+\eta)=\frac{1}{4}(1+\eta) \\
\frac{\partial N_{4}}{\partial \xi} & =\frac{1}{4}(-1)(1+\eta)=\frac{-1}{4}(1+\eta) \\
\frac{\partial N_{1}}{\partial \eta} & =\frac{1}{4}(1-\xi)(-1)=\frac{-1}{4}(1-\xi) \\
\frac{\partial N_{2}}{\partial \eta} & =\frac{1}{4}(1+\xi)(-1)=\frac{-1}{4}(1+\xi) \\
\frac{\partial N_{3}}{\partial \eta} & =\frac{1}{4}(1+\xi)(1)=\frac{1}{4}(1+\xi) \\
\frac{\partial N_{4}}{\partial \eta} & =\frac{1}{4}(1-\xi)(1)=\frac{1}{4}(1-\xi)
\end{aligned}
$$

## Strain-Displacement Matrix [B]

$$
\begin{aligned}
& \quad[B]=\frac{1}{|J|}\left[\begin{array}{cccc}
J_{22} & -J_{12} & 0 & 0 \\
0 & 0 & -J_{21} & J_{11} \\
-J_{21} & J_{11} & J_{22} & -J_{12}
\end{array}\right] \times \frac{1}{4}\left[\begin{array}{cccccccc}
-(1-\eta) & 0 & (1-\eta) & 0 & (1+\eta) & 0 & -(1+\eta) & 0 \\
-(1-\xi) & 0 & -(1+\xi) & 0 & (1+\xi) & 0 & (1-\xi) & 0 \\
0 & -(1-\eta) & 0 & (1-\eta) & 0 & (1+\eta) & 0 & -(1+\eta) \\
0 & -(1-\xi) & 0 & -(1+\xi) & 0 & (1+\xi) & 0 & (1-\xi)
\end{array}\right] \\
& |J|=\left[\left(J_{11} \times J_{22}\right)-\left(J_{12} \times J_{21}\right)\right] \\
& \quad=(2.5 \times 1.25)-0 \\
& |J|=3.125
\end{aligned}
$$

$$
[B]=\frac{1}{3.125}\left[\begin{array}{cccc}
1.25 & 0 & 0 & 0 \\
0 & 0 & 0 & 2.5 \\
0 & 2.5 & 1.25 & 0
\end{array}\right] \times \frac{1}{4}\left[\begin{array}{cccccccc}
-(1-0) & 0 & (1-0) & 0 & (1+0) & 0 & -(1+0) & 0 \\
-(1-0) & 0 & -(1+0) & 0 & (1+0) & 0 & (1-0) & 0 \\
0 & -(1-0) & 0 & (1-0) & 0 & (1+0) & 0 & -(1+0) \\
0 & -(1-0) & 0 & -(1+0) & 0 & (1+0) & 0 & (1-0)
\end{array}\right]
$$

$$
=\frac{1}{3.125}\left[\begin{array}{cccc}
1.25 & 0 & 0 & 0 \\
0 & 0 & 0 & 2.5 \\
0 & 2.5 & 1.25 & 0
\end{array}\right] \times \frac{1}{4}\left[\begin{array}{cccccccc}
-1 & 0 & 1 & 0 & 1 & 0 & -1 & 0 \\
-1 & 0 & -1 & 0 & 1 & 0 & 1 & 0 \\
0 & -1 & 0 & 1 & 0 & 1 & 0 & -1 \\
0 & -1 & 0 & -1 & 0 & 1 & 0 & 1
\end{array}\right]
$$

$$
\left.\begin{array}{l}
=\left[\begin{array}{cccc}
0.4 & 0 & 0 & 0 \\
0 & 0 & 0 & 0.8 \\
0 & 0.8 & 0.4 & 0
\end{array}\right] \times\left[\begin{array}{cccccccc}
-0.25 & 0 & 0.25 & 0 & 0.25 & 0 & -0.25 & 0 \\
-0.25 & 0 & -0.25 & 0 & 0.25 & 0 & 0.25 & 0 \\
0 & -0.25 & 0 & 0.25 & 0 & 0.25 & 0 & -0.25 \\
0 & -0.25 & 0 & -0.25 & 0 & 0.25 & 0 & 0.25
\end{array}\right] \\
=\left[\begin{array}{ccccccc}
(0.4 \times-0.25) & 0 & (0.4 \times 0.25) & 0 & (0.4 \times 2.5) & 0 & (0.4 \times-0.25) \\
0 & (0.8 \times-0.25) & 0 & (0.8 \times-0.25) & 0 & (0.8 \times 0.25) & 0 \\
(0.8 \times-0.25) & (0.4 \times-0.25) & (0.8 \times 0.25) & (0.4 \times 0.25) & (0.8 \times 0.25) & (0.4 \times 0.25) & (0.8 \times 0.25)
\end{array}(0.4 \times-0.25)\right.
\end{array}\right] .
$$

$$
[B]=\left[\begin{array}{cccccccc}
-0.1 & 0 & 0.1 & 0 & 0.1 & 0 & -0.1 & 0 \\
0 & -0.2 & 0 & -0.2 & 0 & 0.2 & 0 & 0.2 \\
-0.2 & -0.1 & -0.2 & 0.1 & 0.2 & 0.1 & 0.2 & -0.1
\end{array}\right]
$$

## Element Stress $\{\sigma\}$

$\{\sigma\}=[D][B]\{\delta\}$
[D] - Stress-strain relationship matrix
For a plane stress condition,

$$
[D]=\frac{E}{1-v^{2}}\left[\begin{array}{llc}
1 & v & 0 \\
v & 1 & 0 \\
0 & 0 & \frac{1-v}{2}
\end{array}\right]=\frac{20.685 \times 10^{6}}{1-(0.3)^{2}}\left[\begin{array}{ccc}
1 & 0.3 & 0 \\
0.3 & 1 & 0 \\
0 & 0 & \frac{1-0.3}{2}
\end{array}\right]
$$

$[D]=22.73 \times 10^{6}\left[\begin{array}{ccc}1 & 0.3 & 0 \\ 0.3 & 1 & 0 \\ 0 & 0 & 0.35\end{array}\right]$

$$
\{\sigma\}=22.731 \times 10^{6}\left[\begin{array}{ccc}
1 & 0.3 & 0 \\
0.3 & 1 & 0 \\
0 & 0 & 0.35
\end{array}\right]_{3 \times 3}\left[\begin{array}{cccccccc}
-0.1 & 0 & 0.1 & 0 & 0.1 & 0 & -0.1 & 0 \\
0 & -0.2 & 0 & -0.2 & 0 & 0.2 & 0 & 0.2 \\
-0.2 & -0.1 & -0.2 & 0.1 & 0.2 & 0.1 & 0.2 & -0.1
\end{array}\right]_{3 \times 8}\left[\begin{array}{c}
0 \\
0 \\
0.05 \\
0.075 \\
0.15 \\
0.8 \\
0 \\
0
\end{array}\right\}_{8 \times 1}
$$

$$
\begin{aligned}
& =22.731 \times 10^{6}\left[\begin{array}{cccccccc}
-0.1 & -0.06 & 0.1 & -0.06 & 0.1 & 0.06 & -0.1 & 0.06 \\
-0.03 & -0.2 & 0.03 & -0.2 & 0.03 & 0.2 & -0.03 & 0.2 \\
-0.07 & -0.035 & -0.07 & 0.035 & 0.07 & 0.035 & 0.07 & -0.035
\end{array}\right]_{3 \times 8}\left[\begin{array}{c}
0 \\
0 \\
0.05 \\
0.075 \\
0.015 \\
0.8 \\
0 \\
0
\end{array}\right]_{8 \times 1} \\
& =22.731 \times 10^{6}\left[\begin{array}{c}
0.05 \\
0.14695 \\
0.028175
\end{array}\right] \\
\{\sigma\} & =\left\{\begin{array}{c}
1.136 \times 10^{6} \\
3.34 \times 10^{6} \\
0.64 \times 10^{6}
\end{array}\right\} \mathrm{N} / \mathrm{cm}^{2}
\end{aligned}
$$

Consider a quadrilateral element as shown in figure. The local coordinates are $\xi=0.5$, $\eta=0.5$, Evaluate Jacobian matrix and strain Displacement matrix.


Figure

## Numerical Integration

Element stiffness matrix is given by,

$$
\begin{gathered}
{[K]=\oint[B]^{T}[D][B] d v} \\
{[k]=\iint_{A}[B]^{T}[D][B] t d x d y}
\end{gathered}
$$

$$
\iint_{A} f(x, y) d x d y=\iint_{A} f(\xi, \eta)|J| d \xi d \eta
$$

$$
[k]=\int_{-1}^{1} \int_{-1}^{1}[B]^{T}[D][B] t|J| d \xi d \eta
$$

Consider Gauss Quadrature formula, $I=\int_{a}^{b} f(x) d x=\sum_{i=1}^{n} W_{i} f\left(x_{i}\right)$

Where,
$W_{i}-$ Weight function
$x_{i}-$ Sampling point
one point technique

Let, $n=1$.

$$
I=\int_{-1}^{1} f(x) d x=W_{1} f\left(x_{1}\right)
$$

The polynomial order is $(2 n-1)$

$$
\begin{aligned}
& =(2(1)-1) \\
& =1
\end{aligned}
$$

Therefore, the polynomial function is given by,

$$
f(x)=a_{1}+a_{2} x
$$

Gauss quadrature of linear polynomial

$$
\begin{array}{rlr}
I= & \int_{-1}^{1}\left(a_{1}+a_{2} x\right) d x=W_{1} f\left(x_{1}\right) \\
& \int_{-1}^{1}\left(a_{1}+a_{2} x\right) d x-W_{1} f\left(x_{1}\right)=0 \\
& \int_{-1}^{1}\left(a_{1}+a_{2} x\right) d x-W_{1}\left(a_{1}+a_{2} x_{1}\right)=0 \\
& a_{1}[x]_{-1}^{1}+a_{2}\left[\frac{x^{2}}{2}\right]_{-1}^{1}-W_{1}\left(a_{1}+a_{2} x_{1}\right)=0 \\
& 2 a_{1}-W_{1} a_{1}-W_{1} a_{2} x_{1}=0 \\
& a_{1}\left(2-W_{1}\right)-W_{1} a_{2} x_{1}=0
\end{array}
$$

## Two Point Technique (Formula)

Consider Gauss Quadrature formula,

$$
I=\int_{a}^{b} f(x) d x=\sum_{i=1}^{n} W_{i} f\left(x_{i}\right)
$$

Let, $n=2$.

$$
I=\int_{-1}^{1} f(x) d x=W_{1} f\left(x_{1}\right)+W_{2} f\left(x_{2}\right)
$$

The polynomial order is $(2 n-1)$

$$
\begin{aligned}
& =(2(2)-1) \\
& =3
\end{aligned}
$$

## Gauss quadrature of Cubic polynomial

Therefore, the polynomial function is given by,

$$
f(x)=a_{1}+a_{2} x+a_{3} x^{2}+a_{4} x^{3}
$$

$$
\begin{aligned}
& {\left[a_{1} x+a_{2} \frac{x^{2}}{2}+a_{3} \frac{x^{3}}{3}+a_{4} \frac{x^{4}}{4}\right]_{-1}^{1}-W_{1} f\left(x_{1}\right)-W_{2} f\left(x_{2}\right)=0} \\
& {\left[2 a_{1}+0+\frac{2}{3} a_{3}+0\right]-W_{1}\left(a_{1}+a_{2} x_{1}+a_{3} x_{1}^{2}+a_{3} x_{1}^{3}\right)-W_{2}\left(a_{1}+a_{2} x_{2}+a_{3} x_{2}^{2}+a_{3} x_{2}^{3}\right)=0} \\
& a_{1}\left[2-\left(W_{1}+W_{2}\right)\right]-a_{2}\left[W_{1} x_{1}+W_{2} x_{2}\right]-a_{3}\left[W_{1} x_{1}^{2}+W_{2} x_{2}^{2}-\frac{2}{3}\right]-a_{4}\left[W_{1} x_{1}^{3}+W_{2} x_{2}^{3}\right]=0
\end{aligned}
$$

satisfies only if,

$$
\begin{array}{lr}
W_{1}+W_{2}=2 \\
W_{1} x_{2}+W_{2} x_{2}=0 & \text { On solving the above equations, } \\
W_{1} x_{1}^{2}+W_{2} x_{2}^{2}=\frac{2}{3} & W_{1}=W_{2}=1 \\
W_{1} x_{1}^{3}+W_{2} x_{2}^{3}=0 & x_{1}=\frac{1}{\sqrt{3}}=0.5773502 \\
x_{2}=\frac{-1}{\sqrt{3}}=-0.5773502
\end{array}
$$

On solving the above equations,

$$
\begin{aligned}
& W_{1}=W_{2}=1 \\
& x_{1}=\frac{1}{\sqrt{3}}=0.5773502 \\
& x_{2}=\frac{-1}{\sqrt{3}}=-0.5773502 \\
& I=\int_{-1}^{1} f(x) d x=(1) f(0.5773502)+(1) f(-0.5773502)
\end{aligned}
$$

## Two Dimensional Problem

Integral to evaluate two dimensional problem,

$$
I=\int_{-1}^{1} \int_{-1}^{1} f(x, y) d x d y
$$

On generalizing the relation,

$$
\begin{aligned}
I & \simeq \int_{-1}^{1}\left[\sum_{i=1}^{n} W_{i} f\left(x_{i}, y\right)\right] d y \\
& \simeq \sum_{j=1}^{n} W_{j}\left[\sum_{i=1}^{n} W_{i} f\left(x_{i}, y_{j}\right)\right] \\
I & \simeq \sum_{i=1}^{n} \sum_{j=1}^{n} W_{i} W_{j} f\left(x_{i}, y_{j}\right)
\end{aligned}
$$

Using two Gauss integration points,
For $n=2$,

$$
\begin{aligned}
& x_{i}=y_{j}= \pm 0.57735 \\
& W_{i}=1
\end{aligned}
$$

i.e.,

$$
\begin{aligned}
& x_{1}=y_{1}=0.5773 j \\
& x_{2}=y_{2}=-0.57735 \\
& W_{1}=W_{2}=1
\end{aligned}
$$

Then, equation (2) becomes,

$$
\begin{aligned}
I & \simeq W_{1}^{2} f\left(x_{1}, y_{1}\right)+W_{2}^{2} f\left(x_{2}, y_{2}\right) \\
\therefore \quad I & \simeq f(0.57735,0.57735)+f(-0.57735,-0.57735)
\end{aligned}
$$

## Unit 4

## Dynamic Analysis:

Formulation of finite element model, element consistent and lumped mass matrices, Evaluation of eigenvalues and eigenvectors, Free vibration analysis.
Steady state heat transfer analysis: one dimensional analysis of a fin.
Introduction to FE software.

## Types of analysis of a Problem

## 1. Static Analysis

In case, if the interia effects are not considered and damping is zero, then it is said to be quasistatic and the analysis is a static analysis.

$$
[K]\{\delta\}=\{F\}
$$

Where,
[ $K$ ] - Global stiffness matrix
$\{\delta\}$ - Global displacement vector
$\{F\}$ - Global load vector

## Types of analysis of a Problem

## 2. Eigen Value Analysis

If the interia effects are taken into consideration with zero damping and applied loads, then the equations of motion is reduced to a generalised eigne value problem. In such cases, following can be used to obtain the solution,

$$
[M]\{\ddot{\delta}\}+[K]\{\delta\}=0
$$

Where,
[ $M$ ] - Global mass matrix
$\{\ddot{\delta}\}$ - Global nodal acceleration vector

## Types of analysis of a Problem

3. Transient Dynamic Analysis

If the loads are arbitrary but known functions of time, then the analysis is known as transient dynamic analysis. In such cases,

$$
[M]\{\ddot{\delta}\}+[C]\{\dot{\delta}\}+[K]\{\delta\}=F(t)
$$

Where,
[C] - Global viscous damping matrix
$\{\dot{\delta}\}$ - Global nodal velocity vector
4. Frequency Response Analysis

If the structure is subjected to harmonic loads, then the analysis is known as frequency response analysis. In such cases,

$$
[M]\{\ddot{\delta}\}+[C]\{\dot{\delta}\}+[K]\{\delta\}=F \sin (\omega t)
$$

## Steps in an FEM Analysis

## Preprocessor/Modeling

## Analysis run/Solve

Post-processing/View results

## 

- Formulation of element stiffness matrices
- Assembly of global stiffness matrix [K]
- formulation of load vector $\{F\}$
- Solution of $[K]\{u\}=\{F\}$ to get nodal displacements $\{u\}$


## in case of dynamic analysis)

$\left.[\mathrm{K}] \Rightarrow[\mathrm{K}]-\mathrm{N}_{[\mathrm{M}}\right]$

| Longitudinal Vibrations |  | Transverse Vibrations |  |
| :--- | :--- | :--- | :--- |
| 1. | The vibrations, in which the particles of the shaft <br> move parallel to the axis of the shaft are called <br> longitudinal vibrations. | 1. | The vibrations, in which particles of the shaft move <br> perpendicular to the axis of the shaft are called <br> transverse vibrations. |
| 2. | Figure: Transverse Vibration |  |  |


| Longitudinal Vibrations |  | Transverse Vibrations |  |
| :---: | :--- | :--- | :--- |
| 3. | In case of longitudinal vibrations, the shaft is <br> elongated and shortened alternately. | 3. | In case of transverse vibrations, the shaft is straight <br> and bent alternately. |
| 4. | Tensile and compressive stresses are induced alternately. <br> The natural frequency of free longitudinal vibrations is, | 4. | 5. |
| Bending stresses are induced. |  |  |  |
| The natural frequency of transverse vibrations is, |  |  |  |
| $f_{n}=\frac{0.4985}{\sqrt{\delta}}$ |  |  |  |$\quad$| $f_{n}=\frac{0.4985}{\sqrt{\delta}}$ |
| :--- |

$$
f_{n 1}=\frac{1}{2 \pi} \sqrt{\frac{g}{\Delta_{1}}}=\frac{\sqrt{9.81}}{2 \pi} \frac{1}{\sqrt{\Delta_{1}}}=\frac{0.4985}{\sqrt{\Delta_{1}}}
$$

longitudinal vibration of bar element,

## Finite Element Equation

$$
\left[[K]-\omega^{2}[M]\right]\{\delta\}=\{F\}
$$

$[K]$ - Stiffness matrix
[ $M$ ] - Mass matrix
$\{\delta\}$ - Displacement vector
$\{F\}$ - Force element
Stiffness Matrix
$[K]=\frac{E A}{L}\left[\begin{array}{cc}1 & -1 \\ -1 & 1\end{array}\right]$
$E$ - Young's modulus of the bar material
$A$ - Area of crossection of the bar
$L$ - Length of the bar
transverse vibration of beam element,

## Finite Element Equation

$$
\left[[K]-\omega^{2}[M]\right]\{\delta\}=\{F\}
$$

Stiffness Matrix

$$
[K]=\frac{E I}{l^{3}}\left[\begin{array}{cccc}
12 & 6 l & -12 & 6 l \\
6 l & 4 l^{2} & -6 l & 2 l^{2} \\
-12 & -6 l & 12 & -6 l \\
6 l & 2 l^{2} & -6 l & 4 l^{2}
\end{array}\right]
$$

longitudinal vibration of bar element,
Consistent Mass Matrix

$$
[M]=\frac{\rho A L}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right]
$$

$\rho$ - Density of the bar element
$A$ - Area of cross-section of bar
$L$ - Length of the bar

Lumped Mass Matrix

$$
[M]=\frac{\rho A L}{2}\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]
$$

transverse vibration of beam element,

## Consistent Mass Matrix

$$
\left.[M]=\frac{\rho A L}{420}\left[\begin{array}{cccc}
156 & 22 l & 54 & -13 l \\
22 l & 4 l^{2} & 13 l & -3 l^{2} \\
54 & 13 l & 156 & -22 l \\
-13 l & -3 l^{2} & -22 l & 4 l^{2}
\end{array}\right] \right\rvert\,
$$

## Lumped Mass Matrix

$$
[M]=\frac{\rho A l}{2}\left[\begin{array}{llll}
1 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 0
\end{array}\right]
$$

Find the natural frequencies for the longitudinal vibrations of the stepped bar. Assume $A_{1}=2 A$, $A_{2}=A$ and $E_{1}=E_{2}=E$.


$$
\left[[K]-\omega^{2}[M]\{\delta\}=\{F\}\right.
$$



Figure (2): Bar Element Model

## Stiffness Matrix

For element (1),

$$
\begin{aligned}
\therefore \quad\left[K_{1}\right] & =\frac{E_{1} A_{1}}{L_{1}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
& =\frac{E(2 A)}{L / 2}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
{\left[K_{1}\right] } & =\frac{2 E A}{L}\left[\begin{array}{cc}
2 & -2 \\
-2 & 2
\end{array}\right]
\end{aligned}
$$

For element (2),

$$
\begin{aligned}
{\left[K_{2}\right] } & =\frac{E_{2} A_{2}}{L_{2}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
& =\frac{E(A)}{L / 2}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
& =\frac{2 E A}{L}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]
\end{aligned}
$$

Global Stiffness Matrix

$$
\left.\begin{array}{rl}
\therefore & {[K]} \\
& =\left[K_{1}\right]+\left[K_{2}\right] \\
& {[K]}
\end{array}\right]\left[\begin{array}{ccc}
2 & -2 & 0 \\
-2 & 3 & -1 \\
0 & -1 & 1
\end{array}\right] \times \frac{2 E A}{L},
$$

## Mass Matrix

For element (1),

$$
\begin{aligned}
{\left[M_{1}\right] } & =\frac{\rho \cdot A_{1} L_{1}}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
& =\frac{\rho(2 A)(L / 2)}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
{\left[M_{1}\right] } & =\frac{\rho A L}{12}\left[\begin{array}{ll}
4 & 2 \\
2 & 4
\end{array}\right]
\end{aligned}
$$

For element (2),

$$
\begin{aligned}
{\left[M_{2}\right] } & =\frac{\rho \cdot A_{2} L_{2}}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
& =\frac{\rho(A)(L / 2)}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
{\left[M_{2}\right] } & =\frac{\rho A L}{12}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right]
\end{aligned}
$$

Global mass matrix,

$$
\begin{aligned}
& {[M]=\left[M_{1}\right]+\left[M_{2}\right]} \\
& {[M]=\left[\begin{array}{lll}
4 & 2 & 0 \\
2 & 6 & 1 \\
0 & 1 & 2
\end{array}\right] \times \frac{\rho A L}{12}}
\end{aligned}
$$

Finite element equation for the longitudinal vibrations of a bar element is given by,

$$
\begin{aligned}
& {\left[[K]-\omega^{2}[M]\right]\{\delta\}=\{F\}} \\
& {\left[\frac{2 A E}{L}\left[\begin{array}{ccc}
2 & -2 & 0 \\
-2 & 3 & -1 \\
0 & -1 & 1
\end{array}\right]-\omega^{2} \cdot \frac{\rho A L}{12}\left[\begin{array}{lll}
4 & 2 & 0 \\
2 & 6 & 1 \\
0 & 1 & 2
\end{array}\right]\right]\left\{\begin{array}{l}
u_{1} \\
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
F_{1} \\
F_{2} \\
F_{3}
\end{array}\right\}}
\end{aligned}
$$

Applying the boundary conditions,

$$
\begin{aligned}
& F_{1}=0, F_{2}=0 . F_{3}=0 \\
& {\left[\frac{2 A E}{L}\left[\begin{array}{ccc}
2 & -2 & 0 \\
-2 & 3 & -1 \\
0 & -1 & 1
\end{array}\right]-\omega^{2} \cdot \frac{\rho A L}{12}\left[\begin{array}{lll}
4 & 2 & 0 \\
2 & 6 & 1 \\
0 & 1 & 2
\end{array}\right]\left[\left\{\begin{array}{l}
u_{1} \\
u_{1} \\
u_{2} \\
u_{3}
\end{array}\right]=0\right.\right.}
\end{aligned}
$$

**One can eliminate the 1st row and column to reduce computation

Characteristic equation is given by,

$$
\begin{aligned}
& \left|[K]-\omega^{2}[M]\right|=0 \\
& \left|\frac{2 A E}{L}\left[\begin{array}{ccc}
2 & -2 & 0 \\
-2 & 3 & -1 \\
0 & -1 & 1
\end{array}\right]-\omega^{2} \cdot \frac{\rho A L}{12}\left[\begin{array}{lll}
4 & 2 & 0 \\
2 & 6 & 1 \\
0 & 1 & 2
\end{array}\right]\right|=0
\end{aligned}
$$

Divide by $\frac{2 A E}{L}$ on both sides.

$$
\begin{aligned}
& \left|\left[\begin{array}{ccc}
2 & -2 & 0 \\
-2 & 3 & -1 \\
0 & -1 & 1
\end{array}\right]-\frac{\frac{\rho A L}{12}}{\frac{2 A E}{L}} \omega^{2}\left[\begin{array}{lll}
4 & 2 & 0 \\
2 & 6 & 1 \\
0 & 1 & 2
\end{array}\right]\right|=0 \\
& \left|\left[\begin{array}{ccc}
2 & -2 & 0 \\
-2 & 3 & -1 \\
0 & -1 & 1
\end{array}\right]-\frac{\rho L^{2} \omega^{2}}{24 E} \omega^{2}\left[\begin{array}{lll}
4 & 2 & 0 \\
2 & 6 & 1 \\
0 & 1 & 2
\end{array}\right]\right|=0
\end{aligned}
$$

Let,

$$
\begin{gathered}
\alpha=\frac{\rho L^{2} \omega^{2}}{24 E} \\
\left|\left[\begin{array}{ccc}
2 & -2 & 0 \\
-2 & 3 & -1 \\
0 & -1 & 1
\end{array}\right]-\alpha\left[\begin{array}{ccc}
4 & 2 & 0 \\
2 & 6 & 1 \\
0 & 1 & 2
\end{array}\right]\right|=0 \\
\left|\left[\begin{array}{ccc}
2 & -2 & 0 \\
-2 & 3 & -1 \\
0 & -1 & 1
\end{array}\right]-\left[\begin{array}{ccc}
4 \alpha & 2 \alpha & 0 \\
2 \alpha & 6 \alpha & \alpha \\
0 & \alpha & 2 \alpha
\end{array}\right]\right|=0
\end{gathered}
$$

$$
\begin{aligned}
& \left.\left\lvert\, \begin{array}{ccc}
(2-4 \alpha) & (-2-2 \alpha) & 0 \\
(-2-2 \alpha) & (3-6 \alpha) & (-1-\alpha) \\
0 & (-1-\alpha) & (1-2 \alpha)
\end{array}\right.\right]=0 \\
& (2-4 \alpha)[(3-6 \alpha)(1-2 \alpha)-(-1-\alpha)(-1-\alpha)] \\
& \quad-(-2-2 \alpha)[(2-4 \alpha)(1-2 \alpha)]=0 \\
& (2-4 \alpha)\left[3-6 \alpha-6 \alpha+12 \alpha^{2}-1-\alpha-\alpha-\alpha^{2}\right] \\
& \quad-(-2-2 \alpha)\left[\left(2-4 \alpha-4 \alpha+8 \alpha^{2}\right)\right]=0
\end{aligned}
$$

$18 \alpha(1-2 \alpha)(\alpha-2)=0$
$\alpha=0, \frac{1}{2}, 2$

The roots of above equation gives the natural frequencies of the bar,

When,

$$
\begin{aligned}
& \alpha=0 \\
& \alpha \\
&=\frac{\rho L^{2} \omega_{1}^{2}}{2 E} \\
& \omega_{1}^{2}=0 \\
& \omega_{1}=0 \mathrm{rad} / \mathrm{sec}
\end{aligned}
$$

When,

$$
\begin{aligned}
& \alpha=\frac{1}{2} \\
& \alpha=\frac{\rho L^{2} \omega_{2}^{2}}{24 E} \\
& \frac{1}{2}=\frac{\rho L^{2} \omega_{2}^{2}}{24 E} \\
& \omega_{2}^{2}=\frac{12 E}{\rho L^{2}} \\
& \omega_{2}=\sqrt{\frac{12 E}{\rho L^{2}}} \\
& \omega_{2}=3.464 \sqrt{\frac{E}{\rho L^{2}}} \mathrm{rad} / \mathrm{sec}
\end{aligned}
$$

When,

$$
\begin{aligned}
& \alpha=2 \\
& \alpha=\frac{\rho L^{2} \omega_{3}^{2}}{24 E} \\
& 2=\frac{\rho L^{2} \omega_{3}^{2}}{24 E} \\
& \omega_{3}^{2}=\frac{48 E}{\rho L^{2}} \\
& \omega_{3}=\sqrt{\frac{48 E}{\rho L^{2}}} \\
& \omega_{3}=6.928 \sqrt{\frac{E}{\rho L^{2}}} \mathrm{rad} / \mathrm{sec}
\end{aligned}
$$

Therefore, the natural frequencies of longitudinal vibrations of given stepped bar is given by,

$$
\begin{aligned}
& \omega_{1}=0 \mathrm{rad} / \mathrm{sec} \\
& \omega_{2}=3.464 \sqrt{\frac{E}{\rho L^{2}}} \mathrm{rad} / \mathrm{sec} \\
& \omega_{3}=6.928 \sqrt{\frac{E}{\rho L^{2}}} \mathrm{rad} / \mathrm{sec}
\end{aligned}
$$

## Equations of Motion Using Lagrange's Approach

We define the Lagrangean by

$$
L=T-\Pi
$$

If $T$ represents the kinetic energy of a system and $\pi$ represents potential energy, then Lagrange's equations of motion

$$
\frac{d}{d t}\left(\frac{\partial L}{\partial \dot{q}_{i}}\right)-\frac{\partial L}{\partial q_{i}}=0 \quad \frac{d}{d t}\left(\frac{\partial T}{\partial \dot{q}}\right)-\frac{\partial T}{\partial q}+\frac{\partial \pi}{\partial q}=\left.F\right|_{q}
$$

where $\left.F\right|_{q}$ represents the generalised force in the coordinate $q$ and $\dot{q}=d q / d t$.

$$
\begin{aligned}
& \text { Lagrange's equations of motion, we obtain } \\
& \qquad \begin{array}{l}
{[M]\{\ddot{\delta}\}+[K]\{\delta\}-\{F\}=\{0\}} \\
\\
{[M]\{\ddot{\delta}\}+[K]\{\delta\}=\{F\}}
\end{array}
\end{aligned}
$$

## Solid Body with Distributed Mass

we express
displacement $\mathbf{u}$ in terms of the nodal displacements $\mathbf{q}$,

$$
\dot{\mathbf{u}}=\left[\begin{array}{lll}
\dot{u} & \dot{v} & \dot{w}
\end{array}\right]^{\mathrm{T}}
$$

$$
\mathbf{u}=\mathbf{N q}
$$

using shape functions $\mathbf{N}$

In dynamic analysis, the elements of $\mathbf{q}$ are dependent on time, while $\mathbf{N}$ represents (spatial) shape functions.

The velocity vector is then given by

$$
\dot{\mathbf{u}}=\mathbf{N} \dot{\mathbf{q}}
$$

The kinetic energy $\quad T=\frac{1}{2} \int_{V} \dot{\mathbf{u}}^{\mathrm{T}} \dot{\mathbf{u}} \rho d V$

$$
T_{e}=\frac{1}{2} \dot{\mathbf{q}}^{\mathrm{T}}\left[\int_{e} \rho \mathbf{N}^{\mathrm{T}} \mathbf{N} d V\right] \dot{\mathbf{q}}
$$

$$
\mathbf{m}^{e}=\int_{e} \rho \mathbf{N}^{\mathrm{T}} \mathbf{N} d V
$$

This mass matrix is consistent with the shape functions chosen and is called the consistent mass matrix.


## Figure: Bar Element

The above figure shows a bar element of length ' $l$ '. Let $u_{1}, u_{2}$ are the nodal displacements at nodal points 1,2 .

$$
N_{1}=\frac{l-x}{l}, N_{2}=\frac{x}{l}
$$

Shape functions for a $1-D$ bar element with two nodes is given by,

$$
\mathbf{m}^{e}=\int_{e} \rho \mathbf{N}^{\mathrm{T}} \mathbf{N} d V
$$

Mass matrix,

$$
\begin{aligned}
& {[M]=\int_{v} \rho \cdot[N]^{T}[N] d V=\rho \int_{0}^{l}\left[\begin{array}{l}
N_{1} \\
N_{2}
\end{array}\right]\left[\begin{array}{ll}
N_{1} & N_{2}
\end{array}\right] A \cdot d x} \\
& =\rho A \int_{0}^{l}\left[\begin{array}{l}
N_{1} \\
N_{2}
\end{array}\right]\left[\begin{array}{ll}
N_{1} & N_{2}
\end{array}\right] \cdot d x \\
& =\rho A \int_{0}^{l}\left[\frac { l - x } { \frac { l } { x } } \left[\left[\frac{l-x}{l} \frac{x}{l}\right] d x\right.\right. \\
& =\rho A \int_{0}^{l}\left[\begin{array}{cc}
\left(\frac{l-x}{l}\right)^{2} & \frac{x}{l}-\frac{x^{2}}{l^{2}} \\
\frac{x}{l}-\frac{x^{2}}{l^{2}} & \frac{x^{2}}{l^{2}}
\end{array}\right] d x \\
& =\rho A\left[\begin{array}{cc}
\frac{\left(1-\frac{x}{l}\right)^{2}}{3\left(\frac{-1}{l}\right)} & \frac{x^{2}}{2 l}-\frac{x^{3}}{3 l^{2}} \\
\frac{x^{2}}{2 l}-\frac{x^{3}}{3 l^{2}} & \frac{x^{3}}{3 l^{2}}
\end{array}\right]_{0} \\
& =\rho A\left[\begin{array}{cc}
\frac{\left(1-\frac{l}{l}\right)^{2}}{3\left(\frac{-1}{l}\right)} & \frac{l^{2}}{2 l}-\frac{l^{3}}{3 l^{2}} \\
\frac{l^{2}}{2 l}-\frac{l^{3}}{3 l^{2}} & \frac{l^{3}}{3 l^{2}}
\end{array}\right] \\
& =\rho A\left[\begin{array}{ll}
\frac{l}{3} & \frac{l}{6} \\
\frac{l}{6} & \frac{l}{3}
\end{array}\right]=\frac{\rho A l}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right]
\end{aligned}
$$

Bar element $\quad N_{1}=\frac{l-x}{l}, N_{2}=\frac{x}{l}$

$$
[m]^{e}=\int_{v} \rho[N]^{T}[N] d v=\rho A \int_{0}^{\ell}\left[\begin{array}{c}
1-\frac{x}{\ell} \\
\frac{x}{\ell}
\end{array}\right]\left[\left(1-\frac{x}{\ell}\right)\left(\frac{x}{\ell}\right)\right] d x=\frac{\rho A \ell}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right]
$$

$$
\text { Beam element } \quad \begin{aligned}
& N_{1}=1-3 x^{2} / \ell^{2}+2 x^{3} / \ell^{3}, \quad N_{2}=x-2 x^{2} / \ell+x^{3} / \ell^{2} \\
& N_{3}=3 x^{2} / \ell^{2}-2 x^{3} / \ell^{3}, \quad N_{4}=-x^{2} / \ell+x^{3} / \ell^{2}
\end{aligned}
$$

$$
[m]^{e}=\int_{v} \rho[N]^{T}[N] d v=\rho A \int_{0}^{\ell}\left[\begin{array}{l}
N_{1} \\
N_{2} \\
N_{3} \\
N_{4}
\end{array}\right]\left[\begin{array}{llll}
N_{1} & N_{2} & N_{3} & N_{4}
\end{array}\right] d x=\frac{\rho A \ell}{420}\left[\begin{array}{ccc}
156 & 22 \ell & 4 \ell^{2} \\
54 & 13 \ell & 156 \\
-13 \ell & -3 \ell^{2} & -22 \ell
\end{array}\right.
$$

Determine the natural frequencies and mode shapes of a stepped bar as shown in figure, using the characteristic polynomial technique. Assume $\mathrm{E}=250 \mathrm{GPa}$ and density is $7850 \mathrm{~kg} / \mathrm{m}^{3}$.


Figure
Given that,
Young's modulus, $E=250 \mathrm{GPa}=250 \times 10^{9} \mathrm{~Pa}$
Determine the natural frequencies of a stepped bar as shown in the figure. ( $\mathrm{E}=250 \mathrm{GPa}$, density 7850 kg/m^3

Density, $\rho=7850 \mathrm{~kg} / \mathrm{m}^{3}$
Evaluation is done using characteristic polynomial technique.


Figure (1): Bar Element Model

Finite element equation of a bar element is given by,

## Stiffness Matrix

$$
\left[[K]-\omega^{2}[M]\right]\{\delta\}=\{F\}
$$

For element (1),

$$
\left.\begin{array}{rl}
{[1),} \\
{\left[K_{1}\right]=} & \frac{E A_{1}}{L_{1}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
= & \frac{250 \times 10^{9} \times 400 \times 10^{-6}}{0.3}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
= & 333.333 \times 10^{6}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
1 & 2
\end{array}\right] \begin{aligned}
& 1 \\
& \\
& \quad\left[K_{1}\right]=10^{6}\left[\begin{array}{cc}
333.333 & -333.333 \\
-333.333 & 333.333
\end{array}\right]
\end{aligned}
$$

For element (2),

$$
\begin{array}{rl}
{\left[K_{2}\right]} & =\frac{E A_{2}}{L_{2}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
& =\frac{250 \times 10^{9} \times 300 \times 10^{-6}}{0.3}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
2 & 3 \\
{\left[K_{2}\right]} & =10^{6}\left[\begin{array}{cc}
250 & -250 \\
-250 & 250
\end{array}\right]_{3}^{2}
\end{array}
$$

Global stiffness matrix,

$$
\therefore \quad[K]=\left[K_{1}\right]+\left[K_{2}\right]
$$

$$
\left.[K]=10^{6}\left[\begin{array}{ccc}
333.333 & -333.333 & 0 \\
-333.333 & 583.333 & -250 \\
0 & -250 & 250
\end{array}\right] \begin{array}{l}
1 \\
2 \\
3
\end{array}\right]
$$

## Mass Matrix

For element (1),

$$
\begin{aligned}
{\left[M_{1}\right] } & =\frac{\rho A_{1} L_{1}}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
& =\frac{7850 \times 400 \times 10^{-6} \times 0.3}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
& =0.157\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
{\left[M_{1}\right] } & =\left[\begin{array}{ll}
0.314 & 0.157 \\
0.157 & 0.314
\end{array}\right] 1
\end{aligned}
$$

For element (2),

$$
\begin{aligned}
{\left[M_{2}\right] } & =\frac{\rho A_{2} l_{2}}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
& =\frac{7850 \times 300 \times 10^{-6} \times 0.3}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
& =0.117\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
{\left[\begin{array}{ll}
3
\end{array}\right] } & =\left[\begin{array}{cc}
0.234 & 0.117 \\
0.117 & 0.234
\end{array}\right] 2
\end{aligned}
$$

Global mass matrix,

$$
\begin{array}{rl}
\because & {[M]=\left[M_{1}\right]+\left[M_{2}\right]} \\
1 & 2
\end{array} 3
$$

$$
\left[[K]-\omega^{2}[M]\right]\{u\}=\{F\}
$$

$$
\left[10^{6}\left[\begin{array}{ccc}
333.333 & -333.333 & 0 \\
-333.333 & 583.333 & -250 \\
0 & -250 & 250
\end{array}\right]-\omega^{2}\left[\begin{array}{ccc}
0.314 & 0.157 & 0 \\
0.157 & 0.548 & 0.117 \\
0 & 0.117 & 0.234
\end{array}\right]\right]\left\{\begin{array}{c}
u_{1} \\
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
0 \\
10^{4}
\end{array}\right\} \longrightarrow \text { Eq-1 }
$$

Apply the boundary conditions to the above equation,

$$
u_{1}=0
$$

Eliminating first row and first column of both the matrixes,

$$
\left[10^{6}\left[\begin{array}{cc}
583.333 & -250 \\
-250 & 250
\end{array}\right]-\lambda\left[\begin{array}{cc}
0.548 & 0.117 \\
0.117 & 0.234
\end{array}\right]\right]\left\{\begin{array}{c}
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
10^{4}
\end{array}\right\}
$$

Let, $\lambda=\omega^{2}$.
Characteristic equation is given by,

$$
\begin{aligned}
& |[\mathrm{K}]-\lambda[\mathrm{M}]|=0 \\
& \left|10^{6}\left[\begin{array}{cc}
583.333 & -250 \\
-250 & 250
\end{array}\right]-\lambda\left[\begin{array}{ll}
0.548 & 0.117 \\
0.117 & 0.234
\end{array}\right]\right|=0
\end{aligned}
$$

$$
\begin{aligned}
& \left|10^{6}\left[\begin{array}{cc}
583.333 & -250 \\
-250 & 250
\end{array}\right]-\left[\begin{array}{cc}
0.548 \lambda & 0.117 \lambda \\
0.117 \lambda & 0.234 \lambda
\end{array}\right]\right|=0 \\
& \left|\left[\begin{array}{cc}
583.333 \times 10^{6}-0.548 \lambda & -250 \times 10^{6}-0.117 \lambda \\
-250 \times 10^{6}-0.117 \lambda & 250 \times 10^{6}-0.234 \lambda
\end{array}\right]\right|=0 \\
& \quad\left[\left(583.333 \times 10^{6}-0.548 \lambda\right)\left(250 \times 10^{6}-0.234 \lambda\right)-\left(-250 \times 10^{6}-0.117 \lambda\right)\left(-250 \times 10^{6}-0.117 \lambda\right)\right]=0 \\
& \left(583.333 \times 10^{6} \times 250 \times 10^{6}\right)-\left(583.333 \times 10^{6} \times 0.234 \lambda\right)-\left(0.548 \lambda \times 250 \times 10^{6}\right)+(0.548 \lambda \times 0.234 \lambda) \\
& \quad-\left(250 \times 10^{6} \times 250 \times 10^{6}\right)-\left(250 \times 10^{6} \times 0.117 \lambda\right)-\left(0.117 \lambda \times 250 \times 10^{6}\right)-(0.117 \lambda \times 0.117 \lambda)=0 \\
& 1.45 \times 10^{17}-136.4 \times 10^{6} \lambda-137 \times 10^{6} \lambda+0.128 \lambda^{2}-6.25 \times 10^{16}-29.25 \times 10^{6} \lambda-29.25 \times 10^{6} \lambda-0.0136 \lambda^{2}=0 \\
& 0.115 \lambda^{2}-331.9 \times 10^{6} \lambda+8.25 \times 10^{16}=0
\end{aligned}
$$

$$
\begin{aligned}
& \lambda_{1}=2611368423 \\
& \lambda_{2}=274718533
\end{aligned}
$$

$$
\begin{aligned}
& \lambda_{1}=2611368423 \\
& \lambda_{1}=\omega_{1}^{2}=2611368423 \\
& \omega_{1}=\sqrt{2611368423} \\
& \omega_{1}=51101.55 \mathrm{rad} / \mathrm{sec}
\end{aligned}
$$

When,

$$
\begin{aligned}
& \lambda_{2}=274718533 \\
& \lambda_{2}=\omega_{2}^{2}=274718533 \\
& \omega_{2}=\sqrt{274718533} \\
& \omega_{2}=16574.63 \mathrm{rad} / \mathrm{sec}
\end{aligned}
$$

Natural frequency of element (1) is, $f_{1}=\frac{\omega_{1}}{2 \pi}$

$$
\begin{aligned}
& f_{1}=\frac{51101.55}{2 \pi} \\
& f_{1}=8133.064 \mathrm{~Hz}
\end{aligned}
$$

Natural frequency of element (2),

$$
\begin{aligned}
f_{2} & =\frac{\omega_{2}}{2 \pi} \\
& =\frac{16574.63}{2 \pi} \\
f_{2} & =2637.93 \mathrm{~Hz}
\end{aligned}
$$

## Mode Shapes

From equation (2)

$$
\begin{aligned}
& {\left[10^{6}\left[\begin{array}{cc}
583.33 & -250 \\
-250 & 250
\end{array}\right]-\lambda\left[\begin{array}{cc}
0.548 & 0.117 \\
0.117 & 0.237
\end{array}\right]\left[\left\{\begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right]=\left\{\begin{array}{c}
0 \\
10^{4}
\end{array}\right\}\right.\right.} \\
& {\left[\begin{array}{cc}
583.33 \times 10^{6}-0.548 \lambda & -250 \times 10^{6}-0.117 \lambda \\
-250 \times 10^{6}-0.117 \lambda & 250 \times 10^{6}-0.234 \lambda
\end{array}\right]\left\{\begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
10^{4}
\end{array}\right\}}
\end{aligned}
$$

For $\lambda=\lambda_{1}$

$$
\lambda_{1}=2611368423
$$

$$
\left\{\begin{array}{cc}
583.33 \times 10^{6}-0.548(2611368423) & -250 \times 10^{6}-0.117(2611368423) \\
-250 \times 10^{6}-0.117(2611368423) & 250 \times 10^{6}-0.234(2611368423)
\end{array}\right\}\left\{\begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
10^{4}
\end{array}\right\}
$$

$$
\left.\left.\left[\begin{array}{cc}
-847699895.8 & -555530105.5 \\
-555530105.5 & -361060211
\end{array}\right] \right\rvert\, \begin{array}{c}
u_{2} \\
u_{3} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
10^{4}
\end{array}\right\}
$$

$$
\left.\begin{array}{l}
{\left[\begin{array}{ll}
-847699895.8 & -555530105.5 \\
-555530105.5 & -361060211
\end{array}\right]\left[\begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
10
\end{array}\right\}} \\
\hline
\end{array}\right\},
$$

|  | u2 | u3 |
| :---: | :---: | :---: |
| solution | -0.00218 | 0.00333 |

1st Mode shape

## Mode Shapes

## For second mode shape

From equation (2)

$$
\begin{aligned}
& {\left[10^{6}\left[\begin{array}{cc}
583.33 & -250 \\
-250 & 250
\end{array}\right]-\lambda\left[\begin{array}{cc}
0.548 & 0.117 \\
0.117 & 0.237
\end{array}\right]\right]\left\{\begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
10^{4}
\end{array}\right\}} \\
& {\left[\begin{array}{cc}
583.33 \times 10^{6}-0.548 \lambda & -250 \times 10^{6}-0.117 \lambda \\
-250 \times 10^{6}-0.117 \lambda & 250 \times 10^{6}-0.234 \lambda
\end{array}\right]\left\{\begin{array}{c}
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
10^{4}
\end{array}\right\}}
\end{aligned}
$$

For $\lambda=\lambda_{2}$

$$
\lambda_{2}=274718533
$$

$$
\begin{aligned}
& {\left[\begin{array}{cc}
583.33 \times 10^{6}-0.548(274718533) & -250 \times 10^{6}-0.117(274718533) \\
-250 \times 10^{6}-0.117(274718533) & 250 \times 10^{6}-0.234(274718533)
\end{array}\right]\left[\begin{array}{l}
u_{2} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
10^{4}
\end{array}\right\}} \\
& {\left[\begin{array}{cc}
432784243.9 & -282142068.4 \\
-282142068.4 & 185715863.3
\end{array}\right]\left\{\left\{\begin{array}{c}
u_{2} \\
u_{3} \\
u_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0 \\
10^{4}
\end{array}\right\}\right.}
\end{aligned}
$$

$\left(\begin{array}{ll}43278 & -28214 \\ -28214 & 18571\end{array}\right)\left\{\begin{array}{l}u_{2} \\ u_{3}\end{array}\right\}=\left\{\begin{array}{l}0 \\ 1\end{array}\right\}$

|  | u2 | u3 |
| :---: | :---: | :---: |
| solution | 0.00367 | 0.00563 |



Figure

## Second Mode



$$
\mathbf{K} \mathbf{U}=\omega^{2} \mathbf{M U} \quad[\mathbf{K}-\lambda \mathbf{M}]\{\mathbf{U}\}=\{0\}
$$

This is the generalized eigenvalue problem

$$
\mathbf{K} \mathbf{U}=\lambda \mathbf{M U}
$$

where $\mathbf{U}$ is the eigenvector, representing the vibrating mode, corresponding to the eigenvalue $\lambda$. The eigenvalue $\lambda$ is the square of the circular frequency $\omega$. The frequency $f$ in hertz (cycles per second) is obtained from $f=\omega /(2 \pi)$.

## EVALUATION OF EIGENVALUES AND EIGENVECTORS

The generalized problem in free vibration is that of evaluating an eigenvalue $\lambda\left(=\omega^{2}\right)$, which is a measure of the frequency of vibration together with the corresponding eigenvector $\mathbf{U}$ indicating the mode shape, as in

$$
\mathbf{K U}=\lambda \mathbf{M U}
$$

$[\mathrm{K}-\lambda \mathrm{M}]\{\mathrm{U}\}=\{0\}$

```
\lambda -------- Eigen Value (Frequency)
U or u -------- Eigen Vector (Mode shape)
```


## Eigenvalue-Eigenvector Evaluation

The eigenvalue-eigenvector evaluation procedures fall into the following basic categories:

1. Characteristic polynomial technique
2. Vector iteration methods
3. Transformation methods

Characteristic Polynomial. From Eq. 11.38, we have

$$
(\mathbf{K}-\lambda \mathbf{M}) \mathbf{U}=\mathbf{0}
$$

If the eigenvector is to be nontrivial, the required condition is

$$
\operatorname{det}(\mathbf{K}-\lambda \mathbf{M})=\mathbf{0}
$$

This represents the characteristic polynomial in $\lambda$.

What is study state heat transfer analysis? Write its governing Equation?

Steady state heat transfer is defined as the temperature at any point in the medium does not change with time.

For a one dimensional steady state heat transfer,

$$
K \cdot \frac{d^{2} T}{d x^{2}}+q=0
$$

$K$ - Thermal conductivity
$T$ - Temperature
$q$ - Internal heat source per unit volume

| Steady State Conduction |  |
| :---: | :--- |
| 1. | In steady state conduction, the temperature at any <br> point in the medium does not change with time. <br> It is due to the rate of heat conducted into the medium is <br> equal to the rate of heat conducted out of the medium. <br> It is a function of space coordinates and is given by, <br> $T=T(x, y, z)$ |
| 4. One dimensional heat conduction is given by, |  |
| $\frac{\partial}{\partial x}\left[k \frac{\partial T}{\partial x}\right]+q=0$ |  |

## 3D Conduction heat transfer

General 3D conduction Equation:

$$
\frac{\partial}{\partial x}\left(k \frac{\partial T}{\partial x}\right)+\frac{\partial}{\partial y}\left(k \frac{\partial T}{\partial y}\right)+\frac{\partial}{\partial z}\left(k \frac{\partial T}{\partial z}\right)+\dot{q}=\rho c \frac{\partial T}{\partial \tau}
$$

For constant conductivity:

$$
\begin{aligned}
\frac{\partial^{2} T}{\partial x^{2}} & +\frac{\partial^{2} T}{\partial y^{2}}+\frac{\partial^{2} T}{\partial z^{2}}+\frac{\dot{q}}{k}=\frac{1}{\alpha} \frac{\partial T}{\partial \tau} \\
\alpha & =k / \rho c \\
& =\text { Thermal diffusivity of a material }
\end{aligned}
$$

Q. Give the finite element equation for a one dimensional heat conduction element.

Ans: The finite element equation for a one dimensional heat conduction element is given by,

$$
\left.\begin{array}{rl} 
& \{F\}=\left[K_{c}\right]\{T\} \\
\{F\} & - \text { Force vector } \\
& =\left\{\begin{array}{l}
F_{1} \\
F_{2}
\end{array}\right\} \text { for a two noded element }
\end{array}\right\} \begin{aligned}
{\left[K_{c}\right] } & - \text { Stiffness matrix in case of heat conduction } \\
& =\frac{K A}{l}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
\{T\} & - \text { Nodal temperature vector } \\
& =\left\{\begin{array}{l}
T \\
1 \\
T_{2} \\
2
\end{array}\right\} \text { for a two noded element }
\end{aligned}
$$



## Similar to structural problems

- Formulation of element stiffness matrices
- Assembly of global stiffness matrix [K]
- formulation of load vector \{F\}
- Solution of $[K]\{u\}=\{F\}$ to get nodal displacements $\{u\}$


## Heat Transfer Rate Due to Conduction

$$
Q=-k A \cdot \frac{\Delta T}{L} \mathrm{Watts}
$$

Heat Transfer Rate Due to Convection

$$
Q=h A\left(T_{s}-T_{\infty}\right) \text { Watts }
$$

Where,
$h$-Heat transfer coefficient
$T_{s}$ - Surface temperature
$T_{s \infty}$-Ambient temperature.
Heat Transfer Rate Due to Radiation

$$
Q=\sigma A(\Delta T)^{4} \text { Watts }
$$

Where,
$\sigma$-Stefan-Boltzmann constant

$$
\sigma=5.67 \times 10^{-8} \mathrm{~W} / \mathrm{m}^{2} \mathrm{~K}^{4} .
$$

Where,

$$
\begin{aligned}
& k \text { - Thermal conductivity } \\
& A \text { - Surface area } \\
& \Delta T \text { - Temperature difference } \\
& L \text { - Length. }
\end{aligned}
$$

Derive the stiffness matrix for a one dimensional heat conduction element.


## Figure: Bar element

The above figure shows a bar element of length ' $?$. Let $T_{1}, T_{2}$ are the temperature at nodes 1,2 and ' $K$ ' be the thermal conductivity of the bar element,

Stiffness matrix,

$$
\begin{equation*}
[K]=\int_{V}[B]^{T}[D][B] d V \tag{1}
\end{equation*}
$$

Where,

$$
\begin{aligned}
& N_{1}=\frac{l-x}{l} \\
& N_{2}=\frac{x}{l}
\end{aligned}
$$

Strain-Displacement matrix,

$$
\begin{aligned}
{[B] } & \left.=\left\lvert\, \begin{array}{ll}
\frac{d N_{1}}{d x} & \frac{d N_{2}}{d x}
\end{array}\right.\right] \\
& =\left\lceil\frac{d}{d x}\left[\frac{l-x}{l}| | \frac{d}{d x}\left(\frac{x}{l}\right)\right]=\left[\frac{-1}{l} \frac{1}{l}\right]\right. \\
{[B]^{r} } & \left.=\left|\frac{-1}{l}\right| \frac{1}{l} \right\rvert\,
\end{aligned}
$$

[B] - Strain displacement matrix
[D] - Stress - strain matrix
For a one dimensional bar element,
Temperature function, $T=N_{j} T_{1}+N_{2} T_{2}$

For a one dimensional heat conduction,

$$
[D]=K
$$

From equation (1),

$$
\begin{aligned}
& {[K]=\int_{v}[B]^{T}[D][B] d V} \\
& {\left[K_{c}\right]=\int_{0}^{l}\left[\begin{array}{l}
\frac{-1}{l} \\
\frac{1}{l}
\end{array}\right] \cdot K \cdot\left[\frac{-1}{l} \frac{1}{l}\right] d V}
\end{aligned}
$$

$\left[K_{6}\right]$ - Stiffiness matrix in case of conduction

$$
\begin{aligned}
& =\int_{0}^{l}\left[\begin{array}{cc}
\frac{1}{l^{2}} & \frac{-1}{l^{2}} \\
\frac{-1}{l^{2}} & \frac{1}{l^{2}}
\end{array}\right] K \cdot d V \\
& =\int_{0}^{l}\left[\begin{array}{cc}
\frac{1}{l^{2}} & \frac{-1}{l^{2}} \\
\frac{-1}{l^{2}} & \frac{1}{l^{2}}
\end{array}\right] K \cdot A \cdot d x \\
& =K A\left[\begin{array}{cc}
\frac{1}{l^{2}} & \frac{-1}{l^{2}} \\
\frac{-1}{l^{2}} & \frac{1}{l^{2}}
\end{array} \int_{0}^{l} 1 \cdot d x \quad=\frac{K A}{l}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]\right.
\end{aligned}
$$

## FE Equation for 1D heat conduction element

The finite element equation for a one dimensional heat conduction element is given by,

$$
\begin{array}{r}
\{F\}=\left[K_{c}\right]\{T\} \\
\{F\} \text { - Force vector }
\end{array}
$$

Thermal

$$
=\left\{\begin{array}{l}
F_{1} \\
F_{2} \\
2
\end{array}\right\} \text { for a two noded element }
$$ stiffness $\left[K_{c}\right]$ - Stiffness matrix in case of heat conduction matrix

$$
=\frac{K A}{l}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]
$$

$\{T\}$ - Nodal temperature vector


$$
=\left\{\begin{array}{l}
T_{1} \\
T_{2}
\end{array}\right\} \text { for a two noded element }
$$

FE Equation for 1D heat conduction and Convection element

$$
[K]\{T\}=\{F\}
$$

Stiffness Matrix in Case of a 1D Fin Pro Thermal Stiffness matrix, $[\mathrm{K}]=[\mathrm{Kc}]+[\mathrm{Khe}]$

$$
\begin{aligned}
& {[k]=\frac{K A}{l}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]+h A\left[\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right]} \\
& {[\mathrm{Kc}]=\frac{A K}{l}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]-\text { Conduction matrix }}
\end{aligned}
$$

$$
[\text { Khe }]=h A\left[\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right] \text { - Convection matrix }
$$

$$
h T_{\infty} A\left[\begin{array}{l}
0 \\
1
\end{array}\right\} \text { - Thermal load matrix }
$$

A $\longrightarrow$ Area of the wall
Thermal conductivity of wall
Length of the wall
Heat transfer coefficient
Atmospheric air temperature
$-\mathrm{m}^{2}$

- W/mK
- m
- W/m²K
- K


In case of convection (end), stiffness matrix is given by,

$$
\begin{equation*}
[\text { Khe }]=\iint_{A} h[N]^{T}[N] d A \tag{3}
\end{equation*}
$$

$$
\text { at } x=l \text {, at the end of element }
$$

$N$ - Shape functions
$h$ - Heat transfer coefficient, W/m²K

$$
\begin{aligned}
{[N] } & =\left[\begin{array}{ll}
N_{1} & N_{2}
\end{array}\right] \\
& =\left[\frac{l-x}{l} \frac{x}{l}\right]
\end{aligned}
$$

From equation (3),

$$
\begin{aligned}
{[\text { Khe }] } & =\iint_{A} h \cdot[N]^{T}[N] \cdot d A \\
& =\iint_{A} h \cdot\left[\begin{array}{l}
0 \\
1
\end{array}\right]\left[\begin{array}{ll}
0 & 1
\end{array}\right] d A \\
& =h\left[\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right] \iint_{A} d A \\
{[\text { Khe }] } & =h A\left[\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right]
\end{aligned}
$$

$$
\begin{aligned}
& {[M]=\left[\begin{array}{ll}
N_{1} & N_{2}
\end{array}\right]=\left[\begin{array}{ll}
0 & 1
\end{array}\right]} \\
& {[M]^{T}=\left[\begin{array}{l}
N_{1} \\
N_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
1
\end{array}\right]}
\end{aligned}
$$

Stiffness matrix,

$$
\begin{aligned}
& {[K]=\left[K_{c}\right]+[\mathrm{Khe}]} \\
& {[K]=\frac{A K}{l}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]+h A\left[\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right]}
\end{aligned}
$$

Convective force vector at the free end,

$$
\{\text { Fhe }\}=h T_{\infty} A\left[\begin{array}{l}
N_{1} \\
N_{2}
\end{array}\right]
$$

Where,
$T_{\infty}$ - Fluid temperature

$$
\{\text { Fhe }\}=h T_{\infty} A\left[\begin{array}{l}
0 \\
1
\end{array}\right]
$$

From equation (1),

$$
\{\text { Fhe }\}=[K]\{T\}
$$

$$
h T_{\infty} A\left[\begin{array}{l}
0 \\
1
\end{array}\right]=\left[\frac{A K}{l}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]+h A\left[\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right]\left[\begin{array}{l}
T_{1} \\
T_{2}
\end{array}\right\}\right.
$$

A composite slab consists of three materials of different conductivities is $20 \mathrm{~W} / \mathrm{mK}, 30 \mathrm{~W} / \mathrm{mK}$ and $50 \mathrm{~W} / \mathrm{mK}$ of thickness $0.3 \mathrm{~m}, 0.15 \mathrm{~m}$ and 0.15 m respectively. The outer surface is $20^{\circ} \mathrm{C}$ and the inner surface is exposed to the convective heat transfer coefficient of $25 \mathrm{~W} / \mathrm{m}^{2} \mathrm{~K}$ at $800^{\circ} \mathrm{C}$. Determine the temperature distribution within the wall.


Given that,
For material-1,
Thermal conductivity, $K_{1}=20 \mathrm{~W} / \mathrm{mK}$
Thickness, $L_{1}=0.3 \mathrm{~m}$
For material-2,
Thermal conductivity, $K_{2}=30 \mathrm{~W} / \mathrm{mK}$
Thickness, $L_{2}=0.15 \mathrm{~m}$
For material-3,
Thermal conductivity, $K_{5}=50 \mathrm{~W} / \mathrm{mK}$
Thickness, $L_{3}=0.15 \mathrm{~m}$
Temperature of outer surface, $T_{o}=20^{\circ} \mathrm{C}$.
Convective heat transfer coefficient, $h=25 \mathrm{~W} / \mathrm{m}^{2} \mathrm{~K}$
Ambient temperature, $T_{\infty}=800^{\circ} \mathrm{C}$

Element (1): It is subjected to both conduction and convection.
Convection is only from left side. Therefore, finite element equation for element (1) can be written as,

$$
\left[K_{1}\right]=\frac{A_{1} K_{1}}{l_{1}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]+h A\left[\begin{array}{ll}
1 & 0 \\
0 & 0
\end{array}\right]
$$

Assume, $A_{1}=1 \mathrm{~m}^{2}$

$$
\begin{aligned}
& =\frac{1 \times 20}{0.3}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]+25 \times 1\left[\begin{array}{ll}
1 & 0 \\
0 & 0
\end{array}\right] \\
& =\left[\begin{array}{cc}
66.66 & -66.66 \\
-66.66 & 66.66
\end{array}\right]+\left[\begin{array}{cc}
25 & 0 \\
0 & 0
\end{array}\right] \\
{\left[K_{1}\right] } & =\left[\begin{array}{cc}
91.66 & -66.66 \\
-66.66 & 66.66
\end{array}\right]
\end{aligned}
$$




Since, on the left end convection takes place, load vector on the left is given by,

$$
\begin{aligned}
& \{\text { Fhe }\}=h T_{\infty} A\left\{\begin{array}{l}
1 \\
0
\end{array}\right\} \\
& \{\text { Fhe }\}=\left\{\begin{array}{c}
1 \\
h T_{\infty} A \\
0
\end{array}\right\} 2
\end{aligned}
$$



Element (2)

$$
\begin{aligned}
{\left[K_{2}\right] } & =\frac{A_{2} K_{2}}{l}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
& =\frac{1 \times 30}{0.15}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
{\left[K_{2}\right] } & =\left[\begin{array}{cc}
200 & -200 \\
-200 & 200
\end{array}\right] 2
\end{aligned}
$$

Element (3)

$$
\begin{aligned}
{\left[K_{3}\right] } & =\frac{A_{3} K_{3}}{l_{3}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
& =\frac{1 \times 50}{0.15}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right] \\
{\left[K_{2}\right] } & =\left[\begin{array}{cc}
333.333 & -333.333 \\
-333.333 & 333.333
\end{array}\right] 3
\end{aligned}
$$



Global stiffness matrix,

$$
\left.\left.\begin{array}{l}
{[K]=\left[K_{1}\right]+\left[K_{2}\right]+\left[K_{3}\right]} \\
1 \\
2
\end{array}\right]=\left[\begin{array}{cccc}
91.66 & -66.66 & 0 & 4 \\
-66.66 & 66.66+200 & -200 & 0 \\
0 & -200 & 200+333.33 & -333.33 \\
0 & 0 & -333.33 & 333.33
\end{array}\right] \begin{array}{l}
1 \\
2 \\
3 \\
4
\end{array}\right][\mathrm{K}]=\left[\begin{array}{cccc}
91.66 & -66.66 & 0 & 0 \\
-66.66 & 266.66 & -200 & 0 \\
0 & -200 & 533.33 & -333.33 \\
0 & 0 & -333.33 & 333.33
\end{array}\right],
$$

Finite element equation for the given composite slab is given by,

$$
\begin{aligned}
& {[K]\{T\}=\{F\}} \\
& {\left[\begin{array}{cccc}
91.66 & -66.66 & 0 & 0 \\
-66.66 & 266.66 & -200 & 0 \\
0 & -200 & 533.33 & -333.33 \\
0 & 0 & -333.33 & 333.33
\end{array}\right]\left[\begin{array}{l}
T_{1} \\
T_{2} \\
T_{3} \\
T_{4} \\
4
\end{array}\right]=\left\{\begin{array}{c}
h T_{\infty} A \\
0 \\
0 \\
0
\end{array}\right\}} \\
& h T_{\infty} A=25 \times 800 \times 1 \\
& h T_{\infty} A=20000 \\
& {\left[\begin{array}{cccc}
91.66 & -66.66 & 0 & 0 \\
-66.66 & 266.66 & -200 & 0 \\
0 & -200 & 533.33 & -333.33 \\
0 & 0 & -333.33 & 333.33
\end{array}\right]\left[\begin{array}{l}
T_{1} \\
T_{2} \\
T_{3} \\
T_{4} \\
4
\end{array}\right]=\left\{\begin{array}{c}
20000 \\
0 \\
0 \\
0
\end{array}\right\}}
\end{aligned}
$$

$$
\begin{aligned}
& T_{1}=304.77^{\circ} \mathrm{C} \\
& T_{2}=119.04^{\circ} \mathrm{C} \\
& T_{3}=57.142^{\circ} \mathrm{C} \\
& T_{4}=20^{\circ} \mathrm{C}
\end{aligned}
$$

Temperature distribution within the wall is given by,

$$
\left\{\begin{array}{l}
T_{1} \\
T_{2} \\
T_{3} \\
T_{4}
\end{array}\right\}=\left\{\begin{array}{c}
304.77 \\
119.04 \\
57.142 \\
20
\end{array}\right\}
$$

One-dimensional Heat Transfer: When the temperature and heat transfer in a system vary only in one direction, it is known as one-dimensional heat transfer. In this, the variation of temperature in other two directions is negligible.
Example: Heat transfer from a glass window is considered to be one-dimensional, as the heat transfer takes place in one direction, whereas in other directions it is zero.

## One dimensional analysis of a fin

## Fin (extended surface)

Fin is a metallic strip of rectangular shape or circular shape and integral with the sufface, through which heat is to be transferred. Fin increases the surface area of heat transfer and also, known as extended surface. They are used on engine cylinders, heat exchanger pipes, etc.

## Heat Transfer Enhancement by Fins



Bare surface
Finned surface

TYPES OF FINS


A fin subjected to conduction and convection

$$
\begin{gathered}
{\left[K_{c}\right]=\frac{A k}{l}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]} \\
{\left[K_{h}\right]=\frac{h P l}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right]} \\
{[K]=\left[K_{c}\right]+\left[K_{h}\right]} \\
{[K]\{T\}=\{F\}}
\end{gathered}
$$

$$
\begin{aligned}
& {[K]\{T\}=\{F\}} \\
& {\left[\frac{A k}{l}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]+\frac{h P l}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right]\right]\left\{\begin{array}{l}
T_{1} \\
T_{2}
\end{array}\right\}=\frac{Q A l+P h T_{\infty} l}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\}}
\end{aligned}
$$

If at free end convection exist

$$
[\text { Khe }]=h A\left[\begin{array}{ll}
0 & 0 \\
0 & 1
\end{array}\right] \quad \begin{aligned}
& \text { (Free end convection matrix) } \\
& \text { - Convection matrix }
\end{aligned}
$$



$$
\begin{gathered}
A=\text { Length } * \text { Thickenss }=l * t \\
P=2 * l(\text { Approximately })
\end{gathered}
$$

| $\mathrm{A} \longrightarrow$ Area of the fin | - $\mathrm{m}^{2}$ |
| :---: | :---: |
| $\mathrm{P} \longrightarrow$ Perimeter of the fin | -m |
| $\mathrm{k} \longrightarrow$ Thermal conductivity of fin | - W/mK |
| $l \longrightarrow$ Length of the fin | -m |
| $h \longrightarrow$ Heat transfer coefficient | - W/m $/ \mathrm{m}^{2}$ |
| $T_{\infty} \longrightarrow$ Atmospheric air temperature | - K |
| $Q \longrightarrow$ Heat Generation | - W |

A metallic fin 20 mm wide and 4 mm thick is attached to a furnace whose wall temperature is $180^{\circ} \mathrm{C}$. The length of the fin is 120 mm . If the thermal conductivity of the material of the fin is 120 mm . If the thermal conductivity of the material of the fin is $350 \mathrm{~W} / \mathrm{m}^{\circ} \mathrm{C}$ and convection coefficient is $9 \mathrm{~W} / \mathrm{m}^{20} \mathrm{C}$, determine the temperature distribution assuming that the tip of the fin is open to the atmosphere and that the ambient temperature is $25^{\circ} \mathrm{C}$.


A fin (length $120 \mathrm{~mm}, 20 \mathrm{~mm}$ wide and 4 mm thick) is attached to a furnace wall temperature of 180 C .
Determine the temperature at the midpoint of the fin assuming the tip of the fin is open to atmosphere, which is at 25 C (take fin's conductivity $350 \mathrm{~W} / \mathrm{mK}$ And convection coefficient of atmosphere 9 $\left.\mathrm{W} / \mathrm{m}^{\wedge} 2 \mathrm{~K}\right)$

## Give that,

Width of the fin, $w=20 \mathrm{~mm}=0.02 \mathrm{~m}$
Thickness of the fin, $t=4 \mathrm{~mm}=4 \times 10^{-3} \mathrm{~m}$
Wall temperature, $T_{1}=180^{\circ} \mathrm{C}$
Length of the fin, $L=120 \mathrm{~mm}=0.12 \mathrm{~m}$
Thermal conductivity, $K=350 \mathrm{~W} / \mathrm{m}^{\circ} \mathrm{C}$
Convection coefficient, $h=9 \mathrm{~W} / \mathrm{m}^{2 \circ} \mathrm{C}$
Ambient temperature, $T_{\infty}=25^{\circ} \mathrm{C}$

The fin is divided into two equal elements.


Figure (2): F.E Modal

Give that,
Width of the fin, $w=20 \mathrm{~mm}=0.02 \mathrm{~m}$
Thickness of the fin, $t=4 \mathrm{~mm}=4 \times 10^{-3} \mathrm{~m}$
Wall temperature, $T_{1}=180^{\circ} \mathrm{C}$
Length of the fin, $L=120 \mathrm{~mm}=0.12 \mathrm{~m}$
Thermal conductivity, $K=350 \mathrm{~W} / \mathrm{m}^{\circ} \mathrm{C}$
Convection coefficient, $h=9 \mathrm{~W} / \mathrm{m}^{20} \mathrm{C}$
Ambient temperature, $T_{\infty}=25^{\circ} \mathrm{C}$

Element (1)


Figure (3): Element (1)

Stiffness matrix for element (1).

$$
\begin{aligned}
{[K] } & =\frac{A_{1} K_{1}}{l_{1}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]+\frac{h p l_{1}}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
A_{1} & =\mathrm{W} \times t=0.02 \times 4 \times 10^{-3}=8 \times 10^{-5} \mathrm{~m}^{2} \\
P & =2(\mathrm{~W}+t)=2\left(0.02+\left(4 \times 10^{-3}\right)\right) \\
P & =0.012 \mathrm{~m}
\end{aligned}
$$

Force vector for element (1).

$$
\left\{\begin{array}{l}
F_{1} \\
F_{2}
\end{array}\right\}=\frac{Q A_{11} l_{1}+p h T_{\infty} l_{1}}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\}
$$

Since, 'Q' is not given in the problem, neglect the term $\left(\frac{Q A_{1} l_{1}}{2}\right)$.

$$
\left\{\begin{array}{c}
F_{1} \\
F_{2}
\end{array}\right\}=\frac{p h T_{\infty} l}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\}
$$

Stiffness matrix for element (1).

$$
\begin{aligned}
{[K] } & =\frac{A_{1} K_{1}}{l_{1}}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]+\frac{h p l_{1}}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
& =\frac{8 \times 10^{-5} \times 350}{0.06}\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]+\frac{9 \times 0.012 \times 0.06}{6}\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
& =0.467\left[\begin{array}{cc}
1 & -1 \\
-1 & 1
\end{array}\right]+\left(1.08 \times 10^{-3}\right)\left[\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right] \\
& =\left[\begin{array}{cc}
0.467 & -0.467 \\
-0.467 & 0.467
\end{array}\right]+\left[\begin{array}{ll}
2.16 \times 10^{-3} & 1.08 \times 10^{-3} \\
1.08 \times 10^{-3} & 2.16 \times 10^{-3}
\end{array}\right] \\
& {\left[K_{1}\right]=\left[\begin{array}{cc}
1 & 2 \\
0.46916 & -0.46592 \\
-0.46592 & 0.46916
\end{array}\right] }
\end{aligned}
$$

$$
\begin{aligned}
\left\{\begin{array}{l}
F_{1} \\
F_{2}
\end{array}\right\} & =\frac{p h T_{\infty} l_{1}}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\} \\
& =\frac{0.012 \times 9 \times 25 \times 0.06}{2}\left\{\begin{array}{l}
1 \\
1
\end{array}\right\}
\end{aligned}
$$

$$
\left\{\begin{array}{l}
F_{1} \\
F_{2}
\end{array}\right\}=\left\{\begin{array}{l}
0.081 \\
0.081
\end{array}\right\}
$$

Finite element equation for element (1).

$$
\begin{aligned}
& {\left[\begin{array}{cc}
0.46916 & -0.46592 \\
-0.46592 & 0.46916
\end{array}\right]\left[\begin{array}{l}
T_{1} \\
T_{2}
\end{array}\right\}=\left\{\begin{array}{l}
F_{1} \\
F_{2}
\end{array}\right\}} \\
& 1 \\
& \left.\left[\begin{array}{cc}
0.46916 & -0.46592 \\
-0.46592 & 0.46916
\end{array}\right] 1 \begin{array}{l}
T_{1} \\
2 \\
T_{2}
\end{array}\right\}=\left\{\begin{array}{l}
0.081 \\
0.081
\end{array}\right\}
\end{aligned}
$$

## Element (2)

Since all the parameters and properties are same finite element equation for element (2) is given by,

$$
\begin{aligned}
& {\left[\begin{array}{cc}
0.46916 & -0.46592 \\
-0.46592 & 0.46916
\end{array}\right]\left\{\begin{array}{c}
T_{2} \\
T_{3}
\end{array}\right\}=\left\{\begin{array}{c}
F_{2} \\
F_{3}
\end{array}\right\}} \\
& 2 \\
& 3 \\
& {\left[\begin{array}{cc}
0.46916 & -0.46592 \\
-0.46592 & 0.46916
\end{array}\right] 2\left[\begin{array}{c}
T_{2} \\
3 \\
T_{3}
\end{array}\right\}=\left\{\begin{array}{l}
0.081 \\
0.081
\end{array}\right\}}
\end{aligned}
$$



## Global finite element equation

$$
\begin{aligned}
& \begin{array}{lll}
1 & 2 & 3
\end{array} \\
& {\left[\begin{array}{ccc}
0.46916 & -0.46592 & 0 \\
-0.46592 & 0.46916+0.46916 & -0.46592 \\
0 & -0.46592 & 0.46916
\end{array}\right] 1\left[\begin{array}{c}
T_{1} \\
2 \\
T_{2} \\
2 \\
T_{3}
\end{array}\right\}=\left\{\begin{array}{c}
0.081 \\
0.081+0.081 \\
0.081
\end{array}\right\}} \\
& {\left.\left[\begin{array}{ccc}
0.46916 & -0.46592 & 0 \\
-0.46592 & 0.93832 & -0.46591 \\
0 & -0.46592 & 0.46916
\end{array}\right] \right\rvert\,\left\{\begin{array}{l}
T_{1} \\
T_{2} \\
T_{3}
\end{array}\right\}=\left\{\begin{array}{l}
0.081 \\
0.162 \\
0.081
\end{array}\right\}}
\end{aligned}
$$

$$
\mathrm{T} 1=180 \mathrm{C}
$$



Writing the 2nd and 3rd rows
$-0.46592 \mathrm{~T} 1+0.93832 \mathrm{~T} 2-0.46591 \mathrm{~T} 3=0.162$
$-0.46592 \mathrm{~T} 2+0.46916 \mathrm{~T} 3=0.081$

$$
\left\{\begin{array}{c}
T_{1} \\
T_{2} \\
T_{3}
\end{array}\right\}=\left\{\begin{array}{c}
180 \\
176.83 \\
175.78
\end{array}\right\}
$$



## ME6603 - FINITE ELEMENT ANALYSIS



## FORMULA BOOK


https://www.slideshare.net/ASHOKKUMAR27088700/me6603-finite-element-analysis-formula-book


## Explain about the use of ANSYS in FEA.

Ans: ANSYS is a finite element modeling package and design analysis tool which is used to solve different problems of engineering based on structural analysis, thermal analysis, CFD analysis, etc.

The documentation of a product designed by ANSYS software consists of commands reference, operations guide, modeling guide, element reference, etc.

Solving the problem using ANSYS is carried out in three stages namely,

ANSYS - Analysis of Systems

## NNSYS

Full Form of APDL is ANSYS Parametric Design Language

1. Preprocessing stage
2. Solution stage
3. Postprocessing stage.
$\Lambda$ ANSYS Multiphysics Utility Menu
```
Eile Select List Plot PlotCtrls WorkPlane Parameters Macro MenuCtrls Help
```



ANSYS loolbar
ANSYS Main Menu
e Preferences

## ANSYS

田 Preprocessor
田 Solution
田 General Postproc
$\boxplus$ TimeHist Postpro
田 Topological Opt
⿴囗 ROM Tool
田 DesignXplorer
田Design Opt
田 Prob Design
T Radiation Opt
图 Session Editor
－Finish

## 1. Preprocessing Stage

In this stage, the problem is described or stated clearly by the following steps,
(i) Defining the key points, lines and areas of elements.
(ii) Defining the element type which includes the elements shape, dimensions, degrees of freedom, etc.
(iii) Defining the material properties like Young's modulus, Poisson's ratio, thermal conductivity, etc.
(iv) Stating the mesh lines, areas or volumes as per the requirement.

2．Solution Stage
After the preprocessing stage，the next stage is solution stage，which includes specifying loads，constraints and obtaining the solution．The sequential steps are as follows，
（i）Stating the loads which may include pressure or points．
（ii）Adding the constraints，either translational or rotational．
（iii）Solving the equations，which are associated with the above steps．
＾ANSYS Mechanical Utility Menu（Load＿Step）

| Eile Select List Plot PlotCtris |  |  |
| :---: | :---: | :---: |
|  |  |  |
| Toolbar |  |  |
| SAVE＿DB | RESUM＿DB | QUIT |


| Main Menu | （2） |
| :---: | :---: |
| 圃 Preferences |  |
| Preprocessor |  |
| $\square$ Solution |  |
| （ Analysis Type |  |
| ＊Define Loads |  |
| $\pm$ Load Step Opts |  |
| （ SE Management（CMS） |  |
| 圆 Results Tracking |  |
| ■ Solve |  |
| 圆 Current LS |  |
| 园 FromLS Files |  |
| （ Manual Rezoning |  |
| （ Multi－field Set Up |  |
| （ ADAMS Connection |  |
| $\pm$ Diagnostics |  |
| Unabridged Menu |  |
| $\boxplus$ General Postproc |  |
| （ TimeHist Postpro |  |
| 田 Prob Design |  |
| ® Radiation Opt |  |
| 里 Session Editor |  |
| 圆 Finish |  |

